

ON A RIEMANNIAN MANIFOLD M_{2n}
WITH AN ALMOST TANGENT STRUCTURE

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Professor Eliopoulous studied almost tangent structures on manifolds M_{2n} in [1; 2]. An almost tangent structure F is a field of class C^∞ of linear operations on M_{2n} such that at each point x in M_{2n} , F_x maps the complexified tangent space T_x^c into itself and that F_x is of rank n everywhere and satisfies that $F^2 = 0$. In this note, we consider a $(1, 1)$ tensor field F_i^j on a Riemannian M_{2n} which satisfies $F_i^j F_j^k = 0$ everywhere and is such that the rank of F is n everywhere. Such F_i^j gives an almost tangent structure F on M_{2n} . We are able to connect such a structure F to an almost product and almost complex structure on M_{2n} (Theorem 1). Furthermore, we study the integrability of F and obtain two results (Theorems 2 and 3). The basic materials used to prove the theorems are stated in §1 and §2, which were developed in Yano [5].

1. Let M_{2n} be a Riemannian manifold of class C^∞ . We assume that a field of class C^∞ of linear operations F_x is defined on M_{2n} such that at each point $x \in M_{2n}$, F_x maps the tangent space T_x of M_{2n} at x into itself; moreover, F_x is of rank n everywhere in M_{2n} , and it satisfies the relation

$$F_x^2 = 0$$

for every $x \in M_{2n}$. Let the image $F_x(T_x)$ be B_x ; then B_x is a

distribution of dimension n on M_{2n} . Let C_x be the orthogonal distribution to B_x with respect to the Riemann metric. Thus

$T_x = B_x \oplus C_x$. Since $F^2 = 0$, we have $F_x(B_x) = 0$ and $B_x = F_x(T_x) = F_x(B_x \oplus C_x) = F_x(C_x)$. In any local coordinate neighborhood at the point x , one can write the operator F_x and distributions B_x, C_x in the tensorial notations:

$$\begin{aligned} F_i^j & \quad (i, j = 1, 2, \dots, 2n); \\ B_\alpha^i & \quad (i = 1, \dots, 2n; \alpha = 1, 2, \dots, n); \\ C_{\bar{\alpha}}^i & \quad (i = 1, \dots, 2n; \bar{\alpha} = \alpha + n). \end{aligned}$$

One can always arrange that

$$(0) \quad F(C_{\bar{\alpha}}) = B_\alpha \quad \text{or} \quad F_j^i C_{\bar{\alpha}}^j = B_\alpha^i,$$

where here and in the following Greek indices $\alpha, \beta, \gamma, \dots$ run over the range $1, 2, \dots, n$; $\bar{\alpha} = \alpha + n$, and Latin indices i, j, k, \dots run over the range $1, 2, \dots, 2n$.

The matrix $(B_\alpha^i, C_{\bar{\alpha}}^i)$ has an inverse, say $\begin{pmatrix} B_i^\alpha \\ C_i^{\bar{\alpha}} \end{pmatrix}$ then

$$(1) \quad B_\alpha^j B_j^\beta = \delta_\alpha^\beta, \quad C_{\bar{\alpha}}^i B_i^\beta = 0, \quad B_\beta^i C_i^{\bar{\alpha}} = 0, \quad C_{\bar{\alpha}}^i C_i^{\bar{\beta}} = \delta_\alpha^\beta,$$

and

$$(2) \quad B_\alpha^j B_i^\alpha + C_{\bar{\alpha}}^j C_i^{\bar{\alpha}} = \delta_i^j.$$

Here and in the following $B_\alpha^j B_i^\alpha = \sum_{\alpha=1}^n B_\alpha^j B_i^\alpha$, $C_{\bar{\alpha}}^j C_i^{\bar{\alpha}} = \sum_{\bar{\alpha}=n+1}^{2n} C_{\bar{\alpha}}^j C_i^{\bar{\alpha}}$.

Put

$$B_i^j = B_i^\alpha B_\alpha^j, \quad C_i^j = C_i^{\bar{\alpha}} C_{\bar{\alpha}}^j,$$

then

$$(3) \quad B_i^j + C_i^j = \delta_i^j, \quad B_i^j B_j^h = B_i^h, \quad C_i^j C_j^h = C_i^h.$$

The following equalities are easy to verify:

$$(4) \quad \begin{cases} B_i^h B_\alpha^i = B_\alpha^h, & B_i^h C_{\bar{\alpha}}^i = 0, \\ C_i^h B_\alpha^i = 0, & C_i^h C_{\bar{\alpha}}^i = C_{\bar{\alpha}}^h; \end{cases}$$

$$(5) \quad \begin{cases} B_i^h B_h^\alpha = B_i^\alpha, & B_i^h C_h^{\bar{\alpha}} = 0, \\ C_i^h B_h^\alpha = 0, & C_i^h C_h^{\bar{\alpha}} = C_i^{\bar{\alpha}}. \end{cases}$$

Let F, B, C be the operators on T_x by the tensors F_i^j, B_i^j, C_i^j ; then it is easy to see that $FB = 0, BF = F, CF = 0$ and $FC = F$. That is:

$$(6) \quad \begin{cases} B_j^i F_h^j = F_h^i, & F_j^i B_h^j = 0, \\ C_j^i F_h^j = 0, & F_j^i C_h^j = F_h^i. \end{cases}$$

2. Let

$$(7) \quad A_{ji} = B_j^\alpha B_i^\alpha + C_j^{\bar{\alpha}} C_i^{\bar{\alpha}} = \sum_{\alpha=1}^n B_j^\alpha B_i^\alpha + \sum_{\bar{\alpha}=n+1}^{2n} C_j^{\bar{\alpha}} C_i^{\bar{\alpha}}.$$

then A_{ji} is symmetric and

$$(8) \quad \begin{cases} A_{ji} B_{\alpha}^j B_{\beta}^i = \delta_{\alpha\beta}, & A_{ji} C_{\bar{\alpha}}^j C_{\bar{\beta}}^i = \delta_{\bar{\alpha}\bar{\beta}}, & A_{ji} B_{\alpha}^j C_{\bar{\beta}}^i = 0; \\ A_{ji} B_{\alpha}^i = B_j^{\alpha}, & A_{ji} C_{\bar{\alpha}}^i = C_j^{\bar{\alpha}}. \end{cases}$$

These identities show that (A_{ji}) is non-singular and that $(B_{\alpha}^i, C_{\bar{\alpha}}^i)$ form an orthonormal frame with respect to A_{ij} .

Let

$$(9) \quad B_{ji} = B_j^h A_{hi}, \quad C_{ji} = C_j^h A_{hi}.$$

Then

$$(10) \quad B_{ji} = B_j^{\alpha} B_i^{\alpha}, \quad C_{ji} = C_j^{\bar{\alpha}} C_i^{\bar{\alpha}}, \quad B_{ji} + C_{ji} = A_{ji},$$

B_{ji}, C_{ji} are both symmetric and

$$(11) \quad B_j^t B_i^s A_{ts} = B_{ji}, \quad B_j^t C_i^s A_{ts} = 0, \quad C_j^t C_i^s A_{ts} = C_{ji}.$$

Consider now the tensor G_{ij} defined by:

$$G_{ji} = \frac{1}{2}(A_{ji} + F_j^t F_i^s A_{ts} + B_{ji}).$$

Noticing that $(F_j^t B_t^{\alpha} - C_j^{\bar{\alpha}}) B_{\beta}^j = 0$, $(F_j^t B_t^{\alpha} - C_j^{\bar{\alpha}}) C_{\bar{\beta}}^j = 0$ and $(B_{\beta}^j, C_{\bar{\beta}}^j)$ is a basis for T_x , we have that $F_j^t B_t^{\alpha} = C_j^{\bar{\alpha}}$. Then by (0), (6) and (8) one has

$$(12) \quad G_{ji} B_{\alpha}^i = B_j^{\alpha}, \quad G_{ji} C_{\bar{\alpha}}^i = C_j^{\bar{\alpha}}, \quad G_{ti} B_j^t = B_{ji}.$$

The first two relations of (12) show that (G_{ij}) is non-singular. Let (G^{ij}) be the inverse matrix of (G_{ji}) , then

$$(12') \quad G^{ij} B_j^\alpha = B_\alpha^i, \quad G^{ij} C_j^{\bar{\alpha}} = C_{\bar{\alpha}}^i.$$

Originally $B_i^\alpha, C_i^{\bar{\alpha}}$ were defined in (1) from the matrix $(B_\alpha^i, C_{\bar{\alpha}}^i)$, now we have found a Riemannian metric G_{ij} so that they are related by (12) and (12').

From (7), (10), (11) and (12) we have

$$(13) \quad G_{st} B_i^s B_j^t = B_{ij}, \quad G_{st} B_i^s C_j^t = 0, \quad G_{st} C_i^s C_j^t = C_{ij}.$$

Finally, we consider the operations F_i^j and $C_i^{\bar{\gamma}} B_\gamma^j = \sum_{\gamma=1}^n C_i^{\gamma+n} B_\gamma^j$ on T_x , these two coincide because

$$F_i^j B_\alpha^i = 0, \quad F_i^j C_{\bar{\alpha}}^i = B_\alpha^j; \quad (C_i^{\bar{\gamma}} B_\gamma^j) B_\alpha^i = 0, \quad (C_i^{\bar{\gamma}} B_\gamma^j) C_{\bar{\alpha}}^i = B_\alpha^j.$$

Thus we have

$$(14) \quad F_i^k = C_i^{\bar{\gamma}} B_\gamma^k.$$

If we defined $F_{ik} = G_{kj} F_i^j$ then by (12):

$$(15) \quad F_{ik} = C_i^{\bar{\gamma}} B_k^\gamma = \sum_{\gamma=1}^n C_i^{\gamma+n} B_k^\gamma.$$

3. Let

$$\psi_{ik} = C_i \bar{\gamma} B_k^\gamma + C_k \bar{\gamma} B_i^\gamma = F_{ik} + F_{ki}$$

$$\varphi_{ik} = C_i \bar{\gamma} B_k^\gamma - C_k \bar{\gamma} B_i^\gamma = F_{ik} - F_{ki},$$

then

$$\psi_i^k \stackrel{\text{def.}}{=} G^{kj} \psi_{ij} = C_i \bar{\gamma} B_\gamma^k + B_i^\gamma C_\gamma^k,$$

$$\varphi_i^k \stackrel{\text{def.}}{=} G^{kj} \varphi_{ij} = C_i \bar{\gamma} B_\gamma^k - B_i^\gamma C_\gamma^k,$$

and

$$\psi_i^k \psi_k^j = (C_i \bar{\gamma} B_\gamma^k + B_i^\gamma C_\gamma^k) (C_k \bar{\delta} B_\delta^j + B_k^\delta C_\delta^j)$$

$$= C_i \bar{\gamma} C_\gamma^j + B_i^\gamma B_\gamma^j = C_i^j + B_i^j = \delta_i^j,$$

$$\varphi_i^k \varphi_k^j = (C_i \bar{\gamma} B_\gamma^k - B_i^\gamma C_\gamma^k) (C_k \bar{\delta} B_\delta^j - B_k^\delta C_\delta^j)$$

$$= -C_i \bar{\gamma} C_\gamma^j - B_i^\gamma B_\gamma^j = -\delta_i^j.$$

Thus ψ is an almost product structure and φ is an almost complex structure on M_{2n} . Furthermore,

$$\psi_i^k \psi_k^j = C_i^j - B_i^j, \quad \psi_i^k \varphi_k^j = -C_i^j + B_i^j,$$

that is,

$$(16) \quad \varphi_i^k \psi_k^j = -\psi_i^k \varphi_k^j.$$

Conversely, if on M_{2n} there are almost product structure ψ and almost complex structure φ which satisfy (16), then if we define

$$F_i^k = \frac{1}{2}(\varphi_i^k + \psi_i^k),$$

F_i^k is a tensor field on M_{2n} which satisfies $F^2 = 0$. Thus we have proved the following theorem.

THEOREM 1. On M_{2n} there is a (1.1) tensor field F satisfying $F^2 = 0$ if and only if there are almost product structure ψ and almost complex structure φ which satisfy $\varphi\psi = -\psi\varphi$.

Let

$$H_i^k = B_i^\gamma C_{\bar{\gamma}}^k \quad \text{or} \quad H_i^k = \frac{1}{2}(\psi_i^k - \varphi_i^k);$$

then H satisfies the following :

$$(17) \quad \begin{cases} H_i^k H_k^j = 0 \\ H_i^k B_k^\delta = 0, \quad H_i^k C_k^{\bar{\alpha}} = B_i^\alpha, \end{cases}$$

and

$$(18) \quad \begin{cases} F + H = \psi, \quad F - H = \varphi, \\ H_i^k F_k^j + F_i^k H_k^j = \delta_i^j. \end{cases}$$

Thus H is again a tensor field on M_{2n} which satisfies $H^2 = 0$.

For the structure (ψ, φ) satisfying (16), Hsu [3, page 441, Corollary 2.4] showed that there is an affine connection Γ so that

$\nabla_k \varphi_i^j = 0, \quad \nabla_k \psi_i^j = 0$ where ∇ is the covariant differentiation with

respect to Γ . In fact, let $\overset{\circ}{\nabla}$ be the covariant differentiation with respect to any given affine connection $\overset{\circ}{\Gamma}$, then one of such affine connection is the following:

$$\Gamma_{ji}^h = \overset{\circ}{\Gamma}_{ji}^h - \frac{1}{2} (\overset{\circ}{\nabla}_j \varphi_i^a) \varphi_a^h + \frac{1}{4} (\overset{\circ}{\nabla}_j \psi_i^a) \psi_a^h - \frac{1}{4} \varphi_i^b [(\overset{\circ}{\nabla}_j \psi_b^d) \psi_d^a] \varphi_a^h.$$

Such a connection, satisfying $\nabla \varphi = 0$, $\nabla \psi = 0$, is called a (φ, ψ) connection.

Hence we have the following corollary:

COROLLARY 1. There is an affine connection for which $\nabla_k F_i^j = 0$, $\nabla_k H_i^j = 0$.

Furthermore, by Hsu [3, page 415, Theorem 3.1] with respect to the (φ, ψ) connection Γ , the holonomy group can be represented as the form $\begin{pmatrix} A_n & 0 \\ 0 & A_n \end{pmatrix}$ referred to a suitable basis where A_n is any $n \times n$ matrix. Thus we have the following corollary:

COROLLARY 2. On M_{2n} with a $(1, 1)$ tensor field F such that $F^2 = 0$, with respect to the connection given in Corollary 1, the holonomy group can be represented as the form $\begin{pmatrix} A_n & 0 \\ 0 & A_n \end{pmatrix}$ referred to a suitable basis.

4. The Nijenhuis tensor formed with F_i^k is defined by:

$$t_{jk}^i(F) = F_j^\ell (\partial_k F_\ell^i - \partial_\ell F_k^i) - F_k^\ell (\partial_j F_\ell^i - \partial_\ell F_j^i).$$

Definition. The tensor field F is called integrable if $t_{jk}^i(F) = 0$.

Now we shall calculate $t_{jk}^i(F)$ by use of $F_i^k = C_i^{\bar{\gamma}} B_\gamma^k$ and (1).

$$F_j^\ell (\partial_k F_\ell^i - \partial_\ell F_k^i) = F_j^\ell (\partial_k C_\ell^{\bar{\gamma}} - \partial_\ell C_k^{\bar{\gamma}}) B_\gamma^i - C_j^{\bar{\delta}} B_\delta^\ell C_k^{\bar{\gamma}} \partial_\ell B_\gamma^i.$$

$$(19) \left\{ \begin{aligned} t_{jk}^i(F) &= F_j^\ell (\partial_k F_\ell^i - \partial_\ell F_k^i) - F_k^\ell (\partial_j F_\ell^i - \partial_\ell F_j^i) \\ &= C_j^{\bar{\delta}} B_\delta^\ell (\partial_k C_\ell^{\bar{\gamma}} - \partial_\ell C_k^{\bar{\gamma}}) B_\gamma^i - C_k^{\bar{\delta}} B_\delta^\ell (\partial_j C_\ell^{\bar{\gamma}} - \partial_\ell C_j^{\bar{\gamma}}) B_\gamma^i \\ &\quad - C_j^{\bar{\delta}} B_\delta^\ell C_k^{\bar{\gamma}} \partial_\ell B_\gamma^i + C_k^{\bar{\delta}} B_\delta^\ell C_j^{\bar{\gamma}} \partial_\ell B_\gamma^i. \end{aligned} \right.$$

$$t_{jk}^i B_\alpha^j = -C_k^{\bar{\delta}} B_\alpha^j B_\delta^\ell (\partial_j C_\ell^{\bar{\gamma}} - \partial_\ell C_j^{\bar{\gamma}}) B_\gamma^i$$

$$t_{jk}^i B_\alpha^j B_i^\beta C_\gamma^k = -B_\alpha^j B_\gamma^\ell (\partial_j C_\ell^{\bar{\beta}} - \partial_\ell C_j^{\bar{\beta}})$$

$$\stackrel{\text{def.}}{=} -\Omega_{\alpha\gamma}^{\bar{\beta}}.$$

For $\Omega_{\alpha\gamma}^{\bar{\beta}}$ there was the following consideration according to Yano [4].

An arbitrary contravariant vector dx^i in the tangent space $T_x(M_{2n})$ at x can be written in

$$(20) \quad dx^i = B_\alpha^i dy^\alpha + C_\alpha^i dy^{\bar{\alpha}}$$

where

$$(20') \quad dy^\alpha \stackrel{\text{def.}}{=} B_j^\alpha dx^j, \quad dy^{\bar{\alpha}} \stackrel{\text{def.}}{=} C_j^{\bar{\alpha}} dx^j.$$

Thus the distribution B is defined by

$$(21) \quad dy^{\bar{\alpha}} = C_j^{\bar{\alpha}} dx^j = 0.$$

The condition for the distribution B to be completely integrable is that

$$(\partial_k C_j^{\bar{\alpha}} - \partial_j C_k^{\bar{\alpha}}) dx^k \wedge dx^j = 0$$

be satisfied by any dx^i satisfying (21), that is

$$\Omega_{\alpha\gamma}^{\bar{\beta}} = B_{\alpha}^j B_{\gamma}^{\ell} (\partial_j C_{\ell}^{\bar{\beta}} - \partial_{\ell} C_j^{\bar{\beta}}) = 0.$$

Using this fact and the relation between t_{jk}^i and $\Omega_{\alpha\gamma}^{\bar{\beta}}$ given above, one has proved Theorem 2.

THEOREM 2. If F is integrable, then the distribution B is completely integrable; if H is integrable, then the distribution C is completely integrable.

The last part of this theorem is a parallel conclusion of the first part.

Conversely, if both distributions B and C are completely integrable, then there are functions $y^{\alpha}(x^i)$, $y^{\bar{\alpha}}(x^i)$ satisfying (20), (20') and hence

$$B_j^{\alpha} = \frac{\partial y^{\alpha}}{\partial x^j}, \quad C_j^{\bar{\alpha}} = \frac{\partial y^{\bar{\alpha}}}{\partial x^j}, \quad B_{\alpha}^i = \frac{\partial x^i}{\partial y^{\alpha}}, \quad C_{\bar{\alpha}}^i = \frac{\partial x^i}{\partial y^{\bar{\alpha}}}.$$

Noticing that $\partial_k C_{\ell}^{\bar{\gamma}} = \partial_k \partial_{\ell} y^{\bar{\alpha}}$ and $B_{\delta}^{\ell} \partial_{\ell} B_{\gamma}^i = \partial_{\delta} B_{\gamma}^i = \partial_{\delta} \partial_{\gamma} x^i$, by (19) we have

$$\begin{aligned} t_{jk}^i(F) &= -C_j^{\delta} B_{\delta}^{\ell} C_k^{\bar{\gamma}} \partial_{\ell} B_{\gamma}^i + C_k^{\delta} B_{\delta}^{\ell} C_j^{\bar{\gamma}} \partial_{\ell} B_{\gamma}^i \\ &= -C_j^{\delta} C_k^{\bar{\gamma}} \partial_{\delta} B_{\gamma}^i + C_k^{\delta} C_j^{\bar{\gamma}} \partial_{\delta} B_{\gamma}^i = 0. \end{aligned}$$

Thus F is integrable. Similarly, it is easy to show that H is also integrable. Hence we have Theorem 3.

THEOREM 3. If B and C are both completely integrable, then F and H are both integrable.

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