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ABELIAN CYCLES IN THE HOMOLOGY OF THE TORELLI GROUP

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Abstract In the early 1980s, Johnson defined a homomorphism $\mathcal{I}_{g}^{1} \to \bigwedge^{3} H_{1}(S_{g},\mathbb{Z})$, where \mathcal{I}_{g}^{1} is the Torelli group of a closed, connected, and oriented surface of genus g with a boundary component and S_{g} is the corresponding surface without a boundary component. This is known as the Johnson homomorphism.

We study the map induced by the Johnson homomorphism on rational homology groups and apply it to abelian cycles determined by disjoint bounding-pair maps, in order to compute a large quotient of $H_n\left(\mathcal{I}_g^1,\mathbb{Q}\right)$ in the stable range. This also implies an analogous result for the stable rational homology of the Torelli group $\mathcal{I}_{g,1}$ of a surface with a marked point instead of a boundary component. Further, we investigate how much of the image of this map is generated by images of such cycles and use this to prove that in the pointed case, they generate a proper subrepresentation of $H_n\left(\mathcal{I}_{g,1}\right)$ for $n \geq 2$ and glarge enough.

Keywords: Torelli group; Johnson homomorphism

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1. Introduction

1.1. The Torelli group

Let $S_{g,r}^s$ denote a connected and oriented surface of genus g, with r marked points and s boundary components. If either of r or s is zero, we simply omit the corresponding index.

Recall that the mapping-class group of $S_{q,r}^s$ is the group

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$$\Gamma_{g,r}^s := \pi_0 \operatorname{Diff}^+ \left(S_{g,r}^s \right)$$

of isotopy classes of orientation-preserving diffeomorphisms of $S_{g,r}^s$ that fix the marked points and boundary components. We define $H := H_1(S_g, \mathbb{Q})$ and $H_{\mathbb{Z}} := H_1(S_g, \mathbb{Z})$. Since $\text{Diff}^+(S_{g,r}^s)$ acts on $S_g, \Gamma_{g,r}^s$ acts on H, which is a symplectic vector space. The $\Gamma_{g,r}^s$ -action preserves the corresponding symplectic form

$$\omega := \sum_{i=1}^{g} a_i \wedge b_i \in \bigwedge^2 H,$$

where $a_1, b_1, \ldots, c, a_g, b_g$ is any symplectic basis of H, and this gives us a group homomorphism

$$\Gamma^s_{q,r} \to \operatorname{Sp}(H),$$

where $\operatorname{Sp}(H)$ denotes the symplectic group of H. We define the *Torelli group* of $S^s_{g,r}$, denoted $\mathcal{I}^s_{g,r}$, to be the kernel of this homomorphism. The Torelli group is thus the subgroup of $\Gamma^s_{g,r}$ consisting of those isotopy classes of diffeomorphisms that act trivially on H. It is a classic result that the image of $\Gamma^s_{g,r} \to \operatorname{Sp}(H)$ is precisely the arithmetic subgroup $\operatorname{Sp}(H_{\mathbb{Z}})$, so we get a short exact sequence

$$1 \to \mathcal{I}_{g,r}^s \to \Gamma_{g,r}^s \to \operatorname{Sp}(H_{\mathbb{Z}}) \to 1.$$

The conjugation action of $\Gamma_{g,r}^s$ on $\mathcal{I}_{g,r}^s$ gives us an action by $\operatorname{Sp}(H_{\mathbb{Z}})$ on $\mathcal{I}_{g,r}^s$ by outer automorphisms, which makes $H_n(\mathcal{I}_{g,r}^s, \mathbb{Q})$ into an $\operatorname{Sp}(H_{\mathbb{Z}})$ -representation, for all $n \geq 0$. Unless explicitly stated, all homology is taken with rational coefficients from now on.

1.2. The Johnson homomorphism

Many basic questions about the groups $H_n(\mathcal{I}_{g,r}^s)$ are unanswered for $n \geq 2$. By contrast, the case n = 1 is completely understood, due to work by Johnson in the early 1980s [7]. Among other things, Johnson contructed a group homomorphism

$$\tau_J: \mathcal{I}_g^1 \to \bigwedge^3 H_{\mathbb{Z}},$$

which is known as the Johnson homomorphism. For $g \ge 3$, he proved that the induced $\operatorname{Sp}(H_{\mathbb{Z}})$ -equivariant map

$$\tau: H_1\left(\mathcal{I}_g^1\right) \to \bigwedge^3 H$$

is an isomorphism. In particular, this proves that $H_1(\mathcal{I}_g^1)$ is finite-dimensional and algebraic as a representation of $\operatorname{Sp}(H_{\mathbb{Z}})$ for $g \geq 3$. Two important open questions about the higher homology groups are thus as follows:

Question 1. Are the homology groups $H_n(\mathcal{I}_{g,r}^s)$ finite-dimensional, for $n \ge 1$ and g sufficiently large?

Question 2. Are the homology groups $H_n(\mathcal{I}_{g,r}^s)$ algebraic $\operatorname{Sp}(H_{\mathbb{Z}})$ -representations, for $n \geq 1$ and g sufficiently large?

In low genera, we know that the answer to both questions is generally no. In genus 2, it was proven by Mess [12] that for a certain permutation matrix $e \in \text{Sp}(2,\mathbb{Z})$, we have

$$H_1(\mathcal{I}_2) \cong \mathbb{Q}[\operatorname{Sp}(2,\mathbb{Z})/(\langle e \rangle \ltimes (\operatorname{SL}(2,\mathbb{Z}) \times \operatorname{SL}(2,\mathbb{Z})))].$$

He also showed that $H_3(\mathcal{I}_3)$ contains a subrepresentation isomorphic to $\mathbb{Q}[\operatorname{Sp}(3,\mathbb{Z})/(\operatorname{Sp}(1,\mathbb{Z}) \times \operatorname{Sp}(2,\mathbb{Z}))]$. It was recently proven by Gaifullin [5] that for $g \geq 3$ and $2g-3 \leq n \leq 3g-6$, $H_n(\mathcal{I}_g,\mathbb{Z})$ contains a free abelian subgroup of infinite rank. This means that $H_n(\mathcal{I}_g)$ is neither algebraic nor finite-dimensional in either of these cases.

For $n \geq 1$, τ_J induces a map

$$(\tau_J)_*: H_n\left(\mathcal{I}_g^1\right) \to H_n\left(\bigwedge^3 H\right) \cong \bigwedge^n\left(\bigwedge^3 H\right),$$

which we will denote by ψ_n . An equivalent definition of ψ_n is as the composition of the comultiplication $H_n(\mathcal{I}_g^1) \to \bigwedge^n H_1(\mathcal{I}_g^1)$ with $\wedge^n \tau$. For $n \geq 2$, it follows from results by Hain that the map ψ_n is not surjective [6], but it is natural to ask the question:

Question 3. Is ψ_n injective for all $n \ge 1$ and g sufficiently large?

It is not clear whether to expect this to be true. A positive answer would of course imply positive answers to Questions 1 and 2 in the corresponding case. It was proven by Kupers and Randal-Williams that the converse is also true [11]. They also gave an explicit description of the kernel of the dual map of ψ_n , under these assumptions.

Remark 1.1. In the case of a surface without a boundary component, Johnson instead defined a homomorphism $\mathcal{I}_g \to \left(\bigwedge^3 H_{\mathbb{Z}}\right)/H_{\mathbb{Z}}$, where $H_{\mathbb{Z}}$ is embedded as a subspace of $\bigwedge^3 H_{\mathbb{Z}}$ by inserting the symplectic form. This is also known as the Johnson homomorphism. This homomorphism induces an $\operatorname{Sp}(H_{\mathbb{Z}})$ -equivariant isomorphism $H_1(\mathcal{I}_g) \to \left(\bigwedge^3 H\right)/H$.

Remark 1.2. Hain [6] has studied the dual of the map ψ_n in the unpointed case and determined its image completely for n = 2. For n = 3, its image was determined by Sakasai, up to one irreducible subrepresentation [14], the presence of which was settled by Kupers and Randal-Williams [11].

The Johnson homomorphisms are closely related to homomorphisms $\Gamma_{g,1} \to \frac{1}{2} \bigwedge^3 H_{\mathbb{Z}} \rtimes$ Sp $(H_{\mathbb{Z}})$ and $\Gamma_g \to \frac{1}{2} \left(\bigwedge^3 H_{\mathbb{Z}}\right) / H_{\mathbb{Z}} \rtimes \text{Sp}(H_{\mathbb{Z}})$ defined by Morita [13], which were used by Kawazumi and Morita to give a description of the tautological subalgebra of $H^*(\Gamma_{g,1})$ in terms of trivalent graphs [10].

We will study how the image of ψ_n decomposes into irreducible $\operatorname{Sp}(H_{\mathbb{Z}})$ -representations and determine a lower bound on this image in the stable range. The representation $\bigwedge^n \left(\bigwedge^3 H\right)$ is algebraic, and this implies that its irreducible subrepresentations viewed as an Sp(H)- and an Sp($H_{\mathbb{Z}}$)-representation agree. Therefore, we will work with the image of ψ_n as an Sp(H)-representation.

1.3. Irreducible representations of symplectic groups

In order to state our results, we need to recall some basics from the representation theory of symplectic groups. The isomorphism classes of irreducible representations of $\operatorname{Sp}(H)$ are indexed by *partitions*. A partition is a decreasing sequence $(\lambda_1 \geq \lambda_2 \geq \cdots \geq 0 \geq 0 \geq \cdots)$ of nonnegative integers that eventually reaches zero. For example, the standard representation H corresponds to the partition $(1 \geq 0 \geq 0 \geq \cdots)$. We will write $(\lambda_1 \geq \lambda_2 \geq \cdots \geq \lambda_k)$ for the partition, where all following entries in the sequence are zero. The sum $\sum_{n\geq 1} \lambda_n$ is called the *weight* of the partition. If $\lambda_{i-1} > \lambda_i = \cdots = \lambda_{i+l-1} > \lambda_{i+l}$, we will often write $(\lambda_1 \geq \cdots > \lambda_i^l > \cdots \geq \lambda_k)$, for brevity, and use the convention that if l = 0, then λ_i does not occur in the partition.

We will write $V_{\lambda} = V_{\lambda_1,...,\lambda_k}$ for the irreducible representation corresponding to the partition $\lambda = (\lambda_1 \ge \lambda_2 \ge \cdots \ge \lambda_k)$. In particular, we have $V_1 \cong H$. We define the weight of V_{λ} to be the weight of λ . For $k \ge 1$, the representation $H^{\otimes k}$ contains irreducible subrepresentations of weight at most k, which means that the top weight irreducible subrepresentations of $\bigwedge^n (\bigwedge^3 H)$ have weight 3n. For more details on the irreducible representations of symplectic groups, see, for example, [4, Chapter 17].

1.4. Results

Our first main result is the following:

Theorem 1.3. For $n \ge 1$ and $g \ge 3n$, the image of $\psi_n : H_n(\mathcal{I}_g^1) \to \bigwedge^n(\bigwedge^3 H)$ contains the Sp(H)-subrepresentation of $\bigwedge^n(\bigwedge^3 H)$ spanned by all irreducible subrepresentations of weight 3n.

Remark 1.4. The kernel of the dual map described by Kupers and Randal-Williams contains no irreducible subrepresentations of top weight, which means that if Questions 1 and 2 have positive answers, Theorem 1.3 follows immediately.

To get a feeling for what Theorem 1.3 means in practice, let us list the irreducible subrepresentations of $\bigwedge^n (\bigwedge^3 H)$ of top weight for $n \leq 4$ in a table, together with their dimension for g = 3n. These can be computed using, for example, SAGE.

\overline{n} Irreducible subrepresentations of top weight in $\bigwedge^n (\bigwedge^3 H)$,	for $g \ge 3n$ Dimension for $g = 3n$
1 V ₁ ³	14
$2 \ V_{2^2,1^2} \oplus V_{1^6}$	19,383
$3 \hspace{0.1 cm} V_{3,2^3} \oplus V_{3^2,1^3} \oplus V_{2^3,1^3} \oplus V_{2^2,1^5} \oplus V_{1^9}$	$\approx 7.5 \cdot 10^7$
$4 V_{4,3,2^2,1} \oplus V_{4^2,1^4} \oplus V_{3^4} \oplus V_{3^2,2^2,1^2} \oplus V_{3^2,1^6} \oplus V_{3,2^3,1^3}$	$\approx 5.3 \cdot 10^{11}$
$\oplus V_{2^6} \oplus V_{2^4,1^4} \oplus V_{2^3,1^3} \oplus V_{2^2,1^8} \oplus V_{1^{12}}$	

Contrary to what the table might suggest, the irreducible subrepresentations are generally not of multiplicity 1 for higher n.

Theorem 1.3 is actually a special case of our second main theorem. This second theorem requires a bit more work to prove, which is why we list and prove the theorems separately. In order to state this result a bit more clearly, we will consider the irreducible Sp(H)-representation V_{1^k} as a \mathbb{Z} -graded vector space concentrated in degree k, and use the Koszul sign convention. Note that in particular, H is concentrated in degree 1. From now on, we will therefore use Λ to denote the free graded commutative algebra functor, rather than the exterior algebra functor of ungraded vector spaces.

Theorem 1.5. Set $n \ge 1$ and let $\lambda = (\lambda_1^{k_1} > \cdots > \lambda_m^{k_m} > 0)$ and $\mu = (\mu_1^{l_1} > \cdots > \mu_{m+2}^{l_{m+2}} > 0)$ be partitions such that $\mu_i = \lambda_i + 2$ for $i = 1, \ldots, m$ and $|\lambda| + |\mu| = n$. Let $k_{m+1} = k_{m+2} = 0$, $k = \sum_{i=1}^{m} k_i$, and $l = \sum_{i=1}^{m+2} l_i$, and suppose that $g \ge n + 2k + l$. Then the subrepresentation of



spanned by all irreducible subrepresentations of weight n + 2k is contained in $Im(\psi_n)$.

In particular, we see that the case m = 1, $\lambda_1 = 1$, $k_1 = n$ gives us back Theorem 1.3, since $\bigwedge^3 H \cong V_{1^3} \oplus V_1$.

Remark 1.6. The Johnson homomorphism $\tau_J : \mathcal{I}_g^1 \to \bigwedge^3 H_{\mathbb{Z}}$ actually factors as $\mathcal{I}_g^1 \to \mathcal{I}_{g,1} \to \bigwedge^3 H_{\mathbb{Z}}$, where the first map is induced by sewing a disk with a marked point into the boundary component and extending diffeomorphisms by the identity. Thus the analogous results to Theorems 1.3 and 1.5 also hold in the pointed case.

Our strategy is to consider *abelian cycles* in $H_n(\mathcal{I}_g^1)$ determined by *n*-tuples of pairwise disjoint *bounding-pair maps*. These notions will be defined in §3. What makes this strategy work is that the coproduct of an abelian cycle has a simple expression in terms of abelian cycles of lower degree (Lemma 3.2), which makes the map ψ_n easy to evaluate.

Remark 1.7. The main inspiration for using abelian cycles is papers by Church and Farb [1] and Sakasai [14]. In the former, these are used to study the images of certain homomorphisms $\tau_n : H_n(\mathcal{I}_{g,1}) \to \bigwedge^{n+2} H$ that generalize τ to higher degree (see [7, pp. 172–173]). Church and Farb prove, among other things, that for $n \ge 1$ and $g \ge n+2$, the homomorphism τ_n gives a surjection $H_n(\mathcal{I}_{g,1}) \twoheadrightarrow V_{1^{n+2}} \oplus V_{1^n}$. This means that the dimension of $H_n(\mathcal{I}_{g,1})$ has a lower bound of order $\sim g^{n+2}$ for all $n \ge 1$. We can note that by Theorem 1.3 and Remark 1.6 we get a lower bound on the dimensions of $H_n(\mathcal{I}_g^1)$ and $H_n(\mathcal{I}_{g,1})$ of order $\sim g^{3n}$ for $g \ge 3n$.

Remark 1.8. Church, Ellenberg, and Farb have used representation stability [2, §7.2] to prove the following theorem:

Theorem 1.9. (Church, Ellenberg, Farb). For each $n \ge 0$ there exists a polynomial $P_n(T)$ of degree at most 3n such that

$$\dim\left(Im\left(\bigwedge^{n} H^{1}\left(\mathcal{I}_{g}^{1}\right) \to H^{n}\left(\mathcal{I}_{g}^{1}\right)\right)\right) = P_{n}(g)$$

for $g \gg n$.

Since $H^n\left(\mathcal{I}_g^1\right) \cong H_n\left(\mathcal{I}_g^1\right)^*$ as $\operatorname{Sp}(H)$ -representations for all $n \ge 0$, it follows that the dimension of the image in Theorem 1.9 is bounded from below by the dimension of the image of $\psi_n : H_n\left(\mathcal{I}_g^1\right) \to \bigwedge^n\left(\bigwedge^3 H\right)$, so the following corollary follows from Theorems 1.3 and 1.9:

Corollary 1.10. The polynomial $P_n(T)$ in Theorem 1.9 has degree 3n.

For the final result of the paper, we will work with a marked point instead of a boundary component, since this gives us something a bit more general. Abusing notation, we will use ψ_n to denote the map $H_n(\mathcal{I}_{g,1}) \to \bigwedge^n \left(\bigwedge^3 H\right)$ as well. Which ψ_n is intended will be clear from context. The last main result concerns the limitations of using abelian cycles of the type we consider. More specifically:

Theorem 1.11. Set $n \geq 2$ and let $A_n(\mathcal{I}_{g,1}) \subset H_n(\mathcal{I}_{g,1})$ be the subrepresentation generated by abelian cycles determined by pairwise disjoint bounding-pair maps. Then $\psi_n(A_n(\mathcal{I}_{g,1}))$ is concentrated in weights $n, n+2, \ldots, 3n$. However, $\psi_n(H_n(\mathcal{I}_{g,1}))$ also contains a subrepresentation of weight n-2 for all $n\geq 2$ and $g\gg 0$. In particular, $A_n(\mathcal{I}_{g,1}) \subset H_n(\mathcal{I}_{g,1})$ is a proper subrepresentation for $n\geq 2$ and $g\gg 0$.

Remark 1.12. We will see that the map induced on homology by $\mathcal{I}_g^1 \to \mathcal{I}_{g,1}$ maps $A_n\left(\mathcal{I}_g^1\right) \to A_n\left(\mathcal{I}_{g,1}\right)$. From this it follows that the first claim of Theorem 1.11 holds for $\psi_n\left(A_n\left(\mathcal{I}_g^1\right)\right)$ as well. To prove the second claim, however, we construct a homology class that requires a marked point rather than a boundary component, and it is not clear how to construct something analogous using a boundary component. However, it follows from results by Kupers and Randal-Williams [11] that $H_3\left(\mathcal{I}_g^1\right)$ contains an irreducible subrepresentation of weight 1, so it seems reasonable to expect Theorem 1.11 to hold in its entirety for $A_n\left(\mathcal{I}_g^1\right)$ as well. If we also believe that Questions 1–3 have positive answers, it follows from the results by Kupers and Randal-Williams that for large *n* there are many classes in $H_n\left(\mathcal{I}_q^1\right)$ that are of lower weight than *n* and thus cannot lie in $A_n\left(\mathcal{I}_q^1\right)$.

1.5. Structure of the paper

In §2 we give a brief description of *Schur–Weyl duality for symplectic groups*, which is a fundamental result in the representation theory of symplectic groups. We then describe some simple but useful consequences of this theorem.

In §3 we recall some basics from the theory of mapping-class groups in order to define bounding-pair maps. We then introduce abelian cycles and show how ψ_n can be easily evaluated on these. This allows us to give a simple proof of Theorem 1.3.

In §4 we prove Theorem 1.5. Here we see how considering surfaces with a boundary component allows us to define a 'multiplication' of abelian cycles and thereby prove the theorem in a systematic way.

Finally, in $\S5$, we prove Theorem 1.11.

2. Schur–Weyl duality for symplectic groups

We want to investigate how the image of $\psi_n : H_n(\mathcal{I}_g^1) \to \bigwedge^n(\bigwedge^3 H)$ decomposes into irreducible representations of $\operatorname{Sp}(H)$, and to this end we first need to understand how $H^{\otimes k}$ decomposes into irreducible subrepresentations for any $k \geq 1$. This decomposition is well understood, by so-called *Schur–Weyl duality* for symplectic groups. The representation $H^{\otimes k}$ contains irreducible subrepresentation of weight at most k (see §1.3), and the irreducibles of top weight form a subrepresentation which we denote by $H^{\langle k \rangle}$. We call the tensors in $H^{\langle k \rangle}$ traceless, since these are precisely those tensors in $H^{\otimes k}$ that are mapped to zero by all of the contraction maps $C_{i,j}^k : H^{\otimes k} \to H^{\otimes k-2}$, given by contracting with the *i*th and *j*th factors with the symplectic form. The space $H^{\langle k \rangle}$ has both a left action by $\operatorname{Sp}(H)$ and a commuting right action by the symmetric group Σ_k , by permuting the tensor factors. Schur–Weyl duality¹ describes how $H^{\langle k \rangle}$ decomposes as a representation of $\operatorname{Sp}(H) \times \Sigma_k$. The irreducible representations of Σ_k are classified by partitions of weight k, as described in [4, Chapter 4], and for such a partition λ we denote the corresponding irreducible representation of Σ_k by σ_λ . Schur–Weyl duality for symplectic groups, which is described for example in [4, Chapter 17.3], then states that

$$H^{\langle k \rangle} \cong \bigoplus_{|\lambda|=k} V_{\lambda} \otimes \sigma_{\lambda}.$$

Up to isomorphism, each V_{λ} can be constructed as the image of a certain Young symmetrizer $\mathbf{c}_{\lambda} : H^{\langle k \rangle} \to H^{\langle k \rangle}$, which is an element of $\mathbb{Q}[\Sigma_k]$. For a proof of this fact, once again see [4, Chapters 6 and 17.3]. Young symmetrizers are defined using Young diagrams, which are diagrams that represent partitions, as illustrated in Figure 1. A Young tableau is a Young diagram with a choice of numbering of the boxes.

To define \mathbf{c}_{λ} , for $|\lambda| = k$, we first need to pick a Young tableau with an underlying Young diagram representing λ . We denote this tableau by T and let $P_{\lambda} = \{\sigma \in \Sigma_k \mid \sigma \text{ preserves each row of } T\}$ and $Q_{\lambda} = \{\sigma \in \Sigma_k \mid \sigma \text{ preserves each column of } T\}$. We now



Figure 1. The Young diagram representing the partition $(5 \ge 3 \ge 1)$ and a corresponding Young tableau.

¹There is also a more general version that describes the decomposition of all of $H^{\otimes k}$, but this requires more background to describe and is not required for our purposes.

define $\mathbf{c}_{\lambda} := \mathbf{a}_{\lambda} \mathbf{b}_{\lambda} \in \mathbb{Q}[\Sigma_k]$, where

$$\mathbf{a}_{\lambda} := \sum_{\sigma \in P_{\lambda}} \sigma, \qquad \mathbf{b}_{\lambda} := \sum_{\sigma \in Q_{\lambda}} \operatorname{sgn}(\sigma) \sigma.$$

If we take the Young tableau of λ that is numbered analogously to the tableau in Figure 1 and let $\mu = (\mu_1 \geq \cdots \geq \mu_l)$ denote the partition we get by transposing the rows and columns of the Young diagram representing λ , it follows that the image of the corresponding Young symmetrizer $\mathbf{c}_{\lambda} : H^{\langle k \rangle} \to H^{\langle k \rangle}$ is the subrepresentation of

$$\bigwedge^{\mu_1} H \otimes \cdots \otimes \bigwedge^{\mu_l} H$$

generated by the tensor

$$(a_1 \wedge \cdots \wedge a_{\mu_1}) \otimes \cdots \otimes (a_1 \wedge \cdots \wedge a_{\mu_l}),$$

for any choice of symplectic basis $\{a_1, b_1, \ldots, a_g, b_g\}$ of H. From this, it follows that as an $\operatorname{Sp}(H)$ -representation, $V_\lambda \otimes \sigma_\lambda$ is generated by tensors of the form

$$((a_1 \wedge \cdots \wedge a_{\mu_1}) \otimes \cdots \otimes (a_1 \wedge \cdots \wedge a_{\mu_l})) \cdot \sigma$$

for $\sigma \in \Sigma_k$. Here the action is defined by viewing $\bigwedge^{\mu_1} H \otimes \cdots \otimes \bigwedge^{\mu_l} H$ as a subspace of $H^{\otimes k}$ – that is, as the composite map

$$\bigwedge^{\mu_1} H \otimes \cdots \otimes \bigwedge^{\mu_l} H \hookrightarrow H^{\otimes k} \stackrel{\cdot \sigma}{\longrightarrow} H^{\otimes k},$$

where the first map is the standard inclusion. This allows us to prove the following lemma:

Lemma 2.1. For all $g \ge k \ge 1$, the Sp(H)-representation $H^{\langle k \rangle}$ is generated by the tensor

$$a_1 \otimes \cdots \otimes a_k,$$
 (1)

for any fixed symplectic basis $\{a_1, b_1, \ldots, a_g, b_g\}$ of H.

Proof. By the foregoing discussion, it suffices to prove that every tensor A of the form

$$a_{i_1} \otimes \cdots \otimes a_{i_k},$$

for $i_1, \ldots, i_j \in \{1, \ldots, g\}$, lies in the $\mathbb{Q}[\operatorname{Sp}(H)]$ -span of the tensor $a_1 \otimes \cdots \otimes a_k$. We may reorder the symplectic basis in such a way that for each $1 \leq j \leq k$, we have $a_{i_j} = a_j$, unless $a_{i_j} = a_{i_l}$ for some l < j. In other words, we may assume that A agrees with $a_1 \otimes \cdots \otimes a_k$ in every tensor factor except for in those where A has a repeating element.

For $p,q \in \{1,\ldots,g\}$ such that $p \neq q$, let $S_{p,q} \in \text{Sp}(H)$ denote the linear map defined by mapping $a_p \mapsto a_p + a_q$ and $b_q \mapsto b_q - b_p$ and mapping all other basis vectors to themselves. If $a_{i_j} = a_l$ for some l < j, we can act on $a_1 \otimes \cdots \otimes a_k$ with $S_{j,l} - \mathbf{1}$ and get

$$a_1 \otimes \cdots \otimes a_{j-1} \otimes a_l \otimes a_{j+1} \otimes \cdots \otimes a_k.$$

We can do this for every repeated tensor factor of A and thus see that A does indeed lie in the $\mathbb{Q}[\operatorname{Sp}(H)]$ -span of $a_1 \otimes \cdots \otimes a_k$.

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This has the following immediate corollary:

Corollary 2.2. Let V be a direct summand of $H^{\otimes k}$ for $g \geq k \geq 1$, as a representation of Sp(H). Let $p: H^{\otimes k} \to V$ be the projection map and fix a symplectic basis $\{a_1, b_1, \ldots, a_g, b_g\}$ of H. Then the subspace of V spanned by irreducible subrepresentations of weight k lies in the $\mathbb{Q}[Sp(H)]$ -span of the tensor

$$p(a_1 \otimes \cdots \otimes a_k).$$

3. Abelian cycles and bounding-pair maps

Computing the *n*-fold coproduct of a general class $\alpha \in H_n(\mathcal{I}_g^1)$ is difficult, but we will restrict our interest to a simple class of homology classes, called *abelian cycles*, for which we can give an explicit expression. For a group G, abelian cycles in $H_n(G)$ are constructed from commuting elements in the group itself, so in order to construct such classes for \mathcal{I}_g^1 in particular, we will first recall some basic notions from the theory of mapping-class groups.

3.1. Dehn twists and bounding-pair maps

First, recall that if α is a simple, closed curve in $S^s_{g,r}$, then a *Dehn twist* along α is a diffeomorphism of $S^s_{g,r}$ that is obtained as follows: start with some tubular neighborhood N around α , with a diffeomorphism $N \cong S^1 \times I$. The corresponding Dehn twist is then given by the identity outside of N, and on $N \cong S^1 \times I$ by $(s,t) \mapsto (se^{i2\pi t},t)$. Pictorially, we can illustrate this as in Figure 2.

Up to isotopy, the choice of tubular neighborhood does not matter, and if α and α' are homotopic, the corresponding Dehn twists are isotopic. Given a homotopy class α of simple, closed curves in $S_{g,r}^s$, it is thus well defined to talk about *the* Dehn twist along α as an element of $\Gamma_{g,r}^s$. We will denote this element by T_{α} . It is a classic result by Dehn, and later independently by Lickorish, that $\Gamma_{g,r}^s$ is generated by Dehn twists.

Dehn twists themselves are generally not elements of $\mathcal{I}_{g,r}^s$, but it is easy to find compositions of Dehn twists which are. A *bounding pair* in $S_{g,r}^s$ is a pair of simple, closed, nonseparating, and homologous curves. If (α,β) is a bounding pair, we call the composition $T_{\alpha}T_{\beta^{-1}}$ the *bounding-pair map* corresponding to this bounding pair. The typical form of a bounding pair is illustrated in Figure 3.



Figure 2. A Dehn twist.



Figure 3. A typical bounding pair.

Bounding-pair maps are elements of $\mathcal{I}_{g,r}^s$, and in fact, Johnson [9] proved that \mathcal{I}_g is generated by a finite number of bounding-pair maps for $g \geq 3$. This was also used to prove that $\tau : H_1(\mathcal{I}_q^1) \to \bigwedge^3 H$ is an isomorphism in this range.

Any bounding pair (α,β) separates S_g^1 into two connected components. Let S' denote the connected component not containing the original boundary component and let $\{a_1,b_1,\ldots,a_g,b_g\}$ be a symplectic basis of H such that for some $1 \leq k \leq g$, $\{a_1,b_1,\ldots,a_k,b_k,[\alpha]\}$ descends to a basis for $H_1(S')$, where $[\alpha]$ denotes the homology class of α .

Proposition 3.1. ([8, Lemma 4B]). Let (α, β) be a bounding pair and let $\{a_1, b_1, \ldots, a_g, b_g\}$ be a symplectic basis of H as in the previous paragraph. If f denotes the bounding-pair map corresponding to (α, β) , we have

$$\tau(f) = \sum_{i=1}^{k} (a_i \wedge b_i) \wedge [\alpha] \in \bigwedge^{3} H.$$

Note that since α and β belong to the same homology class, this formula is well defined.

3.2. Abelian cycles

If two bounding pairs are disjoint, the corresponding bounding-pair maps commute. This gives us a good way to construct elements in $H_n(\mathcal{I}_g^1)$. If G is a group and (f_1, \ldots, f_n) is an *n*-tuple of pairwise commuting elements in G, this determines a map $\mathbb{Z}^n \to G$, which induces a map $H_n(\mathbb{Z}^n) \to H_n(G)$. We have $H_n(\mathbb{Z}^n) \cong H_n(T^n)$, where T^n denotes the *n*torus, and hence $H_n(\mathbb{Z}^n)$ is generated by its fundamental class. We call the image of this class under the map to $H_n(G)$ the *abelian cycle* determined by (f_1, \ldots, f_n) and denote it by $A(f_1, \ldots, f_n)$.

Lemma 3.2. If $\Delta^{n-1}: H_n(G) \to \bigwedge^n H_1(G)$ is the comultiplication map, then

$$\Delta^{n-1}A(f_1,\ldots,f_n) = [f_1] \wedge \cdots \wedge [f_n],$$

where $[f] \in H_1(G)$ denotes the homology class of $f \in G$.

Proof. The homology $H_*(\mathbb{Z}^n)$ is a bialgebra, and if $\{a_1, \ldots, a_n\}$ is a basis for $H_1(\mathbb{Z}^n)$, the fundamental class of $H_n(\mathbb{Z}^n)$ is $a_1 \cdots a_n$. Thus

$$\Delta^{n-1}(a_1\cdots a_n)=a_1\wedge\cdots\wedge a_n.$$

Furthermore, the map $H_1(\mathbb{Z}^n) \to H_1(G)$ maps $a_i \mapsto [f_i]$, for $1 \leq i \leq n$. By functoriality of homology, we have a commutative diagram

$$\begin{array}{cccc}
H_n(\mathbb{Z}^n) & \longrightarrow & H_n(G) \\
\downarrow & & \downarrow \\
\bigwedge^n H_1(\mathbb{Z}^n) & \longrightarrow & \bigwedge^n H_1(G),
\end{array}$$

which means that

$$\Delta^{n-1}A(f_1,\ldots,f_n)=[f_1]\wedge\cdots\wedge[f_n].$$

If (f_1, \ldots, f_n) is an *n*-tuple of bounding pairs corresponding to pairwise disjoint bounding pairs, we will simply call this an *n*-tuple of disjoint bounding-pair maps. We will consider abelian cycles in $H_n(\mathcal{I}_{g,r}^s)$ that are determined by such *n*-tuples and denote by $A_n(\mathcal{I}_{g,r}^s)$ the subrepresentation of $H_n(\mathcal{I}_{g,r}^s)$ generated by these.

Remark 3.3. It is clear from the definition of a bounding-pair map that the map $\mathcal{I}_g^1 \to \mathcal{I}_{g,1}$ induces a map $A_n(\mathcal{I}_q^1) \to A_n(\mathcal{I}_{g,1})$.

If $A(f_1,\ldots,f_n) \in H_n(\mathcal{I}^1_q)$ is an abelian cycle, it follows by Lemma 3.2 that

 $\psi_n(A(f_1,\ldots,f_n)) = \tau([f_1]) \wedge \cdots \wedge \tau([f_n]).$

To evaluate ψ_n on an abelian cycle, this means that we only need Johnson's original result on the image of τ .

Example 3.4. Consider the abelian cycle A determined by the bounding pairs in Figure 4. With an appropriate choice of symplectic basis of H, we see that the image of this abelian cycle is

$$\psi_n(A) = (a_1 \wedge b_1 \wedge a_3) \wedge \dots \wedge (a_{3n-2} \wedge b_{3n-2} \wedge a_{3n}).$$

We can now use this to prove Theorem 1.3:

Proof of Theorem 1.3. Since $\bigwedge^n (\bigwedge^3 H)$ is a direct summand of $H^{\otimes 3n}$, it follows by Corollary 2.2 that it suffices to show that the tensor

$$(a_1 \wedge a_2 \wedge a_3) \wedge \cdots \wedge (a_{3n-2} \wedge a_{3n-1} \wedge a_{3n})$$

lies in the image of ψ_n . To this end, we consider the abelian cycle from Example 3.4. For $i \neq j \in \{1, \ldots, g\}$ let $T_{i,j} : H \to H$ denote the linear map defined by mapping $b_i \mapsto b_i + a_j$ and $b_j \mapsto b_j + a_i$ and mapping all other basis vectors to themselves. Each such map lies



Figure 4. Bounding pairs determining an abelian cycle in $H_n(\mathcal{I}_q^1)$, for $g \geq 3n$.

in Sp(H), and we see that if we act with $\prod_{i=1}^{n} (T_{3i-2,3i-1} - 1) \in \mathbb{Q}[Sp(H)]$ on

$$(a_1 \wedge b_1 \wedge a_3) \wedge \cdots \wedge (a_{3n-2} \wedge b_{3n-2} \wedge a_{3n}),$$

we get

$$(a_1 \wedge a_2 \wedge a_3) \wedge \dots \wedge (a_{3n-2} \wedge a_{3n-1} \wedge a_{3n})$$

4. Proof of Theorem 1.5

Let us start by restating Theorem 1.5:

Theorem 1.5. Set $n \ge 1$ and let $\lambda = (\lambda_1^{k_1} > \cdots > \lambda_m^{k_m} > 0)$ and $\mu = (\mu_1^{l_1} > \cdots > \mu_{m+2}^{l_{m+2}} > 0)$ be partitions such that $\mu_i = \lambda_i + 2$ for $i = 1, \ldots, m$ and $|\lambda| + |\mu| = n$. Let $k_{m+1} = k_{m+2} = 0$, $k = \sum_{i=1}^{m} k_i$, and $l = \sum_{i=1}^{m+2} l_i$, and suppose that $g \ge n + 2k + l$. Then the subrepresentation of

$$\bigotimes_{i=1}^{m+2} \bigwedge^{k_i+l_i} V_{1^{\mu_i}}$$

spanned by all irreducible subrepresentations of weight n+2k is contained in $Im(\psi_n)$.

The idea for proving this is essentially to consider two simple kinds of abelian cycles in $H_n(\mathcal{I}_{g,1})$, whose respective images have components in $V_{1^{n+2}}$ and V_{1^n} , and then taking *products* of such cycles, for a range of *ns* corresponding to the partitions. The special case of Theorem 1.3 is given by taking the *n*th power of a bounding pair in $H_1(\mathcal{I}_3^1)$ whose image has a component in V_{1^3} (see Example 4.2).

4.1. Products of abelian cycles

Given mapping classes $f_1 \in \mathcal{I}_{g_1}^1$ and $f_2 \in \mathcal{I}_{g_2}^1$, we may glue the surfaces $S_{g_1}^1$ and $S_{g_2}^1$ along the dashed segments illustrated in Figure 5, and since both mapping classes fix



Figure 5. Gluing the surfaces along the dashed segments induces the map $H_i(\mathcal{I}_{g_1}^1) \times H_j(\mathcal{I}_{g_1}^1) \to H_{i+j}(\mathcal{I}_{g_1+g_2}^1)$.

the boundaries, this naturally gives us a mapping class in $\mathcal{I}^1_{g_1+g_2}$. Thus we get a map $\mathcal{I}^1_{g_1} \times \mathcal{I}^1_{g_2} \to \mathcal{I}^1_{g_1+g_2}$, which induces a multiplication

$$H_{i}\left(\mathcal{I}_{g_{1}}^{1}\right) \otimes H_{j}\left(\mathcal{I}_{g_{2}}^{1}\right) \to H_{i+j}\left(\mathcal{I}_{g_{1}+g_{2}}^{1}\right)$$

by precomposing with the Künneth map. This makes the space $\bigoplus_{g,n\geq 0} H_n(\mathcal{I}_g^1)$ into a bigraded associative and commutative algebra.

In $H_{i+j}(\mathbb{Z}^{i+j}) \cong H_i(\mathbb{Z}^i) \otimes H_j(\mathbb{Z}^j)$, the fundamental class is the tensor product of the fundamental classes in $H_i(\mathbb{Z}^i)$ and $H_j(\mathbb{Z}^j)$. This proves the following proposition:

Proposition 4.1. If $A(f_1, \ldots, f_i)$ and $A(f_{i+1}, \ldots, f_{i+j})$ are abelian cycles in $H_i(\mathcal{I}_{g_1}^1)$ and $H_j(\mathcal{I}_{g_2}^1)$, respectively, their product is the abelian cycle $A(f_1, \ldots, f_{i+j}) \in H_{i+j}(\mathcal{I}_{g_1+g_2}^1)$.

Example 4.2. With this in mind, we see that the abelian cycle depicted in Figure 4 is simply the product $\sigma_1^n \cdot 1_{g-3n}$, where σ_1 is the class in $H_1(\mathcal{I}_3^1)$ determined by the bounding pair in Figure 6, and $1_{g-3n} \in H_0(\mathcal{I}_{q-3n}^1)$ is the generator.

4.2. Families of truly nested bounding pairs

With Example 4.2 in mind, we will consider abelian cycles determined by families of bounding pairs of the form depicted in Figure 7 and then their products. We call these families of bounding pairs *truly nested*, a terminology that we borrow from [1].

Definition 4.3. A family $(\alpha_1, \beta_1), \ldots, (\alpha_n, \beta_n)$ of bounding pairs is truly nested if

(1) the curves α_i are pairwise nonhomologous and

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Figure 7. Truly nested families of bounding pairs defining the abelian cycles $\sigma_n \in H_n\left(\mathcal{I}_{n+2}^1\right)$ and $\rho_n \in H_n\left(\mathcal{I}_{n+1}^1\right)$ considered in Example 4.5.

(2) after possibly reordering the bounding pairs, cutting the surface along (α_j, β_j) separates the bounding pair (α_i, β_i) from the original boundary component whenever i < j.

We will always assume that the bounding pairs in such a family are numbered so that Definition 4.3(2) is met. By induction it can be seen that a truly nested family of bounding pairs is always of the form depicted in Figure 7, but possibly with holes between each adjacent couple of bounding pairs.

If we are given a truly nested family of bounding pairs $(\alpha_1, \beta_1), \ldots, (\alpha_n, \beta_n)$ and we cut the surface along a bounding pair, we can talk about the resulting connected component that is 'furthest' from the original boundary component. If the surface is cut along (α_1, β_1) , consider that component and let ω_0 denote the symplectic form of a maximal subsurface with only one boundary component. In order to prove Theorem 1.5, we will need the following lemma:

Lemma 4.4. Let $A(f_1, \ldots, f_n)$ be an abelian cycle determined by a truly nested family of bounding-pair maps. Then there exists an Sp(H)-equivariant map $\Phi_n : \bigwedge^n \left(\bigwedge^3 H\right) \to \bigwedge^{n+2} H$, defined by a sequence of n-1 contractions, such that

$$(\Phi_n \circ \psi_n)(A(f_1, \dots, f_n)) = \lambda_n \omega_0 \wedge c_1 \wedge \dots \wedge c_n,$$
(2)

for some nonzero integer λ_n , where c_i is the homology class of the bounding pair corresponding to f_i and ω_0 is as already defined.

We construct the map Φ_n as follows: for n = 1 we set $\Phi_1 = \mathrm{id}_{\bigwedge^3 H}$. For $n \ge 2$, we have a map $\phi_1^n : \bigwedge^n \left(\bigwedge^3 H\right) \to \bigwedge^4 H \otimes \bigwedge^{n-2} \left(\bigwedge^3 H\right)$, given by contracting 'diagonally' with the symplectic form and composing with the multiplication map $\bigwedge^2 \left(\bigwedge^2 H\right) \to \bigwedge^4 H$. For any $2 \le k \le n-1$, we similarly have a composite map

$$\bigwedge^{k+2} H \otimes \bigwedge^{n-k} \left(\bigwedge^{3} H \right) \to \bigwedge^{k+1} H \otimes \bigwedge^{2} H \otimes \bigwedge^{n-k-1} \left(\bigwedge^{3} H \right) \to \bigwedge^{k+3} H \otimes \bigwedge^{n-k-1} \left(\bigwedge^{3} H \right),$$

which we denote by ϕ_k^n . We define $\Phi_n : \bigwedge^n \left(\bigwedge^3 H\right) \to \bigwedge^{n+2} H$ as the composition $\phi_{n-1}^n \circ \cdots \circ \phi_1^n$.

The calculations needed to prove the formula in Lemma 4.4 are simple but tedious and not very enlightening, and we will therefore leave them until the end of this section. Before proving the theorem, let us consider the following example, which illustrates the idea of the proof:

Example 4.5. Let $\sigma_n \in H_n(\mathcal{I}_{n+2}^1)$ denote the abelian cycle determined by the bounding pairs in Figure 7(a). With an appropriate choice of symplectic basis, it follows from Lemma 4.4 that

$$(\Phi_n \circ \psi_n)(\sigma_n) = \lambda_n a_1 \wedge b_1 \wedge a_3 \wedge \dots \wedge a_{n+2},$$

for some nonzero integer λ_n . As in the proof of Theorem 1.3, we can now act by $T_{1,2} - \mathbf{1}$, which gives us

$$\lambda_n a_1 \wedge a_2 \wedge \cdots \wedge a_{n+2}.$$

From Corollary 2.2, we can now draw the conclusion that $V_{1^{n+2}}$ lies in the image, since this is the top weight representation of $\bigwedge^{n+2} H$.

Similarly, let $\rho_n \in H_n(\mathcal{I}_{n+1}^1)$ denote the cycle determined by bounding pairs as in Figure 7(b). Let Ψ_n denote the composition of Φ_n with the contraction map $\bigwedge^{n+2} H \to \bigwedge^n H$. If we once again order the basis appropriately, we see that

$$(\Psi_n \circ \psi_n)(\rho_n) = \lambda_n a_1 \wedge a_2 \wedge \dots \wedge a_n,$$

and hence V_{1^n} lies in the image. We can thus draw the conclusion that if $g \ge n+2$, then $V_{1^{n+2}} \oplus V_{1^n}$ are contained in the image of ψ_n , whereas if $g \ge n$, V_{1^n} is contained in the image.

With λ and μ as in Theorem 1.5, we see that for $i = 1, \ldots, m$, each of σ_{λ_i} and ρ_{μ_i} gives us a copy of $V_{1^{\lambda_i+2}} = V_{1^{\mu_i}}$ in $\psi_{\lambda_i} \left(H_{\lambda_i} \left(\mathcal{I}^1_{\lambda_i+2} \right) \right)$ and $\psi_{\mu_i} \left(H_{\mu_i} \left(\mathcal{I}^1_{\mu_i+1} \right) \right)$, respectively. The weight of these irreducibles is at least 3. Meanwhile, $\rho_{\mu_{m+1}}$ and $\rho_{\mu_{m+2}}$ give us copies of $V_{1^{\mu_{m+1}}}$ and $V_{1^{\mu_{m+2}}}$ in $\psi_{\mu_{m+1}} \left(H_{\mu_{m+1}} \left(\mathcal{I}^1_{\mu_{m+1}+1} \right) \right)$ and $\psi_{\mu_{m+2}} \left(H_{\mu_{m+2}} \left(\mathcal{I}^1_{\mu_{m+2}+1} \right) \right)$, respectively, and here the weight may be lower.

The idea for Theorem 1.5 is therefore to consider the product $\sigma_{\lambda_1}^{k_1}\rho_{\mu_1}^{l_1}\cdots\sigma_{\lambda_m}^{k_m}\rho_{\mu_m}^{l_m}\rho_{\mu_{m+1}}^{l_{m+1}}\rho_{\mu_{m+2}}^{l_{m+1}}\cdot\mathbf{1}_{g-n-2k-l}$. This is what gives us the condition that

$$g \ge \sum_{i=1}^{m+2} (k_i(\lambda_i+2) + l_i(\mu_i+1)) = |\lambda| + |\mu| + 2k + l = n + 2k + l.$$

4.3. Proof of the theorem

Using the foregoing, we may now prove the theorem:

Proof of Theorem 1.5, assuming Lemma 4.4. As stated previously, we will use the abelian cycle

$$\sigma := \sigma_{\lambda_1}^{k_1} \rho_{\mu_1}^{l_1} \cdots \sigma_{\lambda_m}^{k_m} \rho_{\mu_m}^{l_m} \rho_{\mu_{m+1}}^{l_{m+1}} \rho_{\mu_{m+2}}^{l_{m+1}} \cdot \mathbf{1}_{g-n-2k-l},$$

where σ_{λ_i} and ρ_{μ_i} are defined as in Example 4.5. For $1 \le i \le m+2$ and $1 \le p \le k_i$, we let $H_{\sigma_{\lambda_i}}^{(p)}$ denote the first homology of the subsurface corresponding to the *p*th σ_{λ_i} -factor, with $H_{\sigma_{\lambda_{m+1}}}^{(p)}$ and $H_{\sigma_{\lambda_{m+2}}}^{(p)}$ defined as 0 for all *p*. We define $H_{\rho_{\mu_i}}^{(q)}$ similarly, so that

$$H = \bigoplus_{i=1}^{m+2} \left(\left(\bigoplus_{p=1}^{k_i} H_{\sigma_{\lambda_i}}^{(p)} \right) \oplus \left(\bigoplus_{q=1}^{l_i} H_{\rho_{\mu_i}}^{(q)} \right) \right) \oplus H_1 \left(S_{g-n-2k-l} \right).$$

By Proposition 3.1 we have

$$\psi_n(\sigma) \in \bigotimes_{i=1}^{m+2} \left(\left(\bigotimes_{p=1}^{k_i} \bigwedge^{\lambda_i} \bigwedge^3 H_{\sigma_{\lambda_i}}^{(p)} \right) \otimes \left(\bigotimes_{q=1}^{l_i} \bigwedge^{\mu_i} \bigwedge^3 H_{\rho_{\mu_i}}^{(q)} \right) \right) \subseteq \bigwedge^n \bigwedge^3 H.$$

Define $F := \Phi_{\lambda_1}^{\wedge k_1} \wedge \Psi_{\mu_1}^{\wedge l_1} \wedge \cdots \wedge \Phi_{\lambda_m}^{\wedge k_m} \wedge \Psi_{\mu_m}^{\wedge l_m} \wedge \Psi_{\mu_{m+1}}^{\wedge l_{m+1}} \wedge \Psi_{\mu_{m+2}}^{\wedge l_{m+2}}$. Since the direct summands $H_{\sigma_{\lambda_i}}^{(p)}$ and $H_{\rho_{\mu_i}}^{(q)}$ of H are pairwise orthogonal with respect to the symplectic form ω , it follows from the definitions of Φ_{λ_i} and Ψ_{λ_i} that

$$F(\psi_n(\sigma)) \in \bigotimes_{i=1}^{m+2} \left(\left(\bigotimes_{p=1}^{k_i} \bigwedge^{\mu_i} H_{\sigma_{\lambda_i}}^{(p)} \right) \otimes \left(\bigotimes_{q=1}^{l_i} \bigwedge^{\mu_i} H_{\rho_{\mu_i}}^{(q)} \right) \right) \subseteq \bigotimes_{i=1}^{m+2} \bigwedge^{k_i+l_i} \left(\bigwedge^{\mu_i} H \right).$$

Since the top weight irreducibles of $\bigotimes_{i=1}^{m+2} \bigwedge^{k_i+l_i} (\bigwedge^{\mu_i} H)$ are precisely the top weight irreducibles of $\bigotimes_{i=1}^{m+2} \bigwedge^{k_i+l_i} V_{1^{\mu_i}}$, we now want to apply Corollary 2.2. The representation $\bigotimes_{i=1}^{m+2} \bigwedge^{k_i+l_i} (\bigwedge^{\mu_i} H)$ is a direct summand of $H^{\otimes (n+2k)}$. Let us denote the projection map

by π . If we choose a symplectic basis for each $H_{\sigma_{\lambda_i}}^{(p)}$ and each $H_{\sigma_{\mu_i}}^{(q)}$, extend the union of these to a symplectic basis $\{a_1, b_1, \ldots, a_g, b_g\}$ for H, and order this basis appropriately, we have

$$\pi(a_1 \otimes \dots \otimes a_{n+2k}) \in \bigotimes_{i=1}^{m+2} \left(\left(\bigotimes_{p=1}^{k_i} \bigwedge^{\mu_i} H^{(p)}_{\sigma_{\lambda_i}} \right) \otimes \left(\bigotimes_{q=1}^{l_i} \bigwedge^{\mu_i} H^{(q)}_{\rho_{\mu_i}} \right) \right) \subseteq \bigotimes_{i=1}^{m+2} \bigwedge^{k_i+l_i} \left(\bigwedge^{\mu_i} H \right).$$

From what we saw in Example 4.5, there is an element

$$T \in \mathbb{Q}\left[\left(\prod_{i=1}^{m+2} \left(\prod_{p=1}^{k_i} \operatorname{Sp}\left(H_{\sigma_{\lambda_i}}^{(p)}\right)\right) \times \left(\prod_{q=1}^{l_i} \operatorname{Sp}\left(H_{\rho_{\mu_i}}^{(q)}\right)\right)\right) \times \operatorname{Sp}\left(H_1\left(S_{g-n-2k-l}\right)\right)\right] \subseteq \mathbb{Q}[\operatorname{Sp}(H)]$$

such that

$$T \cdot (F(\psi_n(\sigma))) = \pi(a_1 \otimes \cdots \otimes a_{n+2k}).$$

4.4. Proof of Lemma 4.4

Now let us prove the lemma.

Proof of Lemma 4.4. We prove equation (2) by induction on *n*. For n = 1 we have $\psi_1 = \tau$ and $\Phi_1 = \operatorname{Id}_{\Lambda^3 H}$, so the identity holds with $\lambda_1 = 1$, by Proposition 3.1.

Now suppose $n \geq 2$ and that f_1, \ldots, f_n are bounding-pair maps that correspond to truly nested bounding pairs $(\alpha_1, \beta_1), \ldots, (\alpha_n, \beta_n)$. We may pick a symplectic basis $\{a_1, b_1, \ldots, a_g, b_g\}$ such that $[\alpha_k] = a_{i_k}$ for $1 \leq k \leq n$, where $i_1 \leq \cdots \leq i_n$ and such that $a_1 \wedge b_1 + \cdots + a_{i_k-1} \wedge b_{i_k-1}$ descends to a symplectic form of the homology of the surface not containing the original boundary component that we get by cutting along (α_k, β_k) . For n = 2, a simple calculation shows that

$$\begin{split} \Phi_2 \circ \psi_2(A(f_1, f_2)) \\ &= \phi_1^2 \Big(\Big((a_1 \wedge b_1 + \dots + a_{i_1-1} \wedge b_{i_1-1}) \wedge a_{i_1} \Big) \wedge \Big((a_1 \wedge b_1 + \dots + a_{i_2-1} \wedge b_{i_2-1}) \wedge a_{i_2} \Big) \Big) \\ &= -3 \left(a_1 \wedge b_1 + \dots + a_{i_1-1} \wedge b_{i_1-1} \right) \wedge a_{i_1} \wedge a_{i_2}. \end{split}$$

Now we let $n \geq 3$ and assume that

$$\Phi_{n-1} \circ \psi_{n-1} \left(A \left(f_1, \dots, \hat{f}_i, \dots, f_n \right) \right)$$

= $\lambda_{n-1} \left(a_1 \wedge b_1 + \dots + a_{i_1-1} \wedge b_{i_1-1} \right) \wedge a_{i_1} \wedge \dots \wedge \hat{a}_{i_k} \wedge \dots \wedge a_{i_n}$

and

$$\Phi_{n-1} \circ \psi_{n-1}(A(f_2, \dots, f_n)) = \lambda_{n-1} (a_1 \wedge b_1 + \dots + a_{i_2-1} \wedge b_{i_2-1}) \wedge a_{i_2} \wedge \dots \wedge a_{i_n}$$

for some nonzero $\lambda_{n-1} \in \mathbb{Z}$. Since

$$\Phi_n = \phi_{n-1}^n \circ \left(\Phi_{n-1} \wedge \operatorname{Id}_{\bigwedge^3 H} \right)$$

and

$$\psi_n(A(f_1,\ldots,f_n)) = (-1)^{n-k} \psi_{n-1}\left(A\left(f_1,\ldots,\hat{f}_k,\ldots,f_n\right) \wedge \tau([f_k])\right)$$

for each $k = 1, \ldots, n$, it follows that

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$$\begin{split} & {}_{n} \circ \psi_{n}(A(f_{1}, \dots, f_{n})) = \lambda_{n-1} \cdot \phi_{n-1}^{n} \\ & \times \left((-1)^{n-1} \left(a_{1} \wedge b_{1} + \dots + a_{i_{2}-1} \wedge b_{i_{2}-1} \right) \wedge a_{i_{2}} \wedge \dots \wedge a_{i_{n}} \\ & \otimes \left(a_{1} \wedge b_{1} + \dots + a_{i_{1}-1} \wedge b_{i_{1}-1} \right) \wedge a_{i_{1}} \\ & + \sum_{k=2}^{n} (-1)^{n-k} \left(a_{1} \wedge b_{1} + \dots + a_{i_{1}-1} \wedge b_{i_{1}-1} \right) \wedge a_{i_{1}} \wedge \dots \wedge \hat{a}_{i_{k}} \wedge \dots \wedge a_{i_{n}} \\ & \otimes \left(a_{1} \wedge b_{1} + \dots + a_{i_{k}-1} \wedge b_{i_{k}-1} \right) \wedge a_{i_{k}} \right). \end{split}$$

A simple calculation shows that

$$\phi_{n-1}^n \left((a_1 \wedge b_1 + \dots + a_{i_2-1} \wedge b_{i_2-1}) \wedge a_{i_2} \wedge \dots \wedge a_{i_n} \otimes (a_1 \wedge b_1 + \dots + a_{i_1-1} \wedge b_{i_1-1}) \wedge a_{i_1} \right) \\ = 3(-1)^n \left(a_1 \wedge b_1 + \dots + a_{i_1-1} \wedge b_{i_1-1} \right) \wedge a_{i_1} \wedge \dots \wedge a_{i_n},$$

and similarly

$$\phi_{n-1}^{n}((a_{1} \wedge b_{1} + \dots + a_{i_{1}-1} \wedge b_{i_{1}-1}) \wedge a_{i_{1}} \wedge \dots \wedge \hat{a}_{i_{k}} \wedge \dots \wedge a_{i_{n}}$$

$$\otimes (a_{1} \wedge b_{1} + \dots + a_{i_{k}-1} \wedge b_{i_{k}-1}) \wedge a_{i_{k}})$$

$$= (-1)^{n-k-1}(k+1)(a_{1} \wedge b_{1} + \dots + a_{i_{1}-1} \wedge b_{i_{1}-1}) \wedge a_{i_{1}} \wedge \dots \wedge a_{i_{n}}.$$

Thus we have

$$\Phi_n \circ \psi_n(A(f_1, \dots, f_n)) = -\lambda_{n-1} \sum_{k=0}^n (k+1) (a_1 \wedge b_1 + \dots + a_{i_1-1} \wedge b_{i_1-1}) \wedge a_{i_1} \wedge \dots \wedge a_{i_n}$$
$$= -\lambda_{n-1} \frac{(n+1)(n+2)}{2} \tau_n(A(f_1, \dots, f_n)),$$

which proves the formula.

Remark 4.6. The idea to consider truly nested families of bounding pairs is taken from [1]. There the authors consider a family of $\operatorname{Sp}(H_{\mathbb{Z}})$ -equivariant maps $\tau_n : H_n(\mathcal{I}_{g,1}) \to \bigwedge^{n+2} H$ that generalize the map $\tau : H_1(\mathcal{I}_{g,1}) \to \bigwedge^3 H$ to higher degree. We can note that τ_n and $\Phi_n \circ \psi_n$ have the same codomain. In [1, §3], the image of τ_n on abelian cycles determined by truly nested families of bounding-pair maps is proven to satisfy

$$\tau_n(A(f_1,\ldots,f_n)) = \omega_0 \wedge c_1 \wedge \cdots \wedge c_n,$$

with the notation being the same as in our Lemma 4.4. We can thus see that on such abelian cycles, the two maps agree, up to multiplication by a nonzero scalar. An interesting question is whether these maps agree on all of $H_n(\mathcal{I}_{g,1})$, so that the image of τ_n is simply a subrepresentation of the image of ψ_n , or whether there is some other way to explicitly relate ψ_n and τ_n .

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5. Proof of Theorem 1.11

As we have seen, abelian cycles determined by disjoint bounding-pair maps are a useful and simple tool for computing big parts of the image of ψ_n . In this section we will give an explicit limitation on the images of such cycles. More specifically, we will prove the following theorem:

Theorem 1.11. Set $n \geq 2$ and let $A_n(\mathcal{I}_{g,1}) \subset H_n(\mathcal{I}_{g,1})$ be the subrepresentation generated by abelian cycles determined by pairwise disjoint bounding-pair maps. Then $\psi_n(A_n(\mathcal{I}_{g,1}))$ is concentrated in weights $n, n+2, \ldots, 3n$. However, $\psi_n(H_n(\mathcal{I}_{g,1}))$ also contains a subrepresentation of weight n-2 for all $n\geq 2$ and $g\gg 0$. In particular, $A_n(\mathcal{I}_{g,1}) \subset H_n(\mathcal{I}_{g,1})$ is a proper subrepresentation for $n\geq 2$ and $g\gg 0$.

We will prove the following two claims of the theorem separately:

Claim 5.1. The image $\psi_n(A_n(\mathcal{I}_{q,1}))$ is concentrated in weights $n, n+2, \ldots, 3n$.

Claim 5.2. For $n \ge 2$ and $g \gg 0$, the image of $\psi_n(H_n(\mathcal{I}_{g,1}))$ contains an irreducible subrepresentation of weight n-2.

To prove Claim 5.1, we will need to find a description of the tensor $\psi_n(A(f_1,\ldots,f_n))$ that we can work with, for any $A(f_1,\ldots,f_n) \in A_n(\mathcal{I}_{g,1})$. This will be done using the following lemma:

Lemma 5.3. Set $A(f_1, \ldots, f_n) \in A_n(\mathcal{I}_{g,1})$. Then there exists a Lagrangian subspace $L \subset H$, such that for $k = 1, \ldots, n$, we have

$$\tau([f_k]) \in Im\left(\bigwedge^2 L \otimes H \to \bigwedge^3 H\right).$$

Proof. Let us denote the bounding pair corresponding to f_k by (α_k, β_k) , for k = 1, ..., n. Note that since the bounding pairs are pairwise disjoint, the span of $\{[\alpha_1], \ldots, [\alpha_n]\}$ is an isotropic subspace of H, which we may denote by I. Recalling Proposition 3.1, we want to extend this to a Lagrangian subspace L in such a way that after the surface is cut by (α_k, β_k) , for any k, L restricts to a Lagrangian subspace of the first homologies of both of the closed surfaces obtained by filling in the boundary components with open disks.

Now let us cut the surface along all the bounding pairs and sew in open disks in all resulting boundary components, in order to get n+1 closed surfaces. For each such closed surface we may pick curves whose homology classes span a Lagrangian subspace of its first homology. These curves can be chosen so that they do not intersect the sewn-in disks or their boundaries and thus, viewed as curves in S_g they do not intersect any of our bounding pairs. This means that taking the span of the homology classes of these curves and all bounding pairs will still be an isotropic subspace of H. Let us denote this subspace by L. To see that L is Lagrangian, it suffices to note that the total genus of the subsurfaces we get by cutting along all bounding pairs is $g - \dim I$, which means that this is also the dimension of the subspace of H spanned by the chosen curves. Since their classes are linearly independent of the classes in I, it follows that dim L = g.

By the construction of L through a selection of pairwise nonintersecting curves, it is clear that if the surface is cut along one bounding pair, the remaining curves all end up in only one of the connected components and are still pairwise nonintersecting. Since the genera of the resulting two subsurfaces add to g-1, it follows that the span of the homology classes of the curves in either boundary component still form Lagrangian subspaces in their respective components, and hence L has the sought-after property.

This means that when the surface is cut along the bounding pair (α_k, β_k) , for any k = 1, ..., n, we can find a symplectic basis for the first homology of the closed surface corresponding to the connected component not containing the original boundary component, such that its symplectic forms can be written

$$\omega_1 = \sum_{i=1}^{g_1} a_i \wedge b_i,$$

with $a_i \in L$. It follows by Proposition 3.1 that

 $\tau([f_k]) = \omega_1 \wedge [\alpha_k],$

which lies in $\operatorname{Im}\left(\bigwedge^2 L \otimes H \to \bigwedge^3 H\right)$, since $[\alpha_k] \in L$.

Using this we may now prove the claim:

Proof of Claim 5.1. By Lemma 5.3, we can find a Lagrangian $L \in H$ such that

$$\psi_n(A(f_1,\ldots,f_n)) \in \operatorname{Im}\left(\bigwedge^n \left(\bigwedge^2 L \otimes H\right) \to \bigwedge^n \left(\bigwedge^3 H\right)\right).$$

This means that if we apply the standard inclusion $\iota : \bigwedge^n \left(\bigwedge^3 H\right) \hookrightarrow H^{\otimes 3n}$, each term of $\psi_n(A(f_1,\ldots,f_n))$ lies in the subspace

$$L^{\otimes 2n} \otimes H^{\otimes n}$$

up to a permutation of the factors. From this it follows that $\iota(\psi_n(A(f_1,\ldots,f_n)))$ must be mapped to zero by any composition of more than *n* contractions, and thus it cannot have a component in an irreducible subrepresentation of weight less than *n*.

Remark 5.4. We have stated Claim 5.1 in the pointed case, but it holds in the unpointed case as well. To see this, define $U := \left(\bigwedge^3 H\right)/H$, let $A_n(\mathcal{I}_g) \subseteq H_n(\mathcal{I}_g)$ be defined similarly as $A_n(\mathcal{I}_{g,1})$, and let $\psi'_n : H_n(\mathcal{I}_g) \to \bigwedge^n U$ be defined similarly as ψ_n . We then have the commutative diagram

$$\begin{array}{ccc} A_n(\mathcal{I}_{g,1}) & \longrightarrow & H_n(\mathcal{I}_{g,1}) & \stackrel{\psi_n}{\longrightarrow} & \bigwedge^n \bigwedge^3 H \\ \downarrow & & \downarrow & & \downarrow \\ A_n(\mathcal{I}_g) & \longmapsto & H_n(\mathcal{I}_g) & \stackrel{\psi'_n}{\longrightarrow} & \bigwedge^n U, \end{array}$$

which proves that $\psi'_n(A_n(\mathcal{I}_g))$ cannot contain any irreducibles of weight less than *n*. In [14], Sakasai uses elements of $A_3(\mathcal{I}_g)$ to determine the stable image of (the dual map of)

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 ψ'_3 , but his result leaves open whether or not the image contains a copy of V_1 .² Claim 5.1 illustrates why this could not be determined using Sakasai's method.

Now it only remains to prove Claim 5.2. Before proving this, we need some additional background. Let \mathcal{T}_g denote the *Torelli space*, the moduli space of Riemann surfaces of genus g endowed with a homology marking – that is, a choice of basis for the first homology group. This is the quotient of the Teichmüller space Teich_g by the free action of \mathcal{I}_g , which means that $H_*(\mathcal{I}_g) \cong H_*(\mathcal{T}_g)$. In a similar way, we can construct the Torelli space $\mathcal{T}_{g,1}$ of pointed Riemann surfaces. The map given by forgetting the marked point, which we will denote by π , makes

$$S_g \to \mathcal{T}_{g,1} \xrightarrow{\pi} \mathcal{T}_g$$

into a universal S_q -bundle over \mathcal{T}_q . We will use this to prove the claim.

Proof of Claim 5.2. First we consider the case n = 2. Let $[*] \in H_0(\mathcal{T}_g)$ be the class given by choosing a point $* \hookrightarrow \mathcal{T}_g$. Pulling this map back along the universal bundle, we get the trivial bundle



which by integration over the fibers gives us a class $\pi^!([*])$ in $H_2(S_g)$, which is just the fundamental class $[S_g]$. By abuse of notation, we will use $[S_g]$ to denote the corresponding class $\pi^!([*]) \in H_2(\mathcal{T}_{g,1}) \cong H_2(\mathcal{I}_{g,1})$ as well.

To compute $\psi_2([S_g])$, we may compute its coproduct in $H_2(S_g)$ and then use functoriality. The coproduct can be computed using the diagram

$$\begin{array}{c} H_2(S_g) \longrightarrow H_2(Z) \\ \downarrow \qquad \qquad \downarrow \\ \bigwedge^2 H_1(S_g) \longrightarrow \bigwedge^2 H_1(Z) \end{array}$$

where Z is the wedge of tori in Figure 8, that we get from S_g by collapsing a subsurface of genus 0 with g boundary components to a point.

The horizontal arrow on the second row of the diagram is an isomorphism, so we only need to compute the coproduct of the fundamental class [Z], which is equal to the sum $[Z_1] + \cdots + [Z_g] \in H_2(Z) \cong H_2(Z_1) \oplus \cdots \oplus H_2(Z_g)$ of the fundamental classes of the g tori. Since the homology of each torus is a bialgebra, we get $\Delta([Z_k]) = a_k \wedge b_k$, where $\{a_k, b_k\}$ is a symplectic basis for $H_1(Z_k)$. The union of these bases corresponds in a natural way to a symplectic basis $\{a_1, b_1, \ldots, a_g, b_g\}$ of $H_1(S_g)$, so we get that $\Delta([S_g]) = \omega$, where ω is the symplectic form of $H_1(S_g)$.

²He also uses one abelian cycle of a different kind, but whose image under ψ'_3 can easily be generated using elements of $A_3(\mathcal{I}_g)$, as it is concentrated in top weight.



Figure 8. The quotient surface obtained from S_g by collapsing a subsurface of genus 0 with g boundary components to a point.

The map $S_g \to \mathcal{T}_{g,1}$ induces the so called 'point-pushing' map $\pi_1(S_g) \to \mathcal{I}_{g,1}$ [3, Chapters 4.2 and 6.4.2], and using this it follows that if $c \in H_1(S_g)$ and we denote its image by $[c] \in H_1(\mathcal{I}_{g,1})$, then $\tau([c]) = \omega \wedge c \in \bigwedge^3 H$. Thus it follows by functoriality that $\psi_2([S_g])$ is the image of ω by the map $\bigwedge^2 H \to \bigwedge^2 (\bigwedge^3 H)$, given by inserting the symplectic form in each factor. The class $\omega \in \bigwedge^2 H$ generates the trivial representation, which means that $\psi_2([S_g])$ has weight 0. This proves the claim in the case n = 2.

The case n = 3 has already been treated by Kupers and Randal-Williams [11, §8], who proved that for $g \gg 0$, the image of the dual of the map $\psi_3 : H_3\left(\mathcal{I}_g^1\right) \to \bigwedge^3\left(\bigwedge^3 H\right)$ contains V_1 , which implies that V_1 is also contained in the image of $\psi_3 : H_3\left(\mathcal{I}_{g,1}\right)$, by the factorization $\mathcal{I}_g^1 \to \mathcal{I}_{g,1} \to \bigwedge^3 H_{\mathbb{Z}}$. To prove the result for $n \ge 4$, we will use the result for n = 2. Let us suppose that $n \ge 4$

To prove the result for $n \ge 4$, we will use the result for n = 2. Let us suppose that $n \ge 4$ and $g \ge 2n-4$. If we once again define $U := \left(\bigwedge^3 H\right)/H$, we have a commutative diagram

$$\begin{split} H_n\left(\mathcal{T}_{g,1}\right) & \longrightarrow \bigwedge^{n-2} H_1\left(\mathcal{T}_{g,1}\right) \otimes H_2\left(\mathcal{T}_{g,1}\right) & \longrightarrow \bigwedge^n H_1\left(\mathcal{T}_{g,1}\right) & \xrightarrow{\cong} & \bigwedge^n \bigwedge^3 H \\ & \downarrow^{\pi_* \otimes \mathrm{id}} & \downarrow & \downarrow^{\#_* \otimes \mathrm{id}} & \downarrow^{\#_* \otimes \mathrm{i$$

where the leftmost square commutes by the projection formula and the middle square commutes by functoriality of homology.

The composition of the arrows in the upper row is ψ_n . Since the rightmost vertical arrow is surjective, and since we now know that the image of the composition $\psi_2 \circ \pi^!$: $H_0(\mathcal{T}_g) \to \bigwedge^2 \bigwedge^3 H$ lies in weight 0, it now suffices to find a class in $H_{n-2}(\mathcal{T}_g)$ whose image in $\bigwedge^{n-2} U$ has a nonzero component in weight n-2. Note that this is impossible for n=3, since $H_1(\mathcal{T}_g) \cong U \cong V_{1^3}$, which is why this case had to be treated separately.

For $m = n - 2 \ge 2$, let us consider the image of the cycle ρ_1^m , with ρ_1 defined as in Example 4.5, under the map $H_m(\mathcal{I}_g^1) \to H_m(\mathcal{I}_{g,1})$. We also get a corresponding abelian cycle in $H_m(\mathcal{I}_g)$ by forgetting the marked point. Abusing notation, we will use ρ_1^m to denote this class in either of these homology groups. Let $\psi'_m : H_m(\mathcal{I}_g) \to \bigwedge^3 U$ be defined as in Remark 5.4. Then we have a commutative diagram

$$\begin{array}{ccc} H_m\left(\mathcal{I}_{g,1}\right) & \stackrel{\psi_m}{\longrightarrow} & \bigwedge^m \bigwedge^3 H \\ & & \downarrow & \\ & & \downarrow \\ H_m\left(\mathcal{I}_g\right) & \stackrel{\psi_m'}{\longrightarrow} & \bigwedge^m U, \end{array}$$

which means that we may work with the image of $\psi_m(\rho_1^m)$ in $\bigwedge^m U$. The map $\bigwedge^3 H \to U$ can be described explicitly, by identifying U with the image of the map $p: \bigwedge^3 H \to \bigwedge^3 H$ defined by

$$p(x \wedge y \wedge z) = x \wedge y \wedge z - \frac{1}{g-1}C_3(x \wedge y \wedge z) \wedge \omega,$$

where $C_3 : \bigwedge^3 H \to H$ is the contraction map. Note that this image is precisely the subrepresentation of traceless tensors in $\bigwedge^3 H$. With a suitable choice of basis of H, we get

$$(p^{\wedge m} \circ \psi_m)(\rho_1^m) = \left(a_1 \wedge b_1 \wedge a_{m+1} - \frac{1}{g-1}\omega \wedge a_{m+1}\right) \wedge \dots \wedge \left(a_m \wedge b_m \wedge a_{2m} - \frac{1}{g-1}\omega \wedge a_{2m}\right).$$

We want to show that this has a nonzero component in an irreducible subrepresentation of weight m, so we apply the map $\bigwedge^m \left(\bigwedge^3 H\right) \to \bigwedge^{3m} H$ and then contract m times. If we let $[m] = \{1, \ldots, n\}$, the image in $\bigwedge^{3m} H$ can be written as

$$\left(\sum_{k=0}^{m} \frac{(-1)^{k}}{(g-1)^{k}} \sum_{\{i_{1},\dots,i_{k}\}\subset[m]} \omega^{\wedge k} \wedge a_{1} \wedge b_{1} \wedge \dots \wedge \widehat{a_{i_{1}} \wedge b_{i_{1}}} \wedge \dots \wedge \widehat{a_{i_{k}} \wedge b_{i_{k}}} \wedge \dots \wedge a_{m} \wedge b_{m}\right) \wedge a_{m+1} \wedge \dots \wedge a_{2m}.$$

By contracting each term m times, we get

$$\frac{1}{(g-1)^m} \sum_{k=0}^m (-1)^k \binom{m}{k} (g-1)^{m-k} (g-2m+1)^k a_{m+1} \wedge \dots \wedge a_{2m}$$
$$= \frac{(2m-2)^m}{(g-1)^m} a_{m+1} \wedge \dots \wedge a_{2m},$$

which is nonzero, since $m = n - 2 \ge 2$, and which contracts to zero and thus lies in $V_{1^m} \subset \bigwedge^m H$. This proves the claim for $n \ge 4$.

Remark 5.5. The proof of Claim 5.2 is heavily inspired by [1, §4], where the authors compute the image of $\tau_n \circ \pi^!$ for truly nested families of bounding-pair maps. In our situation, it would be reasonable to expect the formula

$$\psi_n\left(\pi^!\pi_*A(f_1,\ldots,f_{n-2})\right) = \tau([f_1]) \wedge \cdots \wedge \tau([f_{n-2}]) \wedge \psi_2([S_g])$$

to hold for any abelian cycle $A(f_1, \ldots, f_{n-2}) \in H_{n-2}(\mathcal{I}_{g,1})$, which would prove Claim 5.2 for $n \geq 3$ directly from the case n = 2. We have not, however, been able to verify or disprove this.

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