

Accretion and wind dynamics in tidal disruption events

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Abstract. We have constructed self similar models of time dependent and non relativistic accretion disks in both sub and super-Eddington phase of TDEs with wind outflows for a general viscosity prescription which is a function of surface density of the disk Σ_d and radius r . The physical parameters are black hole (BH) mass M_\bullet , specific orbital energy E and angular momentum J , star mass M_\star and radius R_\star . We have considered an accretion disk where matter is lost due to accretion by black hole and out flowing wind (in case of super-Eddington) and added through fallback of the disrupted debris. We have simulated the light curve profiles in various spectral bands and fit them to the observations to determine the above mentioned physical parameters.

Keywords. Galaxy: nucleus, black hole physics, accretion, accretion disks

1. Introduction

A star is tidally disrupted if its pericenter $r_p \leq r_t$ where $r_t \sim R_\star(M_\bullet/M_\star)^{1/3}$ is the tidal radius and these events are called as tidal disruption events (TDEs) (Rees (1988)). Mageshwaran & Mangalam (2015) have constructed a detail stellar dynamical model of TDE using parameters that include black hole (BH) mass M_\bullet , specific orbital energy E and angular momentum J , star mass M_\star and radius R_\star . We define the dimensionless energy $\bar{e} = E/E_m$ and angular momentum $\ell = J/J_{lc}$ where J_{lc} is the loss cone angular momentum and $E_m = GM_\bullet/r_t$ is the maximum orbital energy such that $J_{lc} = 0$.

We have considered a self similar model of a time dependent and non-relativistic accretion disk for both sub-Eddington and super-Eddington disk with wind outflows due to strong radiative pressure. We have also considered a general viscosity prescription given by $\Pi_{r\phi} = -K\Sigma_d^b r^d$ where Σ_d is the surface density of the disk, r is the radius, K , b and d are constants (Mangalam (2001)). We assume a α viscous model for sub-Eddington disk with gas pressure dominating such that $b = 5/3$ and $d = -1/2$ and the radiative viscosity for super-Eddington disk with radiation pressure given by $P_r \propto \rho^\gamma$ where $\gamma = 4/3$ such that $b = 1/3$ and $d = -5/6$ (Mangalam (2003)). The non relativistic disk equations in the limit of $u_r \ll u_\phi$, where u_r is the radial velocity and u_ϕ is the azimuthal velocity, are given by $\dot{\Sigma}_d = -(1/r)\partial_r(ru_r\Sigma_d) - \dot{\Sigma}_w$; $\omega^2(r) = (1/r)\partial_r\Phi(r)$; $u_r\Sigma_d\partial_r(r^2\omega) + \dot{\Sigma}_w r^2\omega(r) = -(1/r)\partial_r(r^2\Pi_{r\phi})$, where $\dot{\Sigma}_d$ and $\dot{\Sigma}_w$ are time derivative of the surface density of disk and wind components respectively and $\dot{\Sigma}_w$ is zero for sub-Eddington disk and $\Phi(r) = -GM_\bullet/r$. For a super-Eddington disk, assuming hydrostatic equilibrium up to the height $Z = Z_{ph}$ we construct the vertical density structure ρ which is then used in the Eddington approximation to derive Z_{ph} . Using vertical momentum equation at Z_{ph} , implies $\dot{\Sigma}_w \propto \Sigma_d^e r^f \sqrt{1 - \eta/\eta_{ph}}$, where $e = 0.12$, $f = -3/2$, η and η_{ph} is the ratio of radiation pressure to total pressure at the mid plane and photosphere of the disk respectively. The self similar form of the accretion disk is

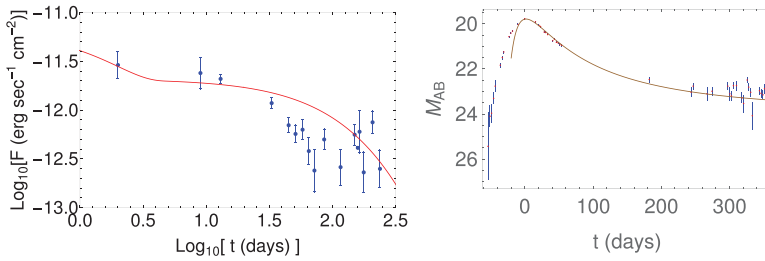


Figure 1. (Left) The sub-Eddington model fit to the XMMSL1 J061927.1-655311 (Saxton (2014)) in X-ray band. (Right) The super-Eddington model fit to the PS1-10jh (Gezari *et al.*(2012)) observations in optical g band.

taken to be $\Sigma_d = \Sigma_0(t/t_0)^\beta A \xi^p$ and $\xi = r/r_0(t/t_0)^{-\alpha}$, where Σ_0 , r_0 , t_0 , A , α , β , p are the constants derived using disk equations and the mass of disk $M_d(t_0) = \int_{t_m}^{t_0} \dot{M}_{fb}(t) dt$, where \dot{M}_{fb} and t_m are the infall rate of disrupted debris after disruption and orbital period of innermost debris respectively as given in Mageshwaran & Mangalam (2015).

We assume a seed accretion disk formed by the circularization of the debris with initial mass $M_d(t_0)$ and the derived self similar structure of Σ_d and $r_0 = q r_{in}$, where r_{in} is the inner radius taken to be innermost stable circular orbit r_{ISCO} which is a function of black hole spin j . This seed disk evolves in effect of mass addition by the infall of disrupted debris and mass loss due to accretion by black hole and wind such that the structure of Σ_d remains same. Applying mass conservation, $\dot{M}_d = \dot{M}_{fb} - \dot{M}_a - \dot{M}_w$ where \dot{M}_d , \dot{M}_a and \dot{M}_w are the rate of the change of disk mass, accretion rate by the black hole and mass loss due to out flowing wind respectively, we calculate the evolution of outer radius $r_{out}(t)$. The temperature of the sub-Eddington disk is then given by $\sigma T_e^4 = (3/8) \omega \Pi_{r\phi}$ where $\omega = \sqrt{GM_\bullet/r^3}$ and for super-Eddington disk, the temperature is derived using radiative Eddington approximation at Z_{ph} . We now discuss the light curve profiles in various spectral bands fit to the observations and estimated physical parameters.

2. Conclusions

We have obtained the self similar structure of a time dependent and non relativistic accretion disk. The parameters $\{\beta, \alpha, p\}$ are $\{-8/7, 5/21, 3/2\}$ for sub-Eddington disk and $\{-0.95, 1.22, -0.77\}$ for super-Eddington disk. The bolometric luminosity for sub-Eddington disk is $L \propto t^{-1.9}(\xi_{out}^{5/2}(t) - \xi_{in}^{5/2}(t))$ and for super-Eddington disk is $L \propto t^{1/2}(\xi_{out}^{0.4}(t) - \xi_{in}^{0.4}(t))$. The Fig 1 (Left) shows the sub-Eddington model fit to the observations in X-rays for XMMSL1 J061927.1-655311 and the derived parameters are $\bar{e} = 0.00316$, $\ell = 0.9$, $M_6 = 3.15$, $m = 1.77$, $q = 2$ and black hole spin $j = 0.15$. The Fig 1 (right) shows the super-Eddington model fit to the PS1-10jh observations in optical g band and the derived parameters are $\bar{e} = 0.0001$, $\ell = 1$, $M_6 = 36$, $m = 1.56$, $q = 2$, $j = 0.5$.

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