

## EXTENDABLE TEMPERATURES

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### Abstract

Let  $E$  and  $D$  be open subsets of  $\mathbb{R}^{n+1}$  such that  $\overline{D}$  is a compact subset of  $E$ , and let  $v$  be a supertemperature on  $E$ . We call a temperature  $u$  on  $D$  *extendable by  $v$*  if there is a supertemperature  $w$  on  $E$  such that  $w = u$  on  $D$  and  $w = v$  on  $E \setminus \overline{D}$ . Such a temperature need not be a thermic minorant of  $v$  on  $D$ . We show that either there is a unique temperature extendable by  $v$ , or there are infinitely many. Examples of temperatures extendable by  $v$  include the greatest thermic minorant  $GM_v^D$  of  $v$  on  $D$ , and the Perron–Wiener–Brelot solution of the Dirichlet problem  $S_v^D$  on  $D$  with boundary values the restriction of  $v$  to  $\partial D$ . In the case where these two examples are distinct, we give a formula for producing infinitely many more. Clearly  $GM_v^D$  is the greatest extendable thermic minorant, but we also prove that there is a least one, which is not necessarily equal to  $S_v^D$ .

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### 1. Introduction

Given an open set  $E$  in  $\mathbb{R}^{n+1}$ , a function  $u \in C^{2,1}(E)$  that satisfies the standard heat equation on  $E$  is called a *temperature*. If

$$W(x, t) = \begin{cases} (4\pi t)^{-n/2} \exp(-|x|^2/4t) & \text{if } t > 0, \\ 0 & \text{if } t \leq 0, \end{cases}$$

then  $W$  is a temperature on  $\mathbb{R}^{n+1} \setminus \{0\}$ . For any point  $p_0 = (x_0, t_0) \in \mathbb{R}^{n+1}$  and any positive number  $c$ , the set

$$\Omega(p_0; c) = \Omega(x_0, t_0; c) = \{(y, s) \in \mathbb{R}^{n+1} : W(x_0 - y, t_0 - s) > (4\pi c)^{-n/2}\}$$

is called the *heat ball* with *centre*  $(x_0, t_0)$  and *radius*  $c$ . The heat ball is of increasing importance and can now be found in several books, including [1, 3, 4, 6]. Temperatures can be characterised in terms of mean values over heat balls, in that a function  $u \in C^{2,1}(E)$  is a temperature if and only if

$$u(x_0, t_0) = (4\pi c)^{-n/2} \iint_{\Omega(x_0, t_0; c)} \frac{|x_0 - x|^2}{4(t_0 - t)^2} u(x, t) dx dt$$

whenever  $\overline{\Omega}(x_0, t_0; c) \subseteq E$ .

An extended real-valued function  $v$  on  $E$  is called a *supertemperature* on  $E$  if it satisfies the following four conditions:

- ( $\delta_1$ )  $-\infty < v(p) \leq +\infty$  for all  $p \in E$ ;
- ( $\delta_2$ )  $v$  is lower semicontinuous on  $E$ ;
- ( $\delta_3$ )  $w$  is finite on a dense subset of  $E$ ;
- ( $\delta_4$ ) given any point  $(x_0, t_0) \in E$  and positive number  $\epsilon$ , there is a positive number  $c < \epsilon$  such that the closed heat ball  $\bar{\Omega}(x_0, t_0; c) \subseteq E$  and the inequality

$$v(x_0, t_0) \geq (4\pi c)^{-n/2} \iint_{\Omega(x_0, t_0; c)} \frac{|x_0 - x|^2}{4(t_0 - t)^2} u(x, t) dx dt$$

holds.

If  $v$  is a supertemperature on  $E$ ,  $D$  is an open subset of  $E$ , and  $u$  is a temperature such that  $u \leq v$  on  $D$ , then  $u$  is called a *thermic minorant* of  $v$  on  $D$ .

Given an open set  $D$  such that  $\bar{D}$  is a compact subset of  $E$ , and a supertemperature  $v$  on  $E$ , it is well known that the restriction of  $v$  to  $D$  can be replaced by a temperature in such a way that the resultant function, perhaps with some modification on  $\partial D$ , is a supertemperature on  $E$ . We shall study this phenomenon, using the following terminology.

**DEFINITION 1.1.** Let  $E$  and  $D$  be open sets such that  $\bar{D}$  is a compact subset of  $E$ , and let  $v$  be a supertemperature on  $E$ . If  $u$  is a temperature on  $D$  such that the function  $w$  defined by

$$w = \begin{cases} u & \text{on } D, \\ v & \text{on } E \setminus \bar{D}, \end{cases}$$

can be extended to a supertemperature on  $E$ , we say that  $u$  is *extendable by  $v$  (to  $E$ )*.

Our starting point is a pair of known results. The first is the following corollary of [7, Theorem 2.5]. In this, a function  $f$  on  $\partial D$  is called *resolutive* if it has a Perron–Wiener–Brelot solution to the generalised Dirichlet problem, in the sense of [6]. That solution is denoted by  $S_f^D$ . Throughout this paper, our notation and terminology follow [6].

**LEMMA 1.2.** Let  $E$  and  $D$  be open sets such that  $\bar{D}$  is a compact subset of  $E$ , and let  $v$  be a supertemperature on  $E$ . Then the restriction of  $v$  to  $\partial D$  is resolutive for  $D$ , and the function  $w$  defined by

$$w = \begin{cases} S_v^D & \text{on } D, \\ v & \text{on } E \setminus \bar{D}, \end{cases}$$

can be extended to a supertemperature  $\bar{w}$  majorised by  $v$  on  $E$ .

The second result is not quite [7, Corollary 3.3] but can be proved by an almost identical method.

**LEMMA 1.3.** *Let  $E$  and  $D$  be open sets such that  $\bar{D}$  is a compact subset of  $E$ , let  $v$  be a supertemperature on  $E$ , and let  $GM_v^D$  be the greatest thermic minorant of  $v$  on  $D$ . Then the function  $w$  defined by*

$$w = \begin{cases} GM_v^D & \text{on } D, \\ v & \text{on } E \setminus \bar{D}, \end{cases}$$

*can be extended to a supertemperature  $\bar{w}$  majorised by  $v$  on  $E$ .*

Thus each of  $S_v^D$  and  $GM_v^D$  is extendable by  $v$  to  $E$ .

We note that the temperature  $u$ , in the definition of ‘extendable’, is not necessarily a thermic minorant of  $v$  on  $D$ . For if  $S_v^D \neq GM_v^D$  and we use the function  $\bar{w}$  of Lemma 1.2 in place of the original supertemperature  $v$ , then  $GM_v^D$  is extendable by  $\bar{w}$ , by Lemma 1.3, but is not a thermic minorant of  $\bar{w}$  on  $D$ .

It is known that, for some open subsets  $D$  such as heat balls, there is only one temperature  $u$  on  $D$  that is extendable by a given supertemperature  $v$ . See [5, Theorem 6] for details. On the other hand, if there are two distinct temperatures  $u_1$  and  $u_2$  on  $D$  that are extendable by  $v$ , then there are infinitely many. This is because, whenever  $0 < \alpha < 1$ , the temperature  $\alpha u_1$  is extendable by  $\alpha v$ , and  $(1 - \alpha)u_2$  is extendable by  $(1 - \alpha)v$ , so that  $\alpha u_1 + (1 - \alpha)u_2$  is extendable by  $\alpha v + (1 - \alpha)v = v$ .

Even if  $S_v^D = GM_v^D$ , there may still be infinitely many temperatures on  $D$  that are extendable by  $v$ , as the following example shows.

**EXAMPLE 1.4.** We take  $E = \mathbb{R}^{n+1}$  and  $v$  the characteristic function of  $\mathbb{R}^n \times ]0, \infty[$ . Given any ball  $B$  in  $\mathbb{R}^n$ , we take  $D = B \times (]-1, 1[ \cup ]1, 2[)$ . Since the set  $B \times \{-1, 1\}$  is a null set for the Riesz measure associated with  $v$ , [8, Theorem 12] implies that  $S_v^D = GM_v^D$ . By Lemma 1.2, the function  $w_1$  defined by

$$w_1 = \begin{cases} S_v^D & \text{on } D, \\ v & \text{on } E \setminus \bar{D}, \end{cases}$$

can be extended to a supertemperature on  $E$ . We now define  $D_+ = B \times ]1, 2[$  and  $D_- = B \times ]-1, 1[$ . Then  $S_v^D = S_v^{D_-}$  on  $D_-$  and  $S_v^D = S_v^{D_+}$  on  $D_+$ . If  $C = B \times ]-1, 2[$ , then  $\bar{C} = \bar{D}$  and the function

$$w_2 = \begin{cases} S_v^C & \text{on } C, \\ v & \text{on } E \setminus \bar{C} = E \setminus \bar{D}, \end{cases}$$

can be extended to a supertemperature on  $E$ . Furthermore  $S_v^C = S_v^{D_-}$  on  $D_-$ , but  $S_v^C < S_v^{D_+}$  on  $D_+$  because  $S_v^C < 1 = v$  on  $B \times \{1\}$ . Therefore  $w_2 < w_1$  on  $D_+$ . Thus  $S_v^D$  and  $S_v^C$  are both temperatures on  $D$  that are extendable by  $v$ , but they are not equal. Hence there are infinitely many such temperatures.

### 2. The least extendable temperature

Example 1.4 also shows that  $S_v^D$  may not be the least temperature on  $D$  that is extendable by  $v$ . However this is because of the choice of  $v$  and a different choice can satisfy this equation. There is a smallest supertemperature  $v_0$  on  $E$  such that  $v_0 = v$  on  $E \setminus \bar{D}$ , and  $S_{v_0}^D = v_0$  on  $D$ , as we now show.

**THEOREM 2.1.** *Let  $E$  and  $D$  be open sets such that  $\bar{D}$  is a compact subset of  $E$ , and let  $v$  be a supertemperature on  $E$ . Then the class  $\mathcal{F}$  of all supertemperatures  $w$  on  $E$  such that  $w = v$  on  $E \setminus \bar{D}$  has a minimal element  $v_0$  for which  $S_{v_0}^D = v_0$  on  $D$ .*

**PROOF.** If  $B$  is an open superset of  $\bar{D}$  such that  $\bar{B}$  is a compact subset of  $E$ , then  $v$  has a lower bound  $K$  on  $\partial B$ . Now the minimum principle shows that the restrictions to  $B$  of all members of  $\mathcal{F}$  are lower bounded by  $K$ . Thus the function  $u = \inf \mathcal{F}$  is locally lower bounded on  $E$ , so that its lower semicontinuous smoothing  $\widehat{u}$  is a supertemperature on  $E$  which equals  $u$  wherever the latter is lower semicontinuous, by [6, Theorem 7.13]. Therefore  $\widehat{u} = v$  on  $E \setminus \bar{D}$ , so that  $\widehat{u} \in \mathcal{F}$  and  $\mathcal{F}$  has a minimal element. Applying Lemma 1.2, we see that the function  $w_0$ , defined by

$$w_0 = \begin{cases} S_{\widehat{u}}^D & \text{on } D, \\ \widehat{u} = v & \text{on } E \setminus \bar{D}, \end{cases}$$

can be extended to a supertemperature majorised by  $\widehat{u}$  on  $E$ . That extension belongs to  $\mathcal{F}$ , so that it also majorises  $\widehat{u}$  and hence the two are equal. □

### 3. A formula for extendable temperatures

For the case where  $S_v^D$  and  $GM_v^D$  are distinct, we can give a formula for an infinity of temperatures on  $D$  that are extendable by  $v$ . The key to this is [8, Theorem 7(b)], which we state below as Lemma 3.1.

We use the following notations. If  $\nu$  is a nonnegative Borel measure on an open set  $E$  and  $A$  is a  $\nu$ -measurable set, we denote the restriction of  $\nu$  to  $A$  by  $\nu_A$ . If  $D$  is an open subset of  $E$ , we denote the extended Green function of  $D$  by  $G_D^-$  (see [2, 8] for details), and define the function  $G_D^- \nu$  by

$$G_D^- \nu(p) = \int_E G_D^-(p; q) d\nu(q)$$

for all  $p \in D$ .

We require a classification of the boundary points of  $E$  in which we use the following notations for upper and lower half-balls. Given any point  $p_0 = (x_0, t_0)$  in  $\mathbb{R}^{n+1}$  and  $r > 0$ , we put  $H(p_0, r) = \{(x, t) : |x - x_0|^2 + (t - t_0)^2 < r^2, t < t_0\}$  and  $H^*(p_0, r) = \{(x, t) : |x - x_0|^2 + (t - t_0)^2 < r^2, t > t_0\}$ . Let  $q$  be a boundary point of the bounded open set  $D$ . In our classification of boundary points, we always suppose that the boundary of  $D$  does not contain any polar set whose union with  $D$  would give another open set. We call  $q$  a *normal* boundary point if every lower half-ball centred at

$q$  meets the complement of  $D$ . If this condition fails, and also for every  $r > 0$  we have  $H^*(q, r) \cap D \neq \emptyset$ , then  $q$  is called a *semisingular* boundary point. The set of all normal boundary points of  $D$  is denoted by  $\partial_n D$  and that of all semisingular points by  $\partial_{ss} D$ . A similar classification is made relative to the adjoint equation, by interchanging  $H$  and  $H^*$  throughout. This leads, in particular, to the idea of a point  $q \in \partial D$  being a *cothermal* normal boundary point if every upper half-ball centred at  $q$  meets the complement of  $D$ ; the set of all such points is denoted by  $\partial_n^* D$ . A point  $q \in \partial_n E$  is called *regular* if  $\lim_{p \rightarrow q} S_f^E(p) = f(q)$  for all  $f \in C(\partial E)$ ; a point  $q \in \partial_n^* E$  is called *coregular* if this holds relative to the adjoint equation. A point  $q = (y, s) \in \partial_{ss} E$  is called *regular* if  $\lim_{(x,t) \rightarrow (y,s+)} S_f^E(x, t) = f(y, s)$  for all  $f \in C(\partial E)$ .

**LEMMA 3.1.** *Let  $E$  and  $D$  be open sets such that  $\bar{D}$  is a compact subset of  $E$ , and let  $Y$  be the complement in  $\partial D$  of the set of coregular points of  $\partial_n^* D$ . Let  $\nu$  be a supertemperature on  $E$ , and let  $\nu$  be its associated Riesz measure. Then we have  $GM_\nu^D = S_\nu^D + G_{\bar{D}}^- \nu_Y$  on  $D$ .*

We can now give our formula for temperatures on  $D$  that are extendable by  $\nu$ , which gives an infinity of such functions in the case where  $S_\nu^D \neq GM_\nu^D$ .

**THEOREM 3.2.** *Let  $E$  and  $D$  be open sets such that  $\bar{D}$  is a compact subset of  $E$ , and let  $Y$  be the complement in  $\partial D$  of the set of coregular points of  $\partial_n^* D$ . Let  $\nu$  be a supertemperature on  $E$  and let  $\nu$  be its associated Riesz measure. If  $\mu$  is a measure on  $E$  such that  $0 \leq \mu \leq \nu_Y$ , then the function  $w$  defined on  $E \setminus \partial D$  by*

$$w = \begin{cases} S_\nu^D + G_{\bar{D}}^- \mu & \text{on } D, \\ \nu & \text{on } E \setminus \bar{D}, \end{cases}$$

is a temperature on  $D$  and can be extended to a supertemperature majorised by  $\nu$  on  $E$ . Moreover,

$$\lim_{p \rightarrow q, p \in D} G_{\bar{D}}^- \mu(p) = 0$$

whenever  $q$  is a regular point of  $\partial_n D \setminus \bar{Y}$ , and

$$\lim_{p \rightarrow q+, p \in D} G_{\bar{D}}^- \mu(p) = 0$$

whenever  $q$  is a regular point of  $\partial_{ss} D \setminus \bar{Y}$ .

**PROOF.** To prove the extendability, we first suppose that  $\nu \geq 0$ , so that we can write  $\nu = G_E \nu + GM_\nu^E$  on  $E$ . Given a measure  $\mu$  on  $E$  such that  $0 \leq \mu \leq \nu_Y$ , we put  $\nu_1 = G_E(\nu - \mu)$  and  $\nu_2 = G_E \mu + GM_\nu^E$  on  $E$ . Then  $\nu_1$  and  $\nu_2$  are supertemperatures on  $E$  and  $\nu = \nu_1 + \nu_2$ . By [8, Lemma 6],  $G_{\bar{D}}^- \mu$  is a temperature on  $D$  because  $\mu$  is supported in  $\partial D$ . By Lemma 1.2, the function  $w_1$  defined by

$$w_1 = \begin{cases} S_{\nu_1}^D & \text{on } D, \\ \nu_1 & \text{on } E \setminus \bar{D}, \end{cases}$$

can be extended to a supertemperature majorised by  $v_1$  on  $E$ . Moreover, by Lemma 1.3, the function  $w_2$  defined by

$$w_2 = \begin{cases} GM_{v_2}^D & \text{on } D, \\ v_2 & \text{on } E \setminus \bar{D}, \end{cases}$$

can be extended to a supertemperature majorised by  $v_2$  on  $E$ . Furthermore, Lemma 3.1 shows that  $GM_{v_2}^D = S_{v_2}^D + G_D^- \mu$  on  $D$ , so that  $w_1 + w_2 = S_v^D + G_D^- \mu = w$  on  $D$ , and hence  $w_1 + w_2 = w$  on  $E \setminus \partial D$ . It follows that  $w$  can be extended to a supertemperature majorised by  $v_1 + v_2 = v$  on  $E$ .

In the general case, we choose an open superset  $C$  of  $\bar{D}$  such that  $\bar{C}$  is a compact subset of  $E$ . Then  $v$  is lower bounded on  $C$ , so that we can apply the case just proved to  $v - m$  on  $C$ , where  $m = \inf_C v$ , noting that  $\nu_C$  is the Riesz measure associated with  $v - m$  on  $C$ . Thus, if  $\mu$  is a measure on  $E$  such that  $0 < \mu \leq \nu_Y$ , then the function  $w_3$  defined by

$$w_3 = \begin{cases} S_{v-m}^D + G_D^- \mu_C & \text{on } D, \\ v - m & \text{on } C \setminus \bar{D}, \end{cases}$$

is a temperature on  $D$  and can be extended to a supertemperature majorised by  $v - m$  on  $C$ . The addition of  $m$  to  $w_3$  gives the restriction of  $w$  to  $C \setminus \partial D$  and the extendability follows.

To prove the last part, we first put  $\nu_\mu = G_E \mu$ . Then, by [8, Theorem 3],

$$\begin{aligned} G_D^- \mu(p) &= \int_E \left( G_E(p; q) - \int_{\partial D} G_E(\cdot; q) d\omega_p^D \right) d\mu(q) \\ &= G_E \mu(p) - \int_{\partial D} G_E \mu d\omega_p^D \\ &= \nu_\mu(p) - S_{\nu_\mu}^D(p) \end{aligned}$$

for all  $p \in D$ . The function  $\nu_\mu$  is a temperature on  $E \setminus \bar{Y}$ , so that it is continuous on  $\partial D \setminus \bar{Y}$ . Therefore

$$\lim_{p \rightarrow q, p \in D} S_{\nu_\mu}^D(p) = \nu_\mu(q)$$

whenever  $q$  is a regular point of  $\partial_n D \setminus \bar{Y}$ , and

$$\lim_{p \rightarrow q^+, p \in D} S_{\nu_\mu}^D(p) = \nu_\mu(q)$$

whenever  $q$  is a regular point of  $\partial_{ss} D \setminus \bar{Y}$ . The result follows. □

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