

## A SINGULAR INTEGRAL ON $L^2(\mathbb{R}^n)$

DASHAN FAN

**ABSTRACT.** We consider a convolution singular integral operator  $T_b^\Omega$  associated to a kernel  $K(x) = b(x)\Omega(x)|x|^{-n}$ , and prove that if  $b \in L^\infty(\mathbb{R}^n)$  is a radial function and  $\Omega \in H(\Sigma_{n-1})$  with mean zero condition (1), then  $T_b^\Omega$  is a bounded linear operator in the space  $L^2(\mathbb{R}^n)$ .

Let  $\mathbb{R}^n$  be the  $n$ -dimensional real Euclidean space. The unit ball  $B_n$  is the set  $\{z \in \mathbb{R}^n, |z| \leq 1\}$ , and the unit sphere  $\Sigma_{n-1}$  is the boundary of  $B_n$ . Let  $d\sigma(x')$  be the element of Lebesgue measure on  $\Sigma_{n-1}$  so that the measure of  $\Sigma_{n-1}$  is 1. Let  $L^p(\mathbb{R}^n)$  and  $L^p(\Sigma_{n-1})$  be the spaces of Lebesgue  $L^p$ -integrable functions on  $\mathbb{R}^n$  and  $\Sigma_{n-1}$ , respectively. Besides considering these  $L^p$  spaces, we are also interested in the Hardy spaces on both  $\mathbb{R}^n$  and  $\Sigma_{n-1}$ .

The Poisson kernel  $P_t(x)$  on  $\mathbb{R}^n$  is defined by

$$P_t(x) = C_n t(t^2 + |x|^2)^{-(n+1)/2}, \quad C_n = \Gamma((n+1)/2)\pi^{-(n+1)/2}.$$

For any  $f \in s'(\mathbb{R}^n)$ , we define the radial maximal function  $P^*f$  by

$$P^*f(x) = \sup_{t>0} |P_t * f(x)|,$$

where  $s'(\mathbb{R}^n)$  is the space of Schwartz distributions on  $\mathbb{R}^n$ .

The Hardy space  $H(\mathbb{R}^n)$  is the linear space of distributions  $f$  with the finite norm  $\|f\|_{H(\mathbb{R}^n)} = \|P^*f\|_{L^1(\mathbb{R}^n)} < \infty$ . More details about the Hardy space on  $\mathbb{R}^n$  can be found in [8].

The Poisson kernel on  $\Sigma_{n-1}$  is defined by

$$P_{ry'}(x') = (1 - r^2)/|ry' - x'|^n,$$

where  $0 \leq r < 1$  and  $x', y' \in \Sigma_{n-1}$ . For any  $f \in s'(\Sigma_{n-1})$ , we define the radial maximal function  $P^+f(x')$  by

$$P^+f(x') = \sup_{0 \leq r < 1} \left| \int_{\Sigma_{n-1}} f(y') P_{rx'}(y') d\sigma(y') \right|,$$

where  $s'(\Sigma_{n-1})$  is the space of Schwartz distributions on  $\Sigma_{n-1}$ . The Hardy space  $H(\Sigma_{n-1})$  is the linear space of distributions  $f \in s'(\Sigma_{n-1})$  with the finite norm  $\|f\|_{H(\Sigma_{n-1})} = \|P^+f\|_{L^1(\Sigma_{n-1})} < \infty$ . Various properties of Hardy space on  $\Sigma_{n-1}$  were studied in [5]. In particular, a well-known result is  $L^1(\Sigma_{n-1}) \supseteq H(\Sigma_{n-1}) \supseteq L^q(\Sigma_{n-1})$  for any  $q > 1$ .

Received by the editors October 5, 1992.

AMS subject classification: 42B99.

© Canadian Mathematical Society 1994.

Suppose  $\Omega$  is a homogeneous function of degree zero and satisfies

$$(1) \quad \int_{\Sigma_{n-1}} \Omega(x') d\sigma(x') = 0.$$

Let  $b$  be a bounded and radial function. We define a kernel  $K$  by

$$(2) \quad K(x) = b(x)\Omega(x)|x|^{-n}$$

and consider the singular integral  $(T_b^\Omega f)(x) = p. v. (K * f)(x)$ . This operator was first studied by Calderón and Zygmund in their pioneer papers [1] and [2] for the case  $b(x) \equiv 1$ . In [2], Calderón and Zygmund proved that if  $b(x) = 1$  and  $\Omega$  satisfies (1), then this operator  $T_1^\Omega$  is  $L^p(\mathbb{R}^n)$  ( $1 < p < \infty$ ) bounded provided  $\Omega \in L^q(\Sigma_{n-1})$  for some  $q > 1$ . Coifman and Weiss [4] improved Calderón and Zygmund's result under a weaker condition  $\Omega \in H(\Sigma_{n-1})$ .

In [6], R. Fefferman generalized this singular operator by considering any  $L^\infty$  function  $b$ . In this Ph.D. thesis, (see [3]) K. Chen proved that  $T_b^\Omega$  is a linear bounded operator in the space  $L^p(\mathbb{R}^n)$  ( $n > 2, 1 < p < \infty$ ) if  $\Omega \in L^q(\Sigma_{n-1})$  for some  $q > 1$  and satisfies the condition (1). The first step of Chen's proof is to prove the  $L^2$  boundedness of  $T_b^\Omega$ ; then the  $L^p$  result follows by an interpolation argument. Recently, Jian and Lu (see [7], p. 140) improved Chen's  $L^2$  result by using a weaker assumption on  $\Omega$ . They proved that if  $\Omega$  satisfies the mean zero condition (1) and  $\Omega$  is in the block space  $B_q(\Sigma_{n-1})$ , then  $T_b^\Omega$  is a bounded operator in  $L^2(\mathbb{R}^n)$  ( $n > 2$ ).

Comparing the early result [4] of Coifman and Weiss, we find a more natural condition on  $\Omega$  should be  $\Omega \in H(\Sigma_{n-1})$ . The following theorem then is the main purpose of this short note.

**THEOREM.** *Suppose that  $\Omega$  is a homogeneous function of degree zero, and satisfies (1). If  $b$  is a bounded radial function and  $\Omega \in H(\Sigma_{n-1})$ ,  $n \geq 2$ , then the operator  $T_b^\Omega$  is bounded in  $L^2(\mathbb{R}^n)$  and its operator norm is bounded by  $C\|b\|_\infty\|\Omega\|_{H(\Sigma_{n-1})}$ , where  $C$  is a constant independent of function  $b(x)$  and  $\Omega(x)$ .*

**PROOF.** By the Plancherel theorem, we need only to prove that

$$(3) \quad \|\widehat{K}\|_\infty \leq C\|b\|_\infty\|\Omega\|_H.$$

In fact,

$$\begin{aligned} |\widehat{K}(x)| &\leq C\|b\|_\infty \int_0^1 \left| \int_{\Sigma_{n-1}} \Omega(\xi') (e^{it\langle x', \xi' \rangle} - 1) d\sigma(\xi') \right| t^{-1} dt \\ &\quad + C\|b\|_\infty \int_1^\infty \left| \int_{\Sigma_{n-1}} \Omega(\xi') e^{it\langle x', \xi' \rangle} d\sigma(\xi') \right| t^{-1} dt. \end{aligned}$$

It is easy to see that the first term above is bounded by  $\|b\|_\infty\|\Omega\|_{L^1} \leq C\|b\|_\infty\|\Omega\|_H$ . Thus now we only have to prove that

$$(4) \quad \int_1^\infty \left| \int_{\Sigma_{n-1}} \Omega(\xi') e^{it\langle x', \xi' \rangle} d\sigma(\xi') \right| t^{-1} dt \leq C\|\Omega\|_{H(\Sigma_{n-1})}.$$

In order to prove (4), we need introduce the atomic decomposition of  $H(\Sigma_{n-1})$ . An exceptional atom is an  $L^\infty$  function  $a(x)$  satisfying  $\|a\|_\infty \leq 1$ . A  $(1, \infty)$  atom is an  $L^\infty$  function  $a(x)$  which satisfies

- (i)  $\text{supp}(a) \subset \{x' \in \Sigma_{n-1}, |x' - x'_0| < \rho \text{ for some } x'_0 \in \Sigma_{n-1} \text{ and } \rho > 0\}$ ,
- (ii)  $\int_{\Sigma_{n-1}} a(\xi') d\sigma(\xi') = 0$ ,
- (iii)  $\|a\|_\infty \leq \rho^{-n+1}$ .

By [5], we know that any  $\Omega \in H(\Sigma_{n-1})$  has an atomic decomposition  $\Omega(\xi') = \sum \lambda_j a_j(\xi')$ , where the  $a_j$ 's are either exceptional atoms or  $(1, \infty)$  atoms and  $\sum |\lambda_j| \leq C \|\Omega\|_{H(\Sigma_{n-1})}$ . Therefore it remains to prove that for all atoms  $a(\xi')$ ,

$$(5) \quad L_a(x') = \int_1^\infty \left| \int_{\Sigma_{n-1}} a(\xi') e^{it\langle x', \xi' \rangle} d\sigma(\xi') \right| t^{-1} dt \leq C$$

with a constant  $C$  independent of  $a(\xi')$  and  $x' \in \Sigma_{n-1}$ . We will prove (5) in the two different cases  $n > 2$  and  $n = 2$ , respectively.

CASE  $n > 2$ . If  $a$  is exceptional, by [3], we have  $\|L_a\|_\infty \leq C\|a\|_2 \leq C$ . Suppose  $a$  is a  $(1, \infty)$  atom; without loss of generality, we may assume that  $\text{supp}(a)$  is contained in the ball  $B(\mathbf{1}, \rho)$ , where  $\mathbf{1} = (1, 0, \dots, 0)$ . By a rotation we also can assume  $x' = \mathbf{1}$ . Under these assumptions, we let

$$\xi' = (s, \xi_1, \xi_2, \dots, \xi_n).$$

Then

$$L_a(x') \leq C \int_0^\infty t^{-1} |\widehat{F}(t)| dt,$$

where  $F(s) = (1-s^2)^{(n-3)/2} \chi_{(-1,1)}(s) \int_{\Sigma_{n-2}} a(s, (1-s^2)^{1/2} y') d\sigma(y')$  and  $\widehat{F}(t)$  is the Fourier transform of  $F(s)$ . Now we easily see that  $\int_{\mathbb{R}} F(s) ds = 0$  and  $\text{supp}(F) \subseteq (1-\rho, 1)$ . Furthermore, we have  $\|F\|_\infty \leq \|a\|_\infty \int_{\Sigma_{n-2} \cap B(\mathbf{1}, \rho)} d\sigma(y') \leq C\rho^{-1}$ . These imply that, up to a constant independent of atom  $a(\xi')$ ,  $F$  is a  $(1, \infty)$  atom on  $\mathbb{R}$ . Thus using the Hardy inequality (see [8]), we obtain that  $L_a(x') \leq C$ . Thus the case  $n > 2$  of (5) is now proved.

CASE  $n = 2$ . In this case  $\Sigma_1 = \mathbb{T}$ , the one-dimension torus. We will first prove (5) for any  $(1, \infty)$  atom  $a(\theta)$ . As before, we may assume that  $\text{supp}(a) \subseteq (-\rho, \rho)$ . Let  $x' = e^{i\alpha}$ ; then

$$L_a(x') = \left\{ \int_{-\rho}^\infty + \int_1^{\rho^{-2}} \right\} t^{-1} \left| \int_{-\pi}^\pi a(\theta) e^{it \cos(\theta-\alpha)} d\theta \right| dt \leq J_1 + J_2.$$

We will only estimate  $J_1$  and  $J_2$  for the case  $\cos \alpha \neq 0$ ; the estimate of the case  $\cos \alpha = 0$  is easier than the prior case.

$$\begin{aligned} J_2 &= \int_1^{\rho^{-2}} t^{-1} \left| \int_{-\pi}^\pi a(\theta) e^{it \cos \theta \cos \alpha} e^{it \sin \theta \sin \alpha} d\theta \right| dt \\ &= \int_{\cos \alpha}^{\rho^{-2} \cos \alpha} t^{-1} \left| \int_{-\pi}^\pi a(\theta) e^{it(\cos \theta - 1)} e^{it \tan \alpha \sin \theta} d\theta \right| dt \\ &\leq \int_{\cos \alpha}^{\rho^{-2} \cos \alpha} t^{-1} \left| \int_{-\pi}^\pi a(\theta) e^{it \tan \alpha \sin \theta} d\theta \right| dt \\ &\quad + \int_0^{\rho^{-2}} \int_{-\pi}^\pi |a(\theta)(\cos \theta - 1)| d\theta dt = I + II. \end{aligned}$$

It is easy to see that  $I \leq C$ .

$$\begin{aligned} I &\leq \int_{\cos \alpha}^{\rho^{-2} \cos \alpha} t^{-1} \left| \int_{-\pi}^{\pi} a(\theta) e^{it \tan \alpha (\sin \theta - \theta)} e^{it \tan \alpha \theta} d\theta \right| dt \\ &\leq \int_0^{\infty} t^{-1} \left| \int_{-\pi}^{\pi} a(\theta) e^{it\theta} d\theta \right| dt \\ &\quad + \int_{\cos \alpha}^{\rho^{-2} \cos \alpha} |\tan \alpha| \int_{-\pi}^{\pi} |a(\theta)(\sin \theta - \theta)| d\theta dt \\ &\leq \left( 1 + \int_0^{\infty} |\widehat{a}(t)| t^{-1} dt \right) \leq C \quad (\text{by the Hardy inequality}). \end{aligned}$$

Next we estimate

$$J_1 = \int_{\rho^{-2}}^{\infty} t^{-1} \left| \int_{-\pi}^{\pi} a(\theta + \alpha) e^{it \cos \theta} d\theta \right| dt.$$

By Hölder’s inequality, we have

$$\begin{aligned} J_1 &\leq 2\rho^{2/q} \left( \int_{\rho^{-2}}^{\infty} \left| \int_0^{\pi} + \int_{-\pi}^0 a(\theta + \alpha) e^{it \cos \theta} d\theta \right|^q dt \right)^{1/q} \\ &\leq 2\rho^{2/q} (J_{1,1} + J_{1,2}), \quad \text{where } 1 < p = q/(q - 1) \leq 3/2. \end{aligned}$$

We will only estimate the above term

$$J_{1,1} = \left( \int_{\rho^{-2}}^{\infty} \left| \int_0^{\pi} a(\theta + \alpha) e^{it \cos \theta} d\theta \right|^q dt \right)^{1/q};$$

the estimate of  $J_{1,2}$  is exactly same. After change variable  $u = \cos \theta$ , we know that

$$J_{1,1} = \left( \int_{\rho^{-2}}^{\infty} \left| \int_{\cos \rho}^1 a(\alpha + \cos^{-1} u) (1 - u^2)^{-1/2} e^{itu} du \right|^q dt \right)^{1/q}.$$

Thus by the Hausdorff-Young inequality, we have

$$\begin{aligned} J_{1,1} &\leq \left( \int_{\mathbb{R}} \chi_{(\cos \rho, 1)}(t) |a(\alpha + \cos^{-1} t) (1 - t^2)^{-1/2}|^p dt \right)^{1/p} \\ &\leq \rho^{-1} \left( \int_{\cos \rho}^1 (1 - t^2)^{-p/2} dt \right)^{1/p} \leq \rho^{-1} \left( \int_0^{\rho} |\sin^{1-p}(\theta)| d\theta \right)^{1/p} \\ &\leq \rho^{-2} \rho^{2/p} = \rho^{-2/q}. \end{aligned}$$

This shows that  $J_1 \leq C$ . Therefore we complete the proof of (5) for any  $(1, \infty)$  atom. Finally we need to prove (5) for an exceptional atom  $a(\theta)$ . But this case easily follows by mimicking the estimate of  $J_1$  in the above argument. Now the theorem is proved.

REFERENCES

1. A. P. Calderón and A. Zygmund, *On the existence of certain singular integrals*, Acta. Math. **88**(1952), 85–139.
2. ———, *On singular integrals*, Amer. J. Math. **18**(1956), 289–309.
3. K. Chen, *On a singular integral*, Studia Math. **LXXXV**(1987), 61–72.
4. R. Coifman and G. Weiss, *Extension of Hardy spaces and their use in analysis*, Bull. Amer. Math. Soc. **83**(1977), 569–645.
5. L. Colzani, *Hardy spaces on Sphere*, Ph.D. thesis, Washington University, St. Louis, 1982.

6. R. Fefferman, *A note on singular integrals*, Proc. Amer. Math. Soc. **74**(1979), 266–270.
7. S. Lu, M. Taibleson and G. Weiss, *Space Generated by Blocks*, Beijing Normal University Mathematics Series, Beijing Normal University Press, 1989.
8. A. Torchinsky, *Real-variable Methods in Harmonic Analysis*, Academic Press, 1986.

*Department of Mathematical Sciences*  
*University of Wisconsin-Milwaukee*  
*Milwaukee, Wisconsin 53201*  
*U.S.A.*  
*e-mail: Fan@csd4.csd.uwm.edu*