

TRINITY RELATIONS IN THE UNIVERSE

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On large scales where linear perturbation theory is valid, the mean square values of the mass excess $(\frac{\delta M}{M})^2$, the peculiar velocity v and the microwave background anisotropy $(\frac{\delta T}{T})_{sw}^2$ due to the Sachs-Wolfe effect, are simply expressed in terms of the present-day power spectrum of the density irregularity $P(k)$:

$$\left(\frac{\delta M}{M}\right)^2(r_M) = \frac{1}{2\pi^2} \int_0^\infty k^2 P(k) \exp(-k^2 r_M^2) dk, \quad (1)$$

$$v^2(r_v) = \frac{\Omega^{1.2} H^2}{2\pi^2} \int_0^\infty P(k) \exp(-k^2 r_v^2) dk, \quad (2)$$

$$\left(\frac{\delta T}{T}\right)_{sw}^2(r_T) = \frac{\Omega^{0.6} H^4}{8\pi^2} \int_0^\infty k^{-2} P(k) \exp(-k^2 r_T^2) dk. \quad (3)$$

For simplicity, we have used a Gaussian window function. If we consider the length scales satisfying that $r_M > r_v > r_T$, one can prove the inequalities among eqs.(1) to (3), *independently* of the specific form of $P(k)$; in the limit of $r_M \gg r_v \gg r_T$, they are reduced to:

$$\left(\frac{v}{1300 \text{ km/sec}}\right) \gtrsim \Omega^{0.6} \left(\frac{r_M}{8h^{-1} \text{ Mpc}}\right) \left(\frac{\delta M}{M}\right), \quad (4)$$

$$\left(\frac{\delta T/T}{5 \times 10^{-6}}\right) \gtrsim \Omega^{0.3} \left(\frac{r_M}{8h^{-1} \text{ Mpc}}\right)^2 \left(\frac{\delta M}{M}\right), \quad (5)$$

$$\left(\frac{\delta T/T}{6 \times 10^{-5}}\right) \gtrsim \Omega^{-0.3} \left(\frac{r_v}{60h^{-1} \text{ Mpc}}\right) \left(\frac{v}{1000 \text{ km/sec}}\right). \quad (6)$$

Interestingly, the ranges of the predicted values for $(\frac{\delta M}{M})$, v and $(\frac{\delta T}{T})$ are very close to the observed ones. With the observational data on $(\frac{\delta M}{M})$, v and $(\frac{\delta T}{T})$, the above *trinity* relations can be used to address the basic question; “ Does the gravitational instability picture account for the formation of large scale structures in the universe ? ” The detailed work is now in progress (Suto *et al.* 1987).