

THE DIFFERENTIAL EMISSION MEASURE OF λ And

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ABSTRACT. Simultaneous X-ray and UV observations of λ And from EXOSAT and IUE are analyzed and the differential emission measure as a function of temperature is derived.

1 - INTRODUCTION

λ And is a well known single line spectroscopic binary with an orbital period of 20.5^d. Its very low mass function (0.006 M_{\odot}) and no evidence of eclipses indicate that the system is seen almost pole on. A photometric optical variability with a period of 54^d has been shown to be produced by a system of two large spots on the primary star which, therefore, does not rotate synchronously with the orbital motion. The primary is a G 7-8 III-IV star which has long been known to have strong chromospheric activity in the Ca II (Gratton, 1950) and Mg II lines (Baliunas and Dupree, 1982). Its chromospheric and coronal activity is also confirmed by microwave (Bath and Wallestein, 1976), X-ray (Walter and Bowyer, 1980; Vaiana et al., 1981) and UV emissions (Linsky et al. 1978). These observations indicate that λ And is a RS CVn star (Hall, 1981) and one of the brightest members of this class.

Observations of λ And with the EXOSAT satellite were included in a large program devoted to the study of activity in late type stars (Pallavicini et al., 1988). On November 11, 1985 simultaneous observations were obtained with EXOSAT LE and ME experiments and with IUE. For the purpose of the present analysis, the ME data have been binned in 3 bands (table 1).

Table 1
Observations of λ And on Nov 11, 1985

EXOSAT ME (keV)	$c s^{-1}$		EXOSAT LE	$c s^{-1}$		IUE	$c s^{-1}$	
1.17-2.14	7.54	10^{-1}	3 Lex	4.815	10^{-1}	N V	2.5	10^{-2}
2.14-3.40	8.58	10^{-1}	4 lex	3.957	10^{-1}	C II	1.23	10^{-1}
3.40-5.26	4.30	10^{-1}	Al/Pa	1.279	10^{-1}	Si IV	1.1	10^{-1}
			Boron	3.253	10^{-2}	C IV	3.11	10^{-1}

The aim of this paper is to analyze these observations in order to deduce the differential emission measure as a function of temperature. Uncertainties are assumed to be $\pm 10\%$ for the EXOSAT data and $\pm 30\%$ for the IUE data.

2 - THE DIFFERENTIAL EMISSION MEASURE DISTRIBUTION.

It is not possible to reproduce both the UV and X-ray observations with a single temperature model; this is well known in the solar case and is clearly shown by fig. 1, where for each temperature we plot the emission measure that would be needed to reproduce the observed emissions. The same figure suggests that most of the emission in the ME and LE bands must come from a "quasi isothermal" plasma with a temperature higher than $20 \cdot 10^6$ °K, while no isothermal solution appears to be possible for the UV lines.

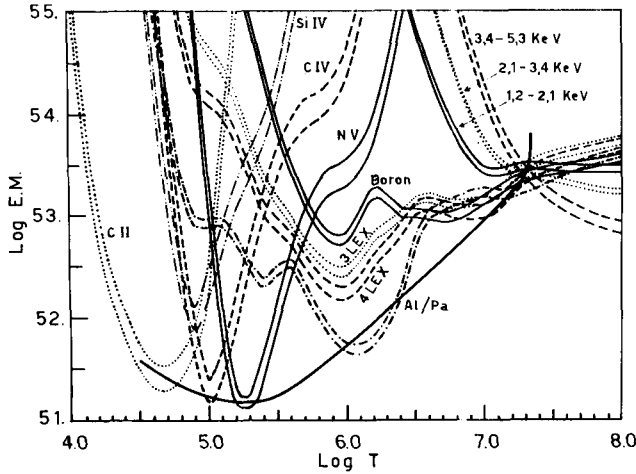


FIG. 1 - Log of the emission measure versus temperature allowed by the observed fluxes (including errors) in binned ME bands (1.17-2.14 KeV, 2.14-3.4 KeV, 3.4-5.26 KeV), LE filters (Boron, 3Lexan, 4 Lexan, Al/Pa) and UV lines (C II, C IV, N V, Si IV).

The envelope of curves shown in fig.1 suggests the shape of the differential emission measure (d.e.m.) as a function of the temperature. We assume the following "parameterized" form of d.e.m.:

$$n^2 (dV/dT) = C e^{T_0/T} (T/T_C)^\beta [1 - (T/T_C)^\gamma]^{-1/2}$$

Apart from the exponential term, which modulates the low temperature region (UV emission), this form is suggested by a theoretical investigation of the energy balance (Landini and Monsignori Fossi, 1975,1981) and has been proved to be useful also for solar active regions (Monsignori Fossi and Landini, 1988).

The plasma emissivity for each spectral region of interest has been integrated over the assumed d.e.m. and a best χ^2 fit procedure has been used to evaluate the free parameters T_0 , T_C , b , g and C ; in the present case the best fit procedure has been applied to all data in table 1 except the C II line which may be optically thick and the Al/Pa filter, which proved to give an unusually low flux (see later).

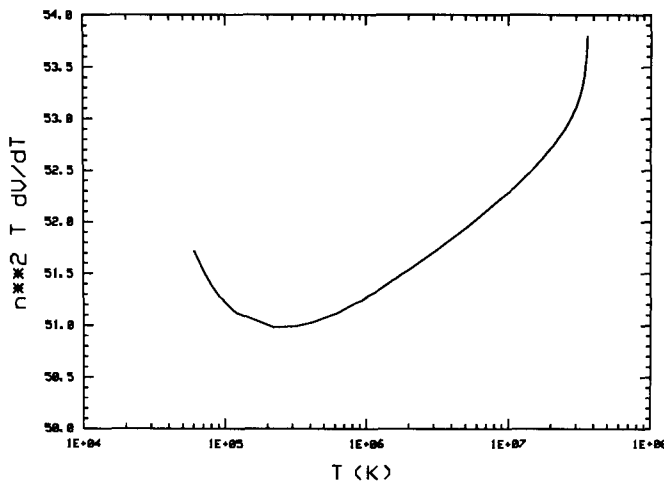


FIG. 2 - The differential emission measure distribution that allows the best χ^2 fit between predicted and observed data.

The following values give a reasonable good fit (reduced $\chi^2 = 1.2$): $\log T_c = 7.56$; $\log T_o = 5.4$; $\gamma = .64$; $\beta = .015$; $C = 1.5 \cdot 10^{45} \text{ cm}^{-3} \text{ K}^{-1}$.

The differential emission measure is plotted in fig.2 and is used to generate a synthetic spectrum (Landini and Monsignori Fossi, 1984) of the radiation emitted by the portion of the loop at temperatures larger than $6 \cdot 10^4 \text{ }^\circ\text{K}$.

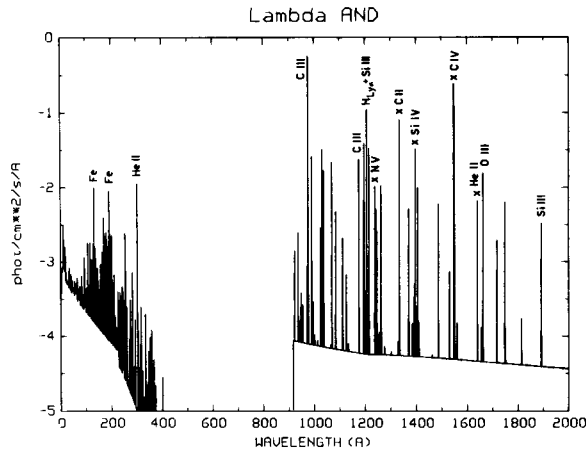


FIG. 3 - The synthetic spectrum from 1 to 2000 A obtained for λ And assuming the parameterized differential emission measure distribution shown in fig.2. The interstellar absorption, evaluated with a column density $N_H = 5.2 \cdot 10^{18} \text{ cm}^{-2}$, affects the region from 200 to 980 A.

The spectral distribution over the region 1-2000 A is shown in fig. 3: the most prominent spectral lines are indicated and crosses give the observed fluxes for the IUE lines. A few lines which have not been used in the χ^2 fit are present in the synthetic spectrum and need comments: Si IV (1397 A) is too low by a factor 3 than predicted: this fact was already noticed and is attributed to the density dependence of ion abundance (Jordan et al., 1987). He II (1640 A) in λ And is much higher than predicted: this fact can be attributed to an overpopulation of the upper level due to recombination of He nuclei following photoionization by photons with $\lambda < 227 \text{ A}$. These photons are much more numerous in λ And than in the Sun, due to the higher coronal temperature of this star (Linsky et al., 1978). Photoionization modifies the ionization balance of the He ions and this fact may also explain the disagreement of the Al/Pa signal which is affected by the strong He II 303 A line visible in the synthetic spectrum. Also a slightly different value for the interstellar absorption may change the expected Al/Pa signal.

In the Sun, transition region and coronal emission originates from loop-like structures which to first approximation can be considered of constant pressure and constant cross section. We can make the same assumption for λ And. Two large spots, covering 12% of the surface of the primary star was inferred by Bopp and Noha (1980), at about 30° in latitude and separated by 160° in longitude. If we assume that they indicate the presence of two "preferential longitudes" of activity we may postulate the existence of a large loop connecting them.

For a star radius is $5-6 R_\odot$, the two spots are about 10^{12} cm apart and this value can also be assumed as a reasonable estimate of the half length L of the loop. If this loop is similar to solar ones, an aspect ratio of 0.1 may be assumed and the loop cross-section is given by:

$$S = \pi \cdot 10^{-2} L^2 \approx 3 \cdot 10^{22} \text{ cm}^2$$

Since the projected disk area of the star is about $4.5 \cdot 10^{23} \text{ cm}^2$, the loop cross-section is somewhat less than 10% of the stellar disk, very near the value obtained from the optical observations.

Assuming a constant pressure and a constant cross-section, the semilength of the loop is given by:

$$L/2 = dh = C (4 k^2/p^2 S) \int_{T_1}^{T_c} T^2 e^{T_0/T} (T/T_c)^\beta [1 - (T/T_c)^\gamma]^{-1/2} dT.$$

With $T_1 = 2 \cdot 10^4$ °K and the best-fit parameters given above, one derives $p = 15 \text{ dyn cm}^{-2}$, using the mentioned values of S and L .

3. CONCLUSIONS

The analysis of X-ray and UV emission of λ And observed with EXOSAT (ME and LE) and IUE shows that:

- a continuous distribution of d.e.m. with temperature up to a maximum temperature $\approx 35 \cdot 10^6$ °K can reproduce (reduced $\chi^2 \approx 1.2$) all available observations. The worst fit occurs for the low energy part of the ME spectrum, which requires a somewhat larger d.e.m. around $5 \cdot 10^6$ °K, and for the Al/Pa filter, which is lower than theoretical predictions (a larger interstellar column density ?);
- the d.e.m. [$T n^2 (dV/dT)$] presents a minimum around $2 \cdot 10^5$ °K and increases at higher temperatures;
- the derived d.e.m. is compatible with a loop model having: half length = 10^{12} cm, cross-section = $3 \cdot 10^{22} \text{ cm}^2$, and pressure = 15 dyn cm^{-2} . In contrast to the model of Rosner et al.(1978) the heating deposition along this loop turns out to be non uniform.

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