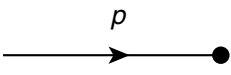
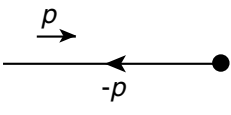


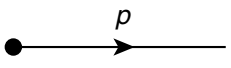
Appendix E

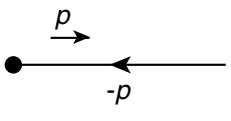
Feynman rules

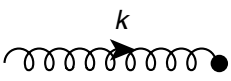
E.1 Factors induced by external or internal lines

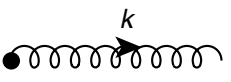
ingoing quark:  $(2\pi)^{-3/2}u(p, \lambda)$

ingoing antiquark:  $(2\pi)^{-3/2}\bar{v}(p, \lambda)$

outgoing quark:  $(2\pi)^{-3/2}\bar{u}(p, \lambda)$

outgoing antiquark:  $(2\pi)^{-3/2}v(p, \lambda)$

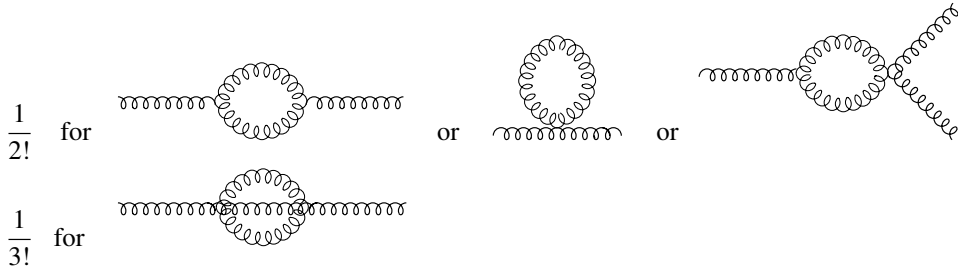
ingoing gluon:  $(2\pi)^{-3/2}\epsilon^\mu(k, \eta)$

outgoing gluon:  $(2\pi)^{-3/2}\epsilon_\mu^*(k, \eta)$

E.2 Factors induced by closed loops

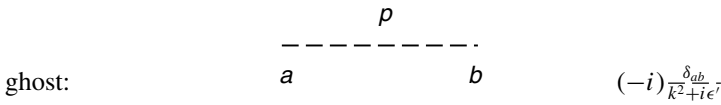
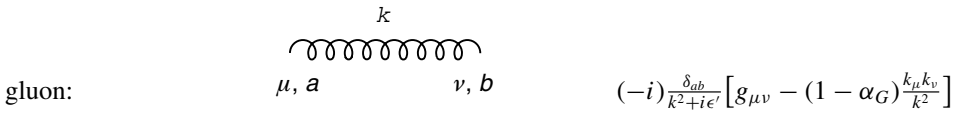
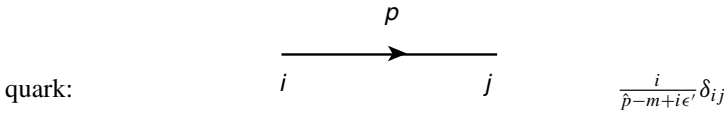
$$\int \frac{d^n p}{(2\pi)^n} \text{ for each loop integration}$$

(-1) for each closed fermion or ghost loop

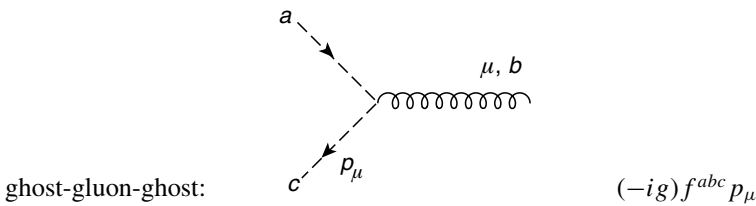
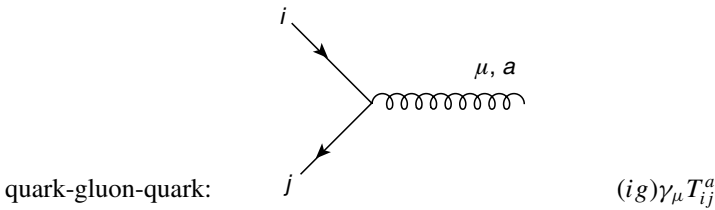


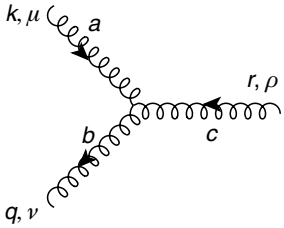
E.3 Propagators and vertices

Propagators

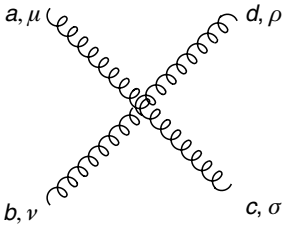


Vertices



3-gluon: 

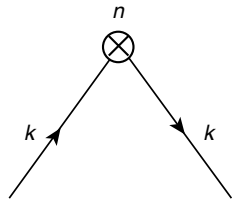
$$(gf^{abc})[g_{\mu\nu}(k+q)_\rho - g_{\nu\rho}(q+r)_\mu + g_{\rho\mu}(r-k)_\nu]$$

4-gluon: 

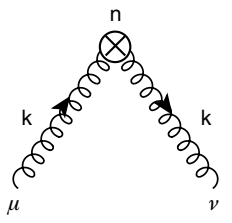
$$(-ig^2)[f^{abe}f^{cde}(g^{\mu\sigma}g^{\nu\rho} - g^{\mu\rho}g^{\nu\sigma}) + f^{ace}f^{bde} \times (g^{\mu\nu}g^{\sigma\rho} - g^{\mu\rho}g^{\nu\sigma}) + f^{ade}f^{cbe} \times (g^{\mu\sigma}g^{\nu\rho} - g^{\mu\nu}g^{\sigma\rho})]$$

E.4 Composite operators in deep-inelastic scattering

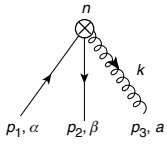
We define $\Gamma \equiv 1$ or γ_5 and Δ to be an arbitrary four-vector with $\Delta^2 = 0$. The composite operators are defined at $x = 0$.



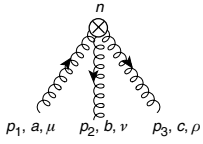
$$: \bar{q} \gamma_{\mu_1} \cdots \partial_{\mu_n} q : \quad \hat{\Delta}(\Delta \cdot k)^{n-1} \Gamma$$



$$: G_{\mu\mu_1} \partial_{\mu_2} \cdots \partial_{\nu} G : \quad g_{\mu\nu}(\Delta \cdot k)^n + k^2 \Delta_\mu \Delta_\nu (\Delta \cdot k)^{n-2} - (k_\mu \Delta_\nu + k_\nu \Delta_\mu)(\Delta \cdot k)^{n-1}$$



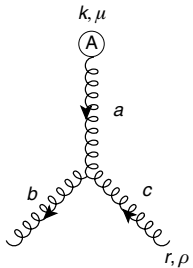
$$: \bar{q}_\alpha \gamma_{\mu_1} \cdots g B_a^\mu T_{ij}^a \cdots \gamma_{\mu_n} \Gamma q_\beta : \quad g T_{\alpha\beta}^a \Delta^\mu \hat{\Delta} \sum_{j=0}^{n-2} (\Delta \cdots p_1)^j \times (\Delta \cdots p_2)^{n-j-2} \Gamma$$



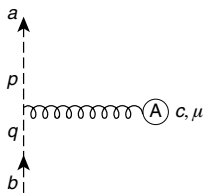
$$: G_{\mu\mu_1} \partial_{\mu_2} \cdots g B_{\mu_i} \cdots G_{\mu_n \nu} : \quad \frac{ig}{3!} f^{abc} \left\{ \Delta_n u [\Delta_\rho p_{3,\mu} (\Delta \cdots p_1) + p_{1,\rho} \Delta_\mu (\Delta \cdot p_3) - g_{\mu\rho} \times (\Delta \cdot p_1) (\Delta \cdot p_3) - \Delta_\mu \Delta_\rho \times (p_3 \cdot p_1)] + \sum_{j=1}^{n-2} (-1)^j \times (\Delta \cdot p_1)^{j-1} (\Delta \cdot p_3)^{n-j-2} + g_{\mu\rho} \Delta_\nu - g_{\nu\rho} \Delta_\mu (\Delta \cdot p_3) + \Delta_\rho (\Delta_\mu p_{3,\nu} - p_{3,\mu} \Delta_\nu) \right] \times (\Delta \cdot p_3)^{n-2} + \text{perm.} \}$$

E.5 Rules in the background field approach

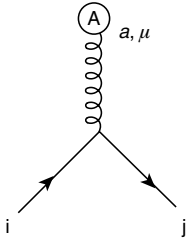
The background field is represented by *A*. The combinations of gauge fields not shown below vanish. For instance, there is no quadrilinear vertices with three or four background fields. We use the conventions in [127].



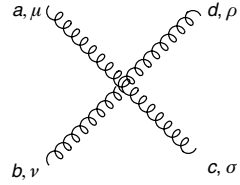
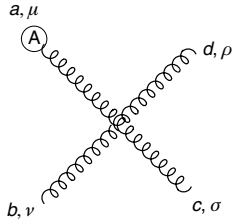
$$(g f^{abc}) \left[g_{\mu\nu} \left(k + q - \frac{r}{\alpha_G} \right)_\rho - g_{\nu\rho} (q + r)_\mu + g_{\rho\mu} \left(r - k - \frac{q}{\alpha_G} \right)_\nu \right]$$



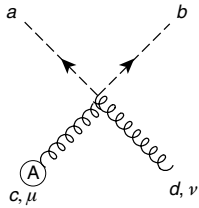
$$(-g) f^{bca} (p + q)_\mu$$



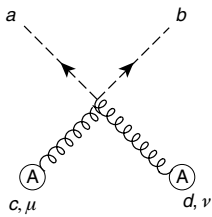
$$(ig)T_{ij}^{(a)} \gamma_\mu$$



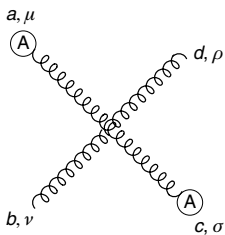
≡



$$(-ig^2)f^{ace}f^{edb}g_{\mu\nu}$$



$$(-ig^2)(f^{ace}f^{edb} + f^{adx}f^{xcb})g_{\mu\nu}$$



$$(-ig^2)[f^{abe}f^{cde}(g^{\mu\sigma}g^{\nu\rho} - g^{\mu\rho}g^{\nu\sigma} + \frac{1}{\alpha_G}g^{\mu\nu}g^{\sigma\rho}) + f^{ace}f^{bde} \times (g^{\mu\nu}g^{\sigma\rho} - g^{\mu\rho}g^{\nu\sigma}) + f^{ade}f^{cbe} \times (g^{\mu\sigma}g^{\nu\rho} - g^{\mu\nu}g^{\sigma\rho} + \frac{1}{\alpha_G}g^{\mu\rho}g^{\nu\sigma})]$$