

"THREE-INTEGRAL", COLLISION-FREE STATISTICAL MECHANICS AND STELLAR SYSTEMS

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In a recent paper, Innanen and Papp (1977) have introduced and discussed the behaviour of a function they called the "super angular momentum," h_t^2 . This quantity, although not formally an integral of motion in the classical definition, is nevertheless a straightforward generalization of the classical total angular momentum. The super angular momentum, together with the other two classical invariants, the total energy E , and the Z component of the angular momentum h_z lead to the introduction for each homogeneous stellar system of a new collision-free distribution function

$$f = f_0 \exp(-Q)$$

$$\text{where } Q = 2E - 2\beta h_z + \gamma h_t^2$$

and f_0 , β and γ are constants. The function f satisfies the fundamental equation of stellar dynamics and shows that the velocity dispersions of stellar systems cannot be isotropic, but rather that

- (a) the velocity dispersion in the radial direction is constant (i. e. "isothermal") and
- (b) the velocity dispersions in the tangential directions follow the

law

$$\sigma_\theta^2 = \sigma_z^2 \alpha (1 + \gamma r^2)^{-1}$$

That is, the tangential components are "isothermal" at the centre of the system, and decay according to the above equation so as to leave a purely radial distribution in the outer parts.

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REFERENCES

Innanen, K. A., and Papp, K. A.: 1977, *Astron. J.* 82, 322.

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W. B. Burton (ed.), The Large-Scale Characteristics of the Galaxy, 461-462.
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DISCUSSION

Contopoulos: I agree that a third integral of motion is very useful in galactic dynamics, e.g., in explaining the three-axial form of the velocity ellipsoid; this has been done by Barbanis, about 15 years ago. Many other applications have been made by several people since that time. In the case of almost spherical systems a third integral has been calculated by Saaf (a student of Vandervoort), which is a generalization of the total angular momentum.

However your "third integral" includes a function $\psi(t)$, which is an integral in time, calculated along the orbit. Therefore in order to evaluate it you have to calculate the whole orbit of a star. But, according to the standard definition, an integral of motion is a function of local quantities, at any point of an orbit, and remains constant along this orbit. That is, the energy of a star in a time-independent potential is a function of the local position and velocity of the star; we do not have to calculate the whole orbit in order to find it. Therefore what you give as a generalization of the square of the total angular momentum should not be called an integral of motion.

Innanen: The difficulty appears to arise over the semantics of the word "integral". The function we have introduced is not a conservative integral in the "standard" definition. The function does, however, reduce to the square of the total angular momentum integral for spherical systems, and is constant to within 2 or 3% during long numerical experiments. Consequently we have called it the "super-angular momentum". The method makes clear predictions about the variation of velocity dispersions within stellar systems and so is subject to direct observational tests.

Lynden-Bell: For the potential $\psi = \psi_1(r) + \psi_2(\theta)/r^2$ the r and θ motions of one star decouple exactly and the third integral is the square of the angular momentum minus $2\psi_2(\theta)$. Woolley has analyzed the local stars using this integral and has calculated $\psi_1(r)$ and $\psi_2(\theta)$ to fit the Galaxy's potential within some 4 kpc from the Sun (Royal Observatory Annals No. 5, 1971). Dr. Innanen's expression will reduce to this integral because this is a good approximation to the galactic potential. Dr. Innanen's work shows that for a different approximation to the galactic potential an integral of this form still works very well.

Innanen: I am aware of the work mentioned by Dr. Lynden-Bell and agree with his comments.