

A NOTE ON THE MATRIX RENEWAL FUNCTION

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ABSTRACT. The Laplace-Stieltjes Transform $m(s)$ of the matrix renewal function $M(t)$ of a Markov Renewal process is expanded in powers of the argument s , in this paper, by using a generalized inverse of the matrix $I - P_0$, where P_0 is the transition probability matrix of the imbedded Markov chain. This helps in obtaining the values of moments of any order of the number of renewals and also of the moments of the first passage times, for large values of t , the time. All the results of renewal theory are hidden under the Laplacian curtain and this expansion helps to lift this curtain at least for large values of t and is thus useful in predicting the number of renewals.

1. Introduction

Let $P_0 = [p_{ij}]$ ($i, j = 1, 2, \dots, m$) be the transition probability matrix of the imbedded Markov Chain of a Markov Renewal Process (M.R.P.), involving a finite number m of states. Let F_{ij} be the distribution function (d.f.) of the time spent in state i by the process before going to the next transitionstate j ($i, j = 1, 2, \dots, m$). Let J_t denote the state of the process at time t and $N_j(t)$ denote the number of visits to state j in the interval $(0, t)$. The matrix $M(t) = [M_{ij}(t)]$, where

$$(1.1) \quad M_{ij}(t) = E\{N_j(t) | J_0 = i\}$$

is called the matrix renewal function of the M.R.P. Its Laplace-Stieltjes transform (L.-S.T.) is given by Pyke [6], [7] as

$$(1.2) \quad m(s) = \int_0^\infty e^{-st} d_t M(t) = (I - q(s))^{-1} - I$$

where

$$(1.3) \quad Q(t) = [Q_{ij}(t)], Q_{ij}(t) = p_{ij} F_{ij}(t), q(s) = \int_0^\infty e^{-st} d_t Q(t).$$

The asymptotic behaviour of the renewal function in an ordinary renewal process ($m = 1$) is well-known (see for example Cox [1] or Smith [9] and the simplest way of deriving it is to expand its Laplace transform in powers of the argument s ,

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in the neighborhood of $s = 0$ and use Tauberian arguments. Since all higher order moments of the number of renewals in an ordinary renewal process are expressible in terms of the renewal function, their asymptotic behaviour can also be investigated from this, provided a sufficient number of terms in the expansion of the Laplace-Stieltjes transform are retained.

These results and methods can be extended to an M.R.P. The present authors demonstrated this by expanding the L.-S.T. $(I - q(s))^{-1} - I$ given by (1.2) in powers of s . The expansion was obtained from

$$(1.4) \quad I - q(s) = I - P_0 + sP_1 - \frac{s^2}{2!} P_2 + \frac{s^3}{3!} P_3 - \dots$$

where

$$(1.5) \quad P_k = \int_0^\infty x^k dQ(x), \quad (k = 0, 1, 2, \dots)$$

subject, of course, to the existence of these moments. Hereafter we assume that all these moments exist. The main difficulty in this is obviously the singularity of the matrix $I - P_0$. The required expansion of $m(s)$ was obtained in terms of certain matrices H_r ($r = 0, 1, 2, \dots$) obtainable from the adjoint of $I - P_0 + sP_1$. Recently Hunter [2] and Keilson have independently solved this problem by different methods. Hunter uses a generalized inverse (Rao [8]) of $I - P_0$ along with the relationship between the moments of Q_{ij} and G_{ij} , the d.f. of the first passage time of the M.R.P. from state i to state j (Pyke [6], [7]). Keilson uses the spectral decomposition of the matrix $q(s)$ and then finds the limiting values of the first two derivatives of $(I - q(s))^{-1}$ with respect to s , at $s = 0$, which are nothing but the coefficients of powers of s in its expansion. In the present paper, the generalized inverse method is used directly on $(I - q(s))^{-1}$. Investigation of higher moments of $N_j(t)$ requires more than two terms in the expansion of $(I - q(s))^{-1}$ and these can be obtained more easily and directly by the method of this paper.

2. Generalized inverse of $I - P_0$

Pyke [6], [7] has described the classification of the states of a M.R.P. in much the same manner as is done for Markov Chains. Following his terminology, we assume that all the states of the M.R.P. are positive and recurrent. The imbedded Markov Chain is thus regular.

Let $\bar{U}' = [U_1, U_2, \dots, U_m]$ be the vector of the stationary state probabilities ($\sum U_i = 1$) of the imbedded Markov Chain and let $\bar{e}' = [1, \dots, 1]$, a vector of m components. Then it is well-known that

$$(2.1) \quad P_0 \bar{e} = \bar{e}, \quad \bar{U}' P_0 = \bar{U}'.$$

Let

$$(2.2) \quad L = \bar{e} \bar{U}',$$

and

$$(2.3) \quad Z = (I - P_0 + L)^{-1}.$$

Z is known as the fundamental matrix of the Markov Chain (Kemmeny and Snell [4]) and the following results are easily derivable (Hunter [2]).

(i) Z is a generalized inverse of $I - P_0$ i.e.

$$(I - P_0)Z(I - P_0) = (I - P_0)$$

(ii) $P_0Z = ZP_0$

$$(2.4) \quad \text{(iii) } Z\bar{e} = \bar{e}$$

$$\text{(iv) } \bar{U}'Z = \bar{U}'$$

$$\text{(v) } ZL = LZ = L$$

$$\text{(vi) } Z(I - P_0) = I - L.$$

We also define

$$(2.5) \quad k_r = \bar{U}'P_r\bar{e} \quad (r = 1, 2, \dots),$$

so that

$$(2.6) \quad LP_rL = k_rL \quad (r = 1, 2, \dots)$$

We shall also need the following result (Rao 1966).

The general solution of the equation $AX = B$ is

$$(2.7) \quad X = A^-B + (I - A^-A)W$$

where A^- is any generalized inverse of A and W is any arbitrary matrix. This result holds if and only if the *consistency condition*

$$(2.8) \quad AA^-B = B$$

is satisfied.

3. Expansion of $(I - q(s))^{-1}$

It is well-known from the theory of ordinary renewal processes (see for example Keilson [3]) that, if the Markov chain is ergodic and if the P_k 's exist, $(I - q(s))^{-1}$ can be expanded in the form

$$(3.1) \quad \frac{1}{s} A_{-1} + A_0 + sA_1 + s^2A^2 + \dots$$

We shall now show how the A 's can be determined. From (1.4),

$$(3.2) \quad \left(I - P_0 + sP_1 - \frac{s^2}{2!} P_2 + \dots \right) \left(\frac{1}{s} A_{-1} + A_0 + sA_1 + \dots \right) \equiv I.$$

Equating coefficients of s^r on both sides,

$$(3.3) \quad (I - P_0)A_r - \sum_{\alpha=1}^{r+1} \frac{(-1)^\alpha}{\alpha!} P_\alpha A_{r-\alpha} = \delta_{r,0} I, \quad (r = -1, 0, 1, 2, \dots)$$

where $\delta_{r,0}$ is the kronecker delta. Premultiply (3.3) by L and observing that $L(I - P_0) = 0$, we obtain

$$(3.4) \quad LD_{r+1} = 0, \quad (r = -1, 0, 1, \dots),$$

where

$$(3.5) \quad D_{r+1} = \sum_{\alpha=1}^{r+1} \frac{(-1)^\alpha}{\alpha!} P_\alpha A_{r-\alpha} + \delta_{r,0} I.$$

Write (3.3) as an equation in A_r , in the form

$$(3.6) \quad (I - P_0)A_r = D_{r+1} \quad (r = -1, 0, 1, \dots).$$

We shall now verify that the consistency condition (2.8) is satisfied for this equation. For this we need (i) and (vi) of (2.4) viz that Z is a generalized inverse of $I - P_0$ and that $Z(I - P_0) = (I - L)$. The consistency condition is then

$$Z(I - P_0)D_{r+1} = D_{r+1},$$

or which is the same as,

$$(I - L)D_{r+1} = D_{r+1},$$

which is true, on account of (3.4). The *general solution* of (3.6) is, therefore, by (2.7)

$$(3.7) \quad A_r = ZD_{r+1} + LW_r, \quad (r = -1, 0, 1, \dots)$$

where W_r is some arbitrary matrix. We shall now obtain the value of W_r . From (3.4), $LD_{r+2} = 0$ ($r = -2, -1, 0, 1, \dots$), i.e.

$$L \left\{ \sum_{\alpha=1}^{r+2} \frac{(-1)^\alpha}{\alpha!} P_\alpha A_{r+1-\alpha} + \delta_{r+1,0} I_0 \right\} = 0$$

i.e.

$$-LP_1 A_r + \sum_{\alpha=2}^{r+2} \frac{(-1)^\alpha}{\alpha!} LP_\alpha A_{r+1-\alpha} + \delta_{r+1,0} L = 0.$$

Substitute for A_r in this equation from (3.7). We get

$$(3.8) \quad -LP_1(ZD_{r+1} + LW_r) + \sum_{\alpha=2}^{r+2} \frac{(-1)^\alpha}{\alpha!} LP_\alpha A_{r+1-\alpha} + \delta_{r+1,0} L = 0.$$

From (2.2) and (2.5), we obtain

$$\begin{aligned} LP_1 L &= eU'P_1 eU' \\ &= (U'P_1 e)eU' \\ &= k_1 L, \end{aligned}$$

and so, (3.8) yields

$$(3.9) \quad LW_r = -\frac{1}{k_1} LP_1 ZD_{r+1} + \frac{1}{k_1} \sum_{\alpha=1}^{r+1} \frac{(-1)^{\alpha+1}}{(\alpha+1)!} LP_{\alpha+1} A_{r-\alpha} + \frac{1}{k_1} \delta_{r+1,0} L.$$

Substituting this back in (3.7), we get finally the general solution of (3.6) as

$$(3.10) \quad A_r = \sum_{\alpha=1}^{r+1} \frac{(-1)^\alpha}{\alpha!} \left\{ ZP_\alpha - \frac{1}{k_1} LP_1 ZP_\alpha - \frac{1}{k_1(\alpha+1)} LP_{\alpha+1} \right\} A_{r-\alpha} + \delta_{r0} \left(I - \frac{1}{k_1} LP_1 \right) Z + \frac{1}{k_1} \delta_{r+1,0} L, \quad (r = -1, 0, 1, \dots).$$

Since this general solution does not include any arbitrary matrix, it is the only unique solution and in particular,

$$(3.11) \quad A_{-1} = \frac{1}{k_1} L.$$

$$(3.12) \quad A_0 = \left(I - \frac{1}{k_1} LP_1 \right) Z \left(I - \frac{1}{k_1} P_1 L \right) + \frac{k_2}{2k_1^2} L,$$

$$(3.13) \quad A_1 = \left\{ -ZP_1 + \frac{1}{k_1} LP_1 ZP_1 + \frac{1}{2k_1} LP_2 \right\} A_0 + \frac{1}{2k_1} \left\{ ZP_2 - \frac{1}{k_1} LP_1 ZP_2 - \frac{1}{3k_1} LP_3 \right\} L.$$

4. Remarks

This expansion of $(I - q(s))^{-1}$ yields immediately, the following asymptotic expression for the renewal function $M(t)$, by inverting the L.-S.T. for large t :

$$(4.1) \quad M(t) = tA_{-1} + A_0 + 0(1).$$

The Laplace-Stieltjes transform of the second factorial moment $E\{N_j(t)N_j(t) - 1\}$ is given by (Kshirsagar and Gupta [5])

$$(4.2) \quad 2\{(I - q(s))^{-1} - I\}_d \{(I - q(s))^{-1} - I\}$$

where ${}_d B$ stands for the diagonal matrix obtained from a matrix B , by removing all the off-diagonal elements. By using (3.7), (3.1), in (4.2), it can be readily shown that, for large t , the second factorial moment, conditional on $J_0 = i$ is of the form

$$(4.3) \quad t^2 G_1 + t G_2 + G_3 + 0(1).$$

The matrices G_1, G_2, G_3 can be easily obtained in terms of A_{-1}, A_0, A_1, \dots and then in terms of P_k 's and Z, L from (3.7). Extension to higher order moments is also possible in exactly a similarly manner, as their L.-S.T.'s all turn up in terms of $(I - q(s))^{-1}$ and $q(s)$.

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