

In situ cryo-electron microscopy of two-dimensional van der Waals magnets

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Magnetism in two-dimensional (2D) materials and their heterostructures have been extensively studied due to the fundamental spin physics and applications for spintronics and quantum information sciences [1-3]. Recently, topologically nontrivial spin structures have been reported in various 2D van der Waals magnets, such as Cr₂Ge₂Te₆, Fe₃GeTe₂, Fe₃Sn₂, and intercalated 2H-NbS₂ [4-7]. Phase imaging in transmission electron microscopy provides a unique opportunity to study such magnetism at the nanoscale and its dynamics under external stimuli. Using off-axis electron holography and Lorentz microscopy combined with liquid helium cooling and electrical biasing, we study topological magnetic spin states in 2D magnets and their dynamic behaviors under external magnetic and/or electric fields.

Figs. 1a-d show the crystal structure of 2H-TaS₂ intercalated with Cr in 1/3 composition with respect to Ta. The Cr atoms occupy the octahedral holes made by S atoms between two neighboring TaS₂ layers, consistent with the $\sqrt{3}\times\sqrt{3}$ structure reported for the Cr_{1/3}NbS₂ [7]. In HAADF STEM image (Fig. 1c), the intercalated Cr atoms are directly observed. This Cr intercalation breaks inversion symmetry and gives rise to chiral spin arrangement along the *c*-axis due to the competition between ferromagnetic exchange interaction (FMI) and antisymmetric exchange interaction (DMI). Lorentz TEM images in Fig. 1e show the alternating bright and dark stripe patterns emerging upon cooling below $T \sim 120$ K, similar to the helical spin states in Cr_{1/3}NbS₂ [7]. However, the periodicity of helical spin states in Cr_{1/3}TaS₂ was about 15.6 nm, much shorter than 48 nm in Cr_{1/3}NbS₂, indicating stronger DMI with respect to the FMI. When an external magnetic field along the imaging direction is applied, the periodicity scales nonlinearly to form chiral soliton states (separation of dark stripes increases).

Interestingly, we observed the helicity (rotation sense) reversal occurs across a structural planar defect, as shown in Figs. 2a and b. The schematic in Fig. 2c shows the spin arrangements across the planar defects. Atomic resolution STEM image in Fig. 2d shows that the 2H phase is partly interrupted by 1T phase; note that the Ta coordinates are mixed with trigonal prismatic (2H) and octahedral (1T) in the structural defects. Interestingly, the partial 1T phase is only found every two TaS₂ layers, forming a superstructure doubling the lattice constant along the *c*-axis. Considering the presence of partial 1T phase, the DMI pathways are illustrated in Fig. 2f. For the next nearest neighbor interaction between two Cr atoms in 2H-TaS₂ phase, two S atoms are off-centered from the 2-fold rotation axis, giving rise to a net DMI along the *c*-axis [8]. However, the S atoms in 1T phase restore inversion symmetry at the center of the Cr-Cr path, vanishing the DMI. This local vanishing DMI at the planar defects effectively changes the DMI sign to reverse the helicity of helical spin states and chiral solitons in Cr_{1/3}TaS₂.

In our cryo-Lorentz phase microscopy study, we directly observed the helicity reversal behavior of the soliton lattices in Cr_{1/3}TaS₂ across structural planar defects. Using atomic resolution imaging, we

show that a partial 1T phase in 2H-TaS₂, forming a superstructure with doubling lattice constant along the *c*-axis, breaking local DMI symmetry, and leading to the spin helicity reversal. Our experimental results show that manipulation of structural phases in intercalated 2D magnets provides a viable way to control the helicity of topologically nontrivial soliton lattices.

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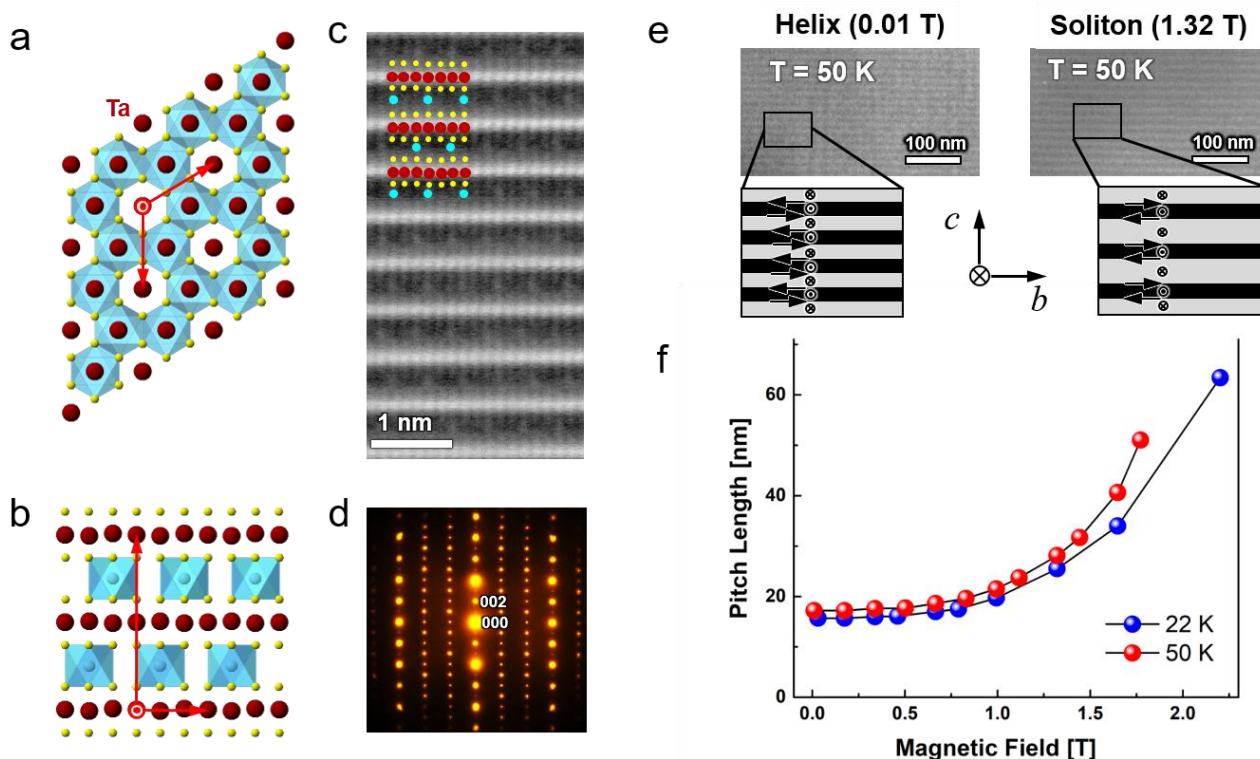


Figure 1. Figure 1| Crystal structures and helical/soliton spin states of Cr-intercalated 2H-TaS₂ Atomic models of Cr-intercalated 2H-TaS₂ viewed along the *c*-axis (a), the *a*-axis (b). HAADF STEM image taken along the *a*-axis (c) show the $\sqrt{3}\times\sqrt{3}$ Cr site-occupancy. Electron diffraction patterns taken along the *a*-axis (d) shows the superlattice peaks associated with Cr-intercalation. Lorentz TEM images (e) obtained at 50 K under various magnetic fields show helical spin states with modulation period ~ 15.6 nm, that are nonlinearly scaled with increasing magnetic field (f). External magnetic field direction is along the imaging direction (the *a*-axis).

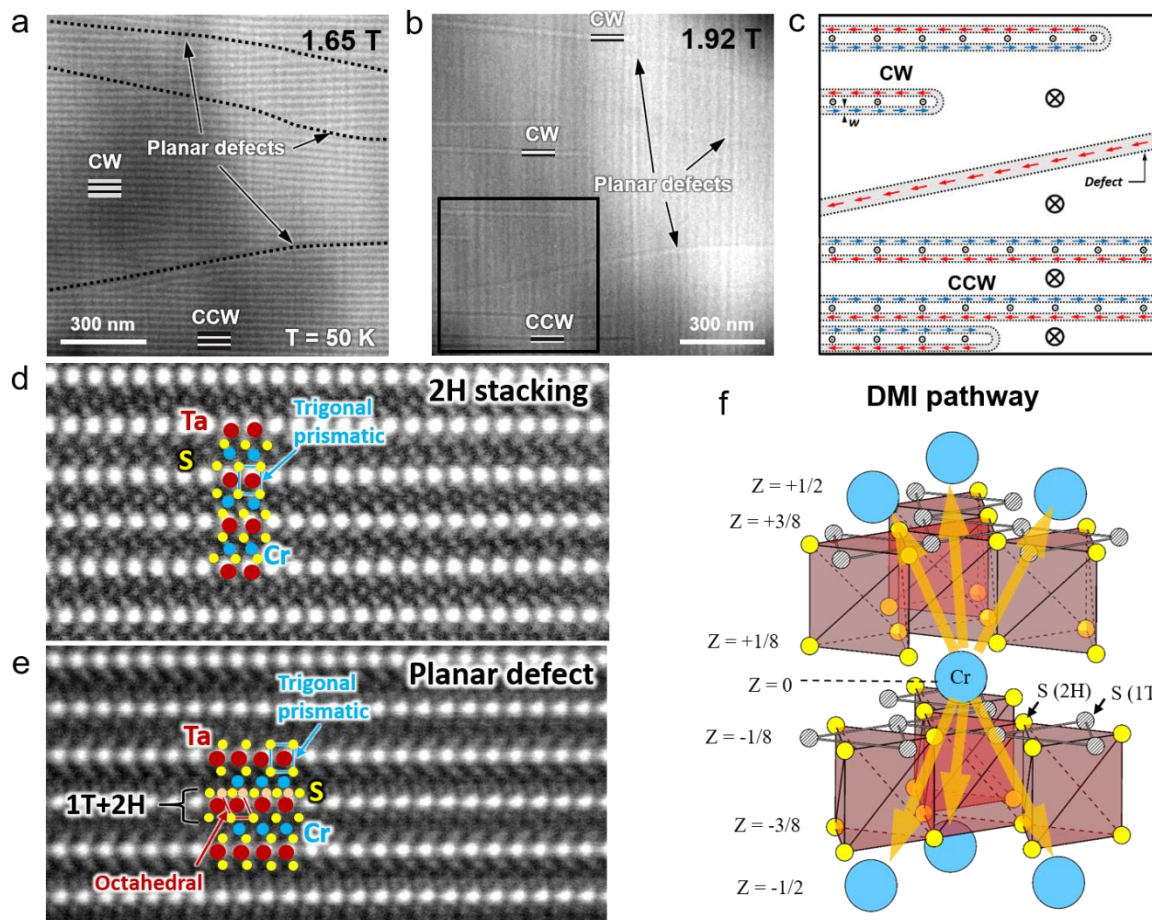


Figure 2. Figure 2| Helicity reversal across planar defects Lorentz TEM images (a-b) under high magnetic field shows reversal of spin helicity across planar defects. Schematic of spin structures (c) in the rectangular box in b. HAADF STEM image taken along the [210] direction shows 2H stacking (d). At the defect, a partial 1T mixing with 2H stacking as seen by S atomic column splitting was found. The schematic (f) shows the modified DMI pathway (six arrows), responsible for the spin helicity reversal observed in b. Ta atoms in the center of trigonal prisms are not shown in f. S atoms in the 1T phase cancel out interlayer DMI associated arising with S atoms in 2H stacking.

References

- [1] K. S. Burch, D. Mandrus, and J. Park, *Nature* **563**, 4 (2018).
- [2] B. Huang, et al., *Nature* **546**, 270 (2017).
- [3] C. Gong, et al., *Nature* **546**, 265 (2017).
- [4] M.-G. Han, et al., *Nano Lett.* **19**, 7859-7865 (2019).
- [5] T. E. Park, et al., arXiv:1907.01425 (2019).
- [6] Q. Du, et al., *Advanced Quantum Technologies* 2000058 (2020).
- [7] Y. Togawa, et al., *Physical Review Letters* **108**, 107202 (2012).
- [8] Y. Cao, et al., *Materials Today Advances* **7**, 100080 (2020).