

RESEARCH ARTICLE

# Two-micron all-fiberized passively mode-locked fiber lasers with high-energy nanosecond pulse

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## Abstract

We report on mode-locked thulium-doped fiber lasers with high-energy nanosecond pulses, relying on the transmission in a semiconductor saturable absorber (SESA) and a carbon nanotube (CNTs-PVA) film separately. A section of an SMF–MMF–SMF structure multimode interferometer with a transmission peak wavelength of  $\sim 2003$  nm was used as a wavelength selector to fix the laser wavelength. When the SESA acted as a saturable absorber (SA), the mode-locked fiber laser had a maximum output power of  $\sim 461$  mW with a pulse energy of  $\sim 0.14$   $\mu$ J and a pulse duration of  $\sim 9.14$  ns. In a CNT-film-based mode-locked fiber laser, stable mode-locked pulses with the maximum output power of  $\sim 46$  mW, pulse energy of  $\sim 26.8$  nJ and pulse duration of  $\sim 9.3$  ns were obtained. To the best of our knowledge, our experiments demonstrated the first 2  $\mu$ m region ‘real’ SA-based dissipative soliton resonance with the highest mode-locked pulse energy from a ‘real’ SA-based all-fiberized resonator.

**Keywords:** high pulse energy; mode-locking; nanosecond pulse; Tm-doped fiber laser

## 1. Introduction

Mode-locked thulium-doped fiber lasers (TDFLs), which operate in the 2  $\mu$ m spectral region with pulse duration ranging from nanoseconds to femtoseconds, have attracted a lot of interest due to their applications in mid-infrared supercontinuum generation, remote sensing, laser processing, medicine and free-space communication<sup>[1–10]</sup>. In the 2  $\mu$ m regime, research efforts associated with mode-locked fiber lasers so far have mainly aimed at ultra-short pulse ( $\sim$ fs or  $\sim$ ps pulse) generation. The nanosecond-scale mode-locked pulses have received less attention. In effect, nanosecond pulses with high pulse energy have played an important role in many potential applications, such as laser cutting and welding, laser ablation for surface cleaning and as a pump source for mid-infrared pulse laser generation<sup>[11]</sup>.

In general,  $Q$ -switching and gain switching are two well-known techniques for nanosecond pulse generation. Usually,

it is difficult for passively  $Q$ -switched fiber lasers to produce pulses with several nanoseconds duration owing to the limitation of cavity length and modulation depth of the saturable absorber (SA)<sup>[12]</sup>. As for actively  $Q$ -switched fiber lasers, stable nanosecond pulses can be obtained. Nevertheless, a complex electronics controlling system is needed, which highly increases the complexity of the fiber laser. Gain-switched fiber laser is another alternative, but a high-energy pulse pump source is a challenge for 2  $\mu$ m gain-switched pulse generation<sup>[13, 14]</sup>. Furthermore, compared with the mode-locked pulse, the  $Q$ -switched pulse and gain-switched pulse have poorer stability and cannot be reshaped since the phase of the pulse of  $Q$ -switching or gain switching is random.

In fact, passively mode-locked fiber lasers also can generate nanosecond pulses with high energy<sup>[15]</sup> in which dissipative soliton resonance (DSR) has been proved as an effective method. Chang *et al.* theoretically predicted that the DSR can be formed in the anomalous dispersion regime<sup>[16]</sup>. In the following year, the existence of

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high-energy nanosecond DSR pulses in long cavity passively mode-locked erbium-doped fiber lasers with anomalous dispersion was experimentally reported by Li *et al.*<sup>[17]</sup>. Relying on the nonlinear polarization rotation (NPR) technique, they obtained the output pulse energy of  $\sim 79.5$  nJ at a pulse repetition rate of  $\sim 278$  kHz, and the pulse duration could be tuned from  $\sim 12.8$  ns to  $\sim 155.4$  ns. After that, several works about DSR pulse generation in anomalous dispersion have been reported. Nevertheless, most of the works are focused on the  $1.5$   $\mu\text{m}$  region<sup>[18–21]</sup>. As for the  $2$   $\mu\text{m}$  region, there are a few reports about the nanosecond passively mode-locked fiber laser. In 2013, Fu *et al.* realized a passively mode-locked Tm, Ho-codoped fiber laser with a pulse duration of  $\sim 122$  ns and pulse energy of  $\sim 35.2$  nJ by utilizing graphene as an SA and lengthening the cavity to  $\sim 213$  m<sup>[22]</sup>. Based on NPR and  $\sim 460$  m long cavity length, fundamental mode-locked pulses with a pulse duration of  $\sim 304$  ns were obtained by Wang *et al.*<sup>[23]</sup>. By adding a section of  $\sim 40$  m long single mode fiber (SMF) in a simple linear cavity TDFL, Liu *et al.* obtained self-mode-locked pulses with a pulse duration of  $\sim 40$  ns and pulse energy of  $\sim 32.7$  nJ<sup>[24]</sup>. Recently, relying on a graphene SA and a narrow bandwidth fiber Bragg grating (FBG), Wang *et al.* realized a nanosecond-scale mode-locked TDFL with  $\sim 54$  m long cavity length, but the laser has a relatively low pulse energy of  $\sim 1.5$  nJ<sup>[25]</sup>. In our previous work, by utilizing a nonlinear optical loop mirror (NOLM) and an F-P fiber interferometer with a  $\sim 1000$  m long cavity, stable pulses with pulse duration ranging from  $\sim 2.4$  ns to  $\sim 21.2$  ns and a maximum pulse energy of  $\sim 1.5$   $\mu\text{J}$  were obtained<sup>[26]</sup>. For DSR in the  $2$   $\mu\text{m}$  region, Xu *et al.* reported a DSR TDFL with net normal dispersion<sup>[27]</sup>. Zhao *et al.* first reported a nanosecond DSR pulse in TDFL with anomalous dispersion based on nonlinear amplification loop mirror (NALM). With a figure-of-eight configuration resonator, the direct output pulse energy is  $\sim 40.5$  nJ<sup>[28]</sup>. With a figure-9 cavity configuration, Kharitonov *et al.* demonstrated a

DSR mode-locked TDFL with an average output power of 670 mW and pulse energy of 400 nJ<sup>[29]</sup>. The highest average output power of DSR from the TDFL oscillator is  $\sim 1.4$  W with pulse energy of  $\sim 353$  nJ, which was obtained by Du *et al.* with an NOLM-based mode-locked  $\sigma$ -shaped cavity TDFL<sup>[30]</sup>. Nevertheless, we noted that the high-energy DSR nanosecond pulses in the  $2$   $\mu\text{m}$  region were almost based on the ‘artificial’ SAs, such as NOLM and NALM. There are a few experimental reports of ‘real’ SA-based DSR TDFL.

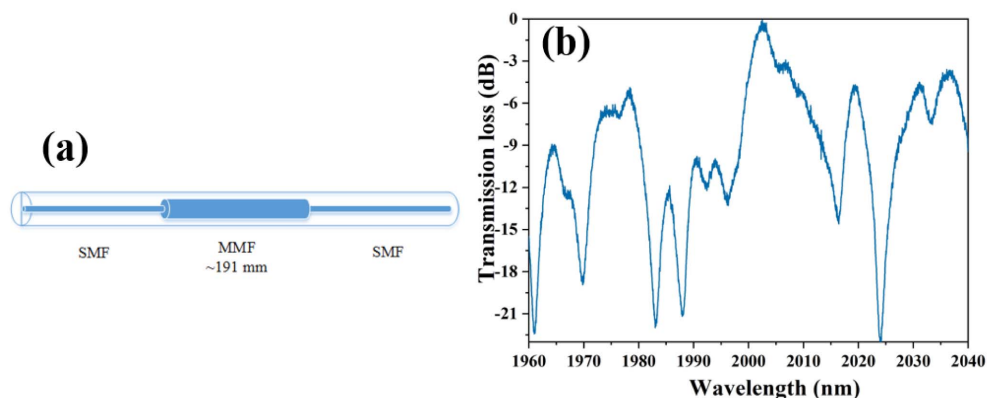
In this paper, we will present high pulse energy nanosecond-scale passively mode-locked TDFLs with semiconductor saturable absorbers (SESAs) and carbon nanotubes (CNTs). A stable rectangular-shape DSR pulse with a maximum pulse energy up to  $\sim 0.14$   $\mu\text{J}$  was first obtained in ‘real’ SA-based mode-locked TDFL. Moreover, our experimental results verified that the SMF-MMF-SMF (SMS, SMF: SMF-28, Corning; MMF: multimode fiber, AFS105/125Y) fiber structure multimode interferometer (MMI)<sup>[31, 32]</sup> can be used as an effective wavelength filter for mode-locked TDFL.

## 2. Experimental setup and result

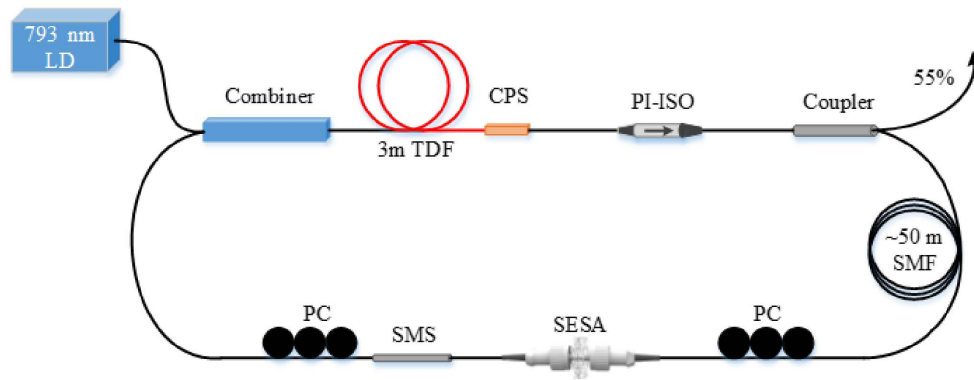
In our experiment, the SMS fiber structure MMI filter was constructed by fusion splicing two pieces of SMFs at both ends of a section of  $\sim 191$  mm long MMF. The structure configuration and transmission characteristics of the fiber MMI are shown in Figures 1(a) and 1(b), respectively. The main transmission peak is at  $\sim 2003$  nm with a transmission loss of  $\sim 0.5$  dB, as shown in Figure 1(b).

### 2.1. SESA-based nanosecond mode-locked fiber laser

The experimental setup for an SESA-based nanosecond mode-locked TDFL is schematically shown in Figure 2. The gain fiber is a section of  $\sim 3$  m length single mode



**Figure 1.** SMS fiber structure MMI: (a) structure configuration; (b) transmission characteristics<sup>[31]</sup>.



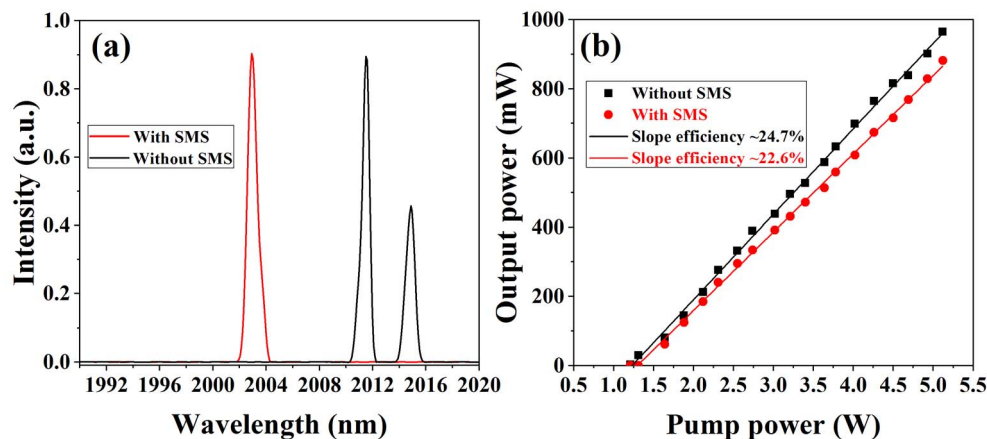
**Figure 2.** Schematic configuration of SESA-based nanosecond TDFL. LD, laser diode; PC, polarization controller; CPS, cladding power stripper; SMF, single mode fiber.

double-clad Tm-doped fiber (SM-TDF-10P/130-HE), which is pumped by a fiber-pigtailed 793 nm multimode laser diode (LD) with a maximum output power of  $\sim 12$  W via a  $(2 + 1) \times 1$  pump/signal combiner. The inner cladding diameter of the TDF is  $\sim 130$   $\mu\text{m}$  with absorption of  $\sim 3$  dB/m at 793 nm. The unabsorbed pump light will effectively strip by a self-made cladding power stripper (CPS) behind the TDF. To ensure the unidirectional propagation of the laser in the cavity, a polarization independent optical isolator with insertion loss of  $\sim 2.6$  dB was used in the cavity. After the isolator, an optical coupler with an output ratio of  $\sim 55\%$  was used for laser outputting. The SESA and SMS fiber structure MMI were placed between two polarization controllers (PCs), which act as a mode locker and a wavelength selector, respectively. The SESA (SA-2000-25-10ps-x, BATOP) has the modulation depth, transmittance and saturation fluence of  $\sim 15\%$ ,  $\sim 74\%$  and  $\sim 2$   $\text{mJ}/\text{cm}^2$ , respectively, at  $\sim 2000$  nm. The PCs are utilized to adjust the polarization state of the light and initiate the mode-locking. Moreover, a section of a  $\sim 50$  m long SMF was added into the cavity to manage the dispersion and implement the nanosecond mode-locked

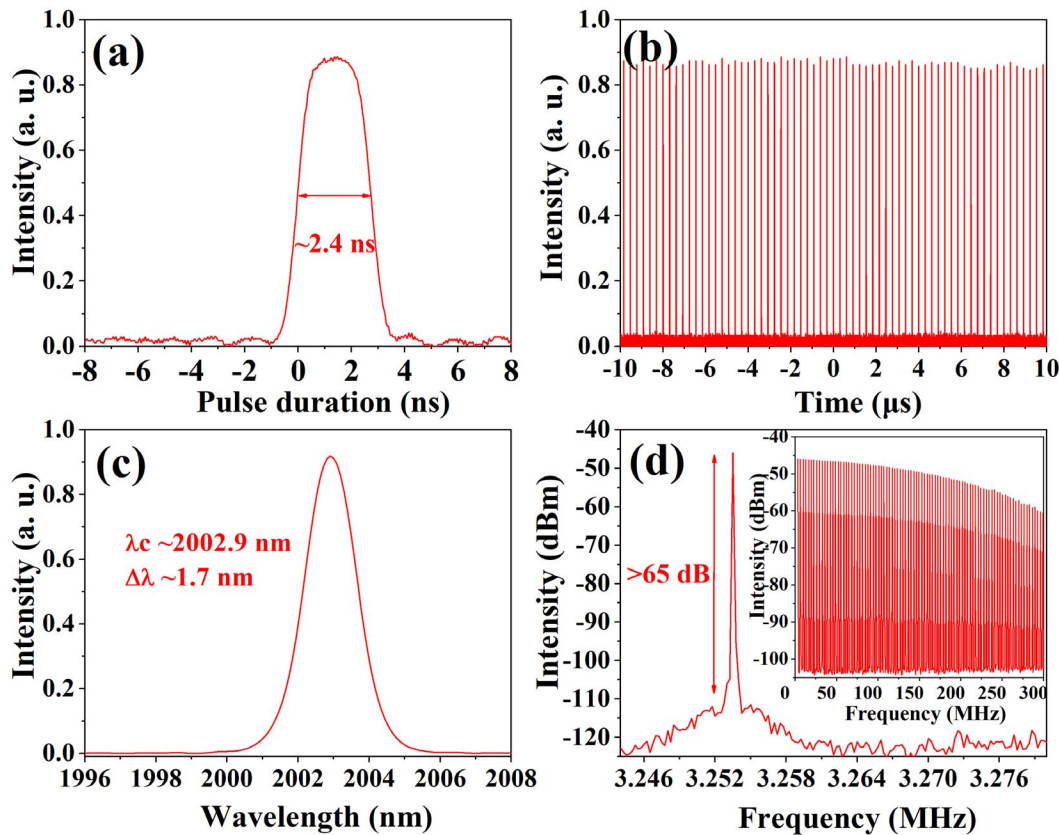
pulse. Thus, a simple ring cavity laser was formed with a total cavity length and a dispersion of  $\sim 63.5$  m and  $-5.8$   $\text{ps}^2$ , respectively.

In the beginning, to verify that the SMS fiber structure MMI can function as an available wavelength selector, the SESA was not inserted into the cavity. The output properties of the fiber laser with and without the SMS fiber structure MMI are shown in Figure 3. When the fiber MMI was inserted into the cavity, the lasing wavelength was  $\sim 2003$  nm and the output efficiency of the fiber laser only had a slight decrease (from  $\sim 24.7\%$  to  $\sim 22.6\%$ ), which means the fiber MMI can be used as an effective wavelength selector. Moreover, no pulses were observed by increasing the pump power and fully manipulating the PCs, which verifies the fiber MMI is not responsible for mode-locking<sup>[33]</sup>.

Then, the SESA was sandwiched between the FC/PC ferrules, and the threshold of the mode-locked operation was about  $\sim 2.06$  W. The typical mode-locked output properties at pump power of  $\sim 2.83$  W are exhibited in Figure 4. Stable mode-locked pulses are shown in Figures 4(a) and 4(b), which were measured by a photodetector (ET-5000,  $\sim 28$  ps



**Figure 3.** Output properties of the fiber laser without SA: (a) typical output spectra with and without the SMS fiber structure MMI; (b) output efficiencies of the fiber laser with and without the SMS fiber structure MMI.



**Figure 4.** Typical mode-locked properties of the fiber laser at pump power of  $\sim 2.83$  W: (a), (b) mode-locked pulses; (c) output spectrum; (d) output RF spectra.

rise/fall time) together with a 20 GS/s high-speed oscilloscope with 1-GHz bandwidth (Tektronix, DPO 7104C digital phosphor oscilloscope). The mode-locked pulses have a pulse duration of  $\sim 2.4$  ns with a rectangular shape, as shown in Figure 4(a). Since the response threshold of our detecting system is less than 1 ns, we confirm that the pulse durations are truly detected. In addition, we also measured the autocorrelation traces of the nanosecond pulses with a time span of 50 ps by a commercial autocorrelator (APE Pulsecheck). But no sub-pulse was observed, which further proves that the nanosecond mode-locked pulses do not have complicated ultra-short sub-pulses. Figure 4(c) shows the output spectrum of the mode-locked fiber laser. The central wavelength is  $\sim 2002.9$  nm with 3 dB bandwidth of  $\sim 1.7$  nm, corresponding to the transmission spectrum of the SMS fiber structure MMI. Thus the time-band product is  $\sim 305$ , which indicates a serious chirp. With the spectrum analyzer (Agilent, E4407B), the radio-frequency (RF) spectra were measured and plotted in Figure 4(d). The repetition frequency of the mode-locked pulse is  $\sim 3.254$  MHz with a signal-to-noise ratio of  $\sim 65$  dB, which verifies that the laser operated at a stable fundamental frequency mode-locked state. The harmonic RF spectra with 300 MHz bandwidth are shown in the inset of Figure 4(d).

In the experiment, we noted that the mode-locked pulse duration was increased with the pump powers, and the details are shown in Figure 5. Figure 5(a) shows the evolution of the pulse envelopes. With the increasing pump power from  $\sim 2.06$  W to  $\sim 4.40$  W, the pulse duration was broadened from  $\sim 949$  ps to  $\sim 9.2$  ns, as given in Figure 5(b). Moreover, the average output powers of the mode-locked pulses were measured, and the corresponding pulse energies and peak powers were calculated, as depicted in Figure 6. The average output power increased linearly as the enhancement of the pump power with a slope efficiency of  $\sim 18.8\%$ . When the pump power reached  $\sim 4.40$  W, the maximum output power of  $\sim 461$  mW and pulse energy of  $\sim 0.14$   $\mu$ J were obtained. The corresponding peak powers of the pulses were calculated and shown in Figure 6(b). We noted that the pulse peak powers were approximately maintained at a stable value of  $\sim 21$  W, when the pump power was in the range of  $\sim 2.57$  W to  $\sim 3.74$  W. Similar results were reported in Refs. [20, 28, 34], and the phenomenon was caused by the peak power clamping (PPC) effect. The decrease in the peak power at higher pump power was considered as the result of performance degrading of the SESA, since the high pulse energy will lead to a large heat accumulation at the SESA. The above characteristics of the



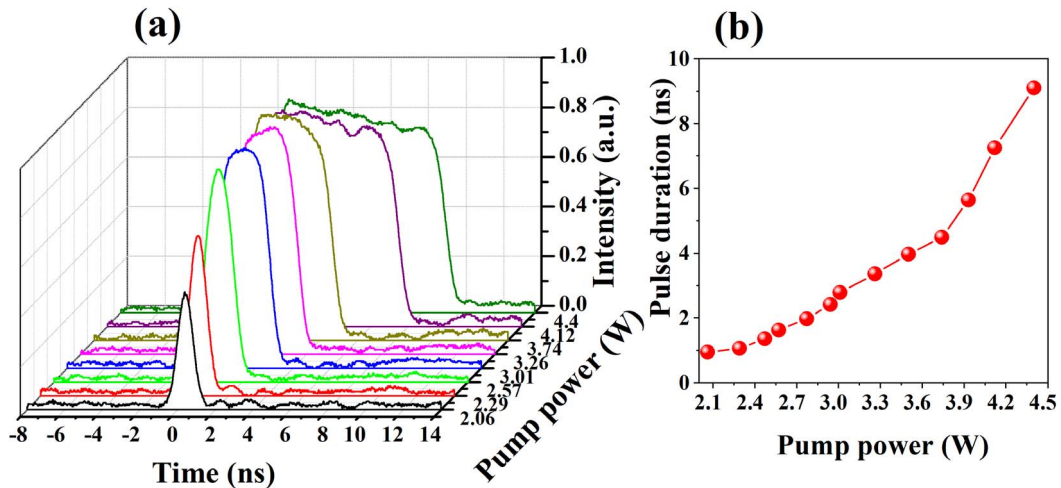


Figure 5. (a) Mode-locked pulse envelopes at different pump powers; (b) pulse duration versus pump power.

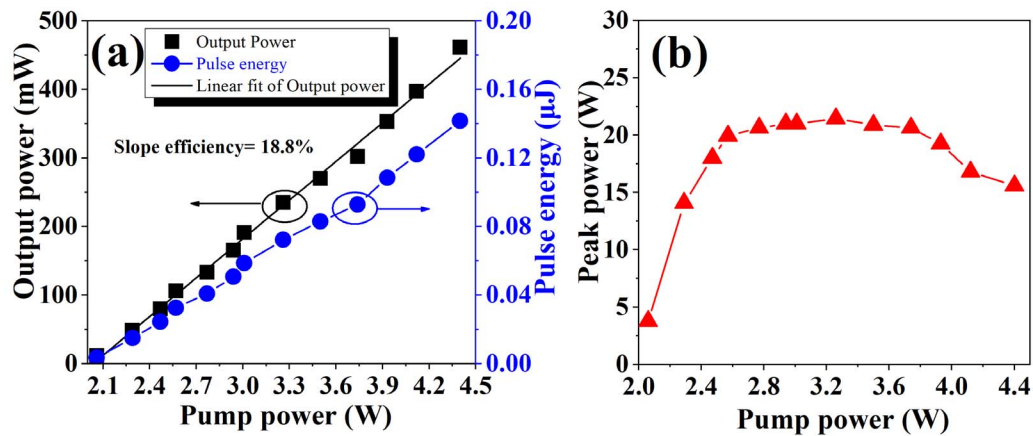


Figure 6. (a) Average output power and pulse energy variance with the pump power; (b) peak power variance with the pump power.

mode-locked pulses, such as pulse duration increase with the pump power, nanosecond rectangular-shaped pulse and PPC effect, indicated that the mode-locked fiber laser operated in the DSR regime<sup>[17–21, 28]</sup>.

## 2.2. CNTs-PVA film based nanosecond mode-locked fiber laser

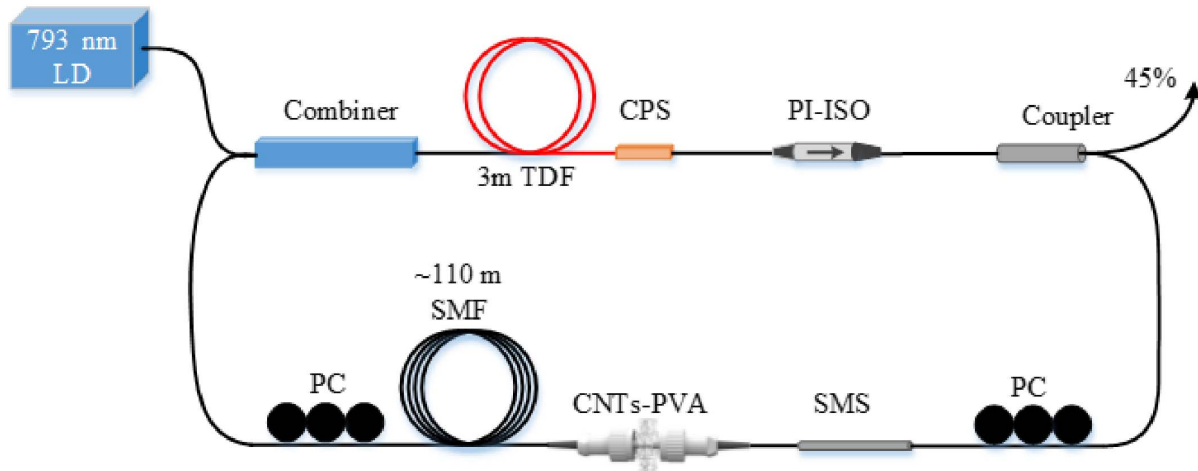
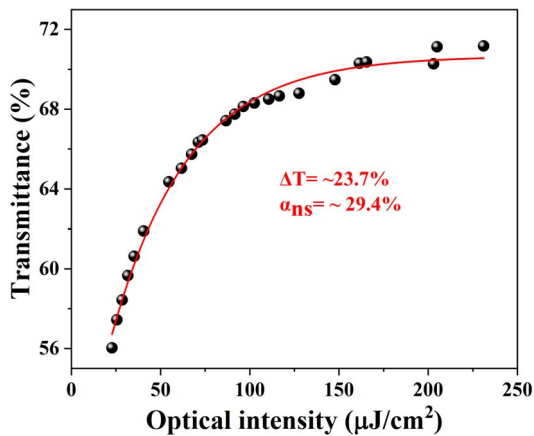
Similarly to the fiber laser configuration in Figure 2, the experimental setup of CNTs-based nanosecond mode-locked TDFL is shown in Figure 7. The laser output ratio is 45%, and a piece of  $\sim 110$  m long SMF was used for managing the dispersion and enhancing the nonlinearity. Thus, the total cavity length is  $\sim 120$  m and the dispersion is  $\sim -10.75$  ps<sup>2</sup>. Through a typical balanced twin-detector method, the saturable absorber property of the CNTs-PVA film was measured and shown in Figure 8. The ultra-pulse light source is a mode-locked fiber laser with central

wavelength, pulse duration and repetition rate of  $\sim 1958$  nm,  $\sim 725$  fs and  $\sim 55$  MHz, respectively. The CNTs-PVA film SA has a modulation depth of  $\sim 23.7\%$  and nonsaturable loss of  $\sim 29.4\%$ .

When the CNTs-PVA film was inserted into the cavity, with proper manipulation of the PCs, the initial mode-locked operation also could be obtained. Here, we characterized the mode-locking performance at the maximum pump power of  $\sim 2.12$  W, as shown in Figure 9. The mode-locked fiber laser centered at  $\sim 2003.1$  nm with 3 dB bandwidth of  $\sim 1.3$  nm, which depends on the transmission spectrum of the SMS fiber structure MMI, is shown in Figure 9(a). Figure 9(b) depicts the RF spectra. The repetition rate of the mode-locked pulse is  $\sim 1.718$  MHz and the signal-to-noise ratio is larger than 63 dB, which indicate a stable fundamental mode-locked operation. Figures 9(c) and 9(d) depict the mode-locked pulses. The single pulse duration of  $\sim 9.3$  ns with an approximate rectangular shape was measured. Together with the spectral width, the time-band product is

**Table 1.** Comparison of output properties of 2  $\mu\text{m}$  nanosecond mode-locked fiber lasers

	SA	Wavelength (nm)	Maximum output power (mW)	Maximum pulse energy (nJ)	Pulse duration (ns)	Reference
'Real' SA	Graphene	1897.7–1930.2	$\sim 33.9$	35.2	$\sim 122$ –143	[22]
	Graphene	2010.15	$\sim 5.5$	1.4	$\sim 3.8$ –94.3	[25]
	CNTs	2003.1	$\sim 46$	26.8	$\sim 9.3$	Our work
	SESA	2002.9	$\sim 461$	141.7	$\sim 1.6$ –9.1	Our work
'Artificial' SA	NPR	1960	$\sim 95$	/	$\sim 304$	[23]
	Self-mode-locking	1985.5	$\sim 66.8$	$\sim 32.7$	$\sim 40$	[24]
	NALM	1975.56	$\sim 43.1$	$\sim 40.5$	$\sim 3.74$ –72.19	[28]
	NALM	1940.2–1969.2	$\sim 60.73$	$\sim 19.51$	$\sim 0.48$ –6.19	[27]
	NALM	1985	$\sim 670$	$\sim 400$	$\sim 50$	[29]
	NOLM	1948.13	$\sim 295$	$\sim 1.5 \mu\text{J}$	$\sim 2.4$ –21.2	[30]
	NOLM	2005.9	$\sim 1.4 \text{ W}$	$\sim 353$	$\sim 1.9$ –13.7	[31]
	Intermode beating	1983	$\sim 1.03 \text{ W}$	$\sim 107$	45	[35]

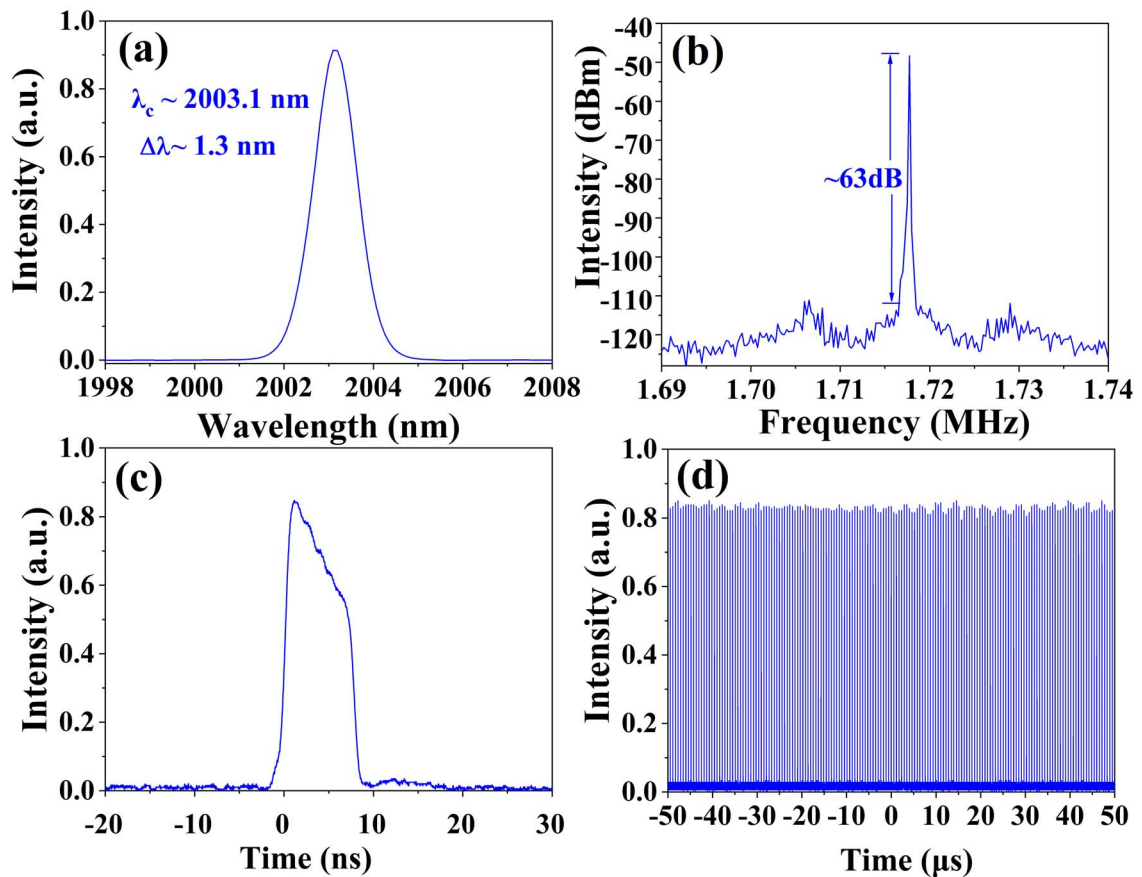
**Figure 7.** Schematic configuration of CNTs-PVA-based nanosecond TDFL.**Figure 8.** Saturable absorber properties of the CNTs-PVA.

$\sim 904$ . At a pump power of  $\sim 2.12 \text{ W}$ , the average output power of the fiber laser is  $\sim 46 \text{ mW}$ , corresponding to the pulse energy of  $\sim 26.8 \text{ nJ}$ .

In addition, Table 1 summarizes the representative output performances of all-fiberized nanosecond mode-locked TDFLs. In our experiment, the maximum average output power is  $\sim 461 \text{ mW}$  corresponding to the pulse energy of  $\sim 141.7 \text{ nJ}$ , which is higher than the previously reported results of 'real' SA-based nanosecond passively mode-locked TDFLs and comparable to the results of 'artificial' SA-based ones. However, owing to the relatively low damage threshold of the 'real' SA, the 'real' SA-based nanosecond mode-locked TDFLs still have the lower output power and pulse energy.

### 3. Conclusion

In summary, we have experimentally presented long cavity nanosecond-scale mode-locked fiber lasers with wavelength above  $2 \mu\text{m}$  based on real SAs (SESA and CNTs) and SMS fiber structure MMI. The DSR pulse with output pulse



**Figure 9.** Typical mode-locking performance of the CNT-based fiber laser at pump power of  $\sim 2.12$  W: (a) output spectrum; (b) output RF spectrum; (c), (d) mode-locked pulses.

energy up to  $\sim 0.14$   $\mu\text{J}$  was reported for the first time in real SA-based TDFL. This nanosecond mode-locked pulse with high pulse energy can be widely used in industrial processing and as a pump source for mid-infrared ( $\sim 3.5$   $\mu\text{m}$ ) pulse laser generation after further amplification. Moreover, our experimental results verified the feasibility of SMS fiber structure MMI functioning as an effective wavelength filter in mode-locked TDFL.

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