

Evolution of magnetic fields in the IGM: kinetic MHD turbulence

Reinaldo S. de Lima¹, E. M. de Gouveia Dal Pino¹, A. Lazarian²
and D. Falceta-Gonçalves³

¹Instituto de Astronomia, Geofísica e C. Atmosféricas, Universidade de São Paulo, Brazil
email: rlima@astro.iag.usp.br, dalpino@astro.iag.usp.br

²Núcleo de Astrofísica Teórica, Universidade Cruzeiro do Sul, Brazil

³Astronomy Department, University of Wisconsin, USA

Abstract. In this work, we present 3D MHD simulations of non-helical, forced turbulence, with an anisotropic thermal pressure with respect to the orientation of the local magnetic field. Such anisotropy arises when the plasma is weakly collisional, i.e., when the Larmor frequency is much greater than the ion-ion collision frequency. In this Kinetic MHD regime (KMHD), there are instabilities that give rise to fast growing magnetic fluctuations in the smallest scales. The plasma that fills the intergalactic and intracluster media has small density ($n \sim 10^{-3} \text{ cm}^{-3}$), hence the effects of these instabilities could be important in the turbulent amplification of the magnetic fields there. In order to study the KMHD turbulence, we have performed 3D numerical simulations employing a godunov-MHD code (e.g., Kowal, Lazarian & Beresnyak 2007; Falceta-Gonçalves, Lazarian & Kowal 2008). The power spectrum of the velocity and magnetic fields were calculated for two cases: when there is a pre-existing mean magnetic field, and when there is only an initial weak magnetic field.

Keywords. Magnetic fields – plasmas – intergalactic medium

Kinetic MHD description

Due the low density, an MHD description of the IGM and ICM is not adequate: the ions Larmor radius r_L is much smaller than its mean free path λ . For instance, in the Hydra A cluster, $r_L \sim 10^5 \text{ km}$ and $\lambda \sim 10^{15} \text{ km}$ (Ensslin & Vogt 2006). A collisionless plasma can be described by an MHD set of equations under the kinematic approximation (the KMHD equations), where the scalar pressure in the ideal MHD equations is replaced by the following tensor:

$$P_{ij} = p_{\perp} \delta_{ij} + (p_{\parallel} - p_{\perp}) B_i B_j / B^2 \quad (0.1)$$

where p_{\parallel} and p_{\perp} are the components of the pressure in the directions parallel and perpendicular to \mathbf{B} . The standard approach for each pressure component is the double-adiabatic law:

$$\frac{d}{dt} \left(\frac{p_{\perp}}{\rho B} \right) = 0 \quad \frac{d}{dt} \left(\frac{p_{\parallel} B^2}{\rho^3} \right) = 0 \quad (0.2)$$

However, in the preliminary numerical simulations shown below, we have adopted an isothermal description for the pressure components in each direction, with sound velocities c_{perp} and c_{par} .

In all simulations, the turbulence is forced in a 3D periodic box (with 128^3 resolution) by a random, non-helical, solenoidal force acting around the forcing wave number k_f . The velocity of the forcing scale is $u_f \sim 1$.

Turbulence in the presence of an external magnetic field: In these simulations, the external magnetic field \mathbf{B}_{ext} is uniform with Alfvénic Mach number $M_A = 1$ (see Fig. 1).

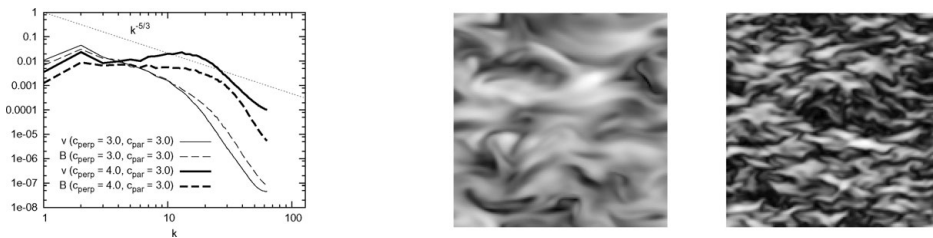


Figure 1. Left: power spectrum of the velocity and magnetic fields for a collisional MHD simulation with $c_{perp} = c_{par} = 3.0$ (MHD) and for a collisionless kinetic MHD simulation $c_{perp} = 4.0, c_{par} = 3.0$ (KMHD). The images in the right show $|\mathbf{B}|$ in the central slices of the cubes, with horizontal \mathbf{B}_{ext} (middle: MHD, right: KMHD). In these simulations $k_f = 2.5$.

Growth of a weak initial magnetic field: For different degrees of pressure anisotropy and the same initial weak magnetic field (see Fig. 2).

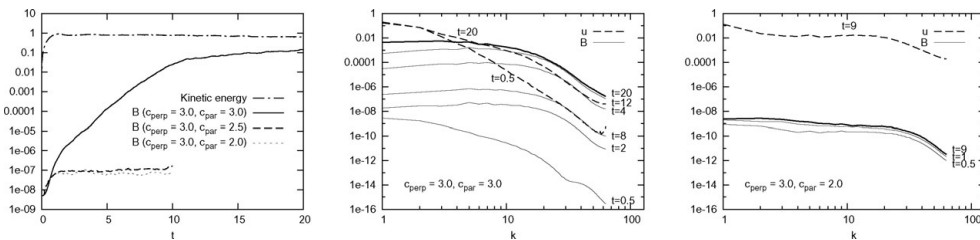


Figure 2. Left: Magnetic energy evolution. Middle and right: magnetic power spectrum evolution for the MHD case and KMHD case, respectively. In these simulations $k_f = 1.5$.

The instabilities that develop due to the anisotropic pressure accumulate energy at the smallest scales (limited by the numerical dissipation), changing the inclination of the spectrum in the inertial range. The small-scale fluctuations of the fields due to these instabilities give rise to a more “wrinkled” field distribution in the KMHD case.

In the simulations with anisotropic pressure, the dynamo action seems to be inhibited. This issue however, still requires deeper investigation with simulations spanning a larger parameter space.

Further simulations including a double-adiabatic law must be performed in order to study the turbulence development with spontaneously raised anisotropic pressures.

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References

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